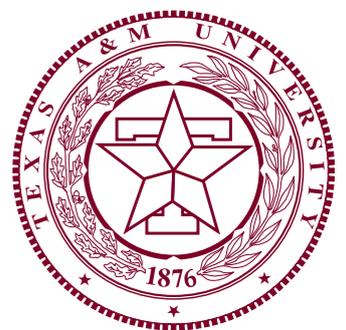


Submesoscale eddies along the Mississippi/Atchafalaya plume front

Rob Hetland

Texas A&M University, USA

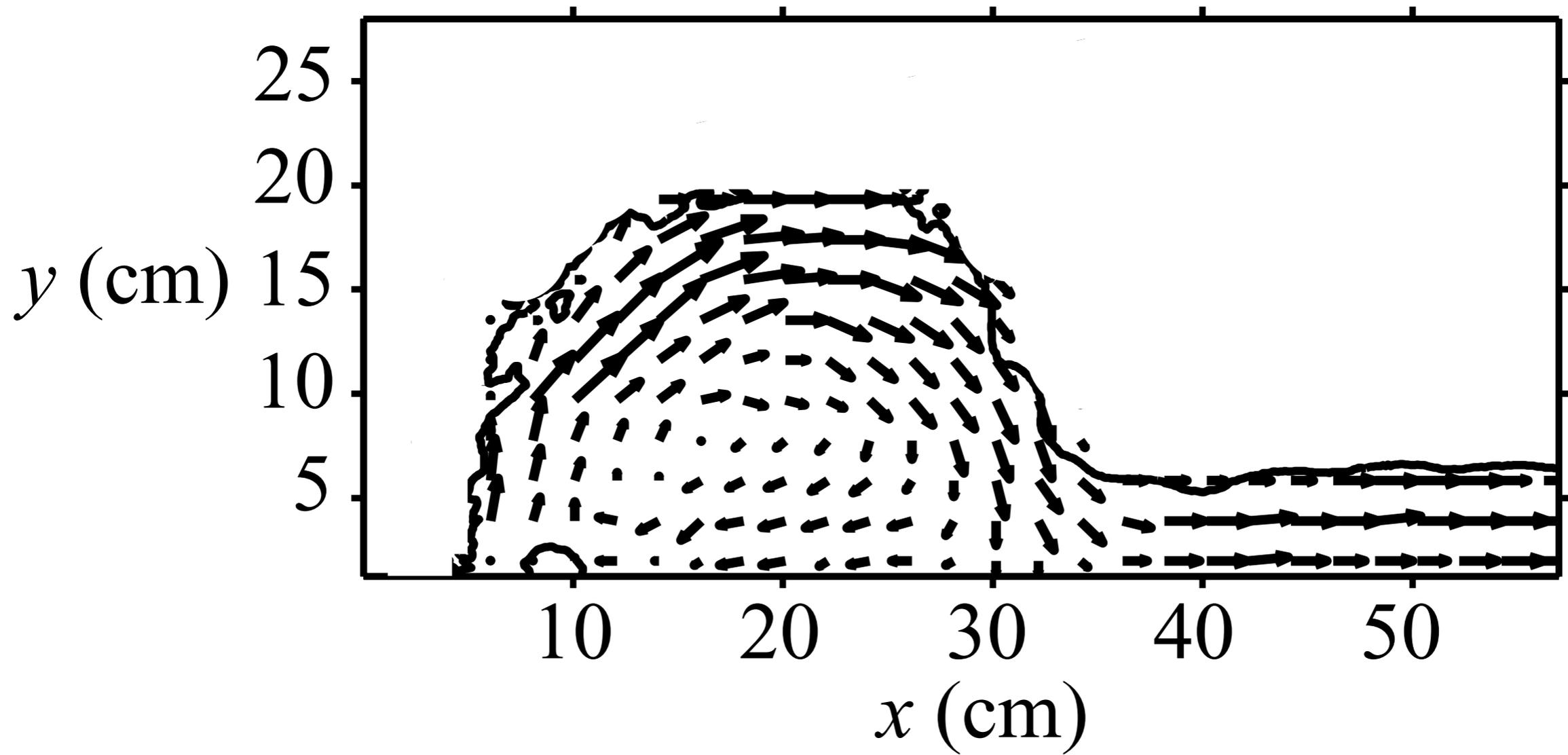


Submesoscale?

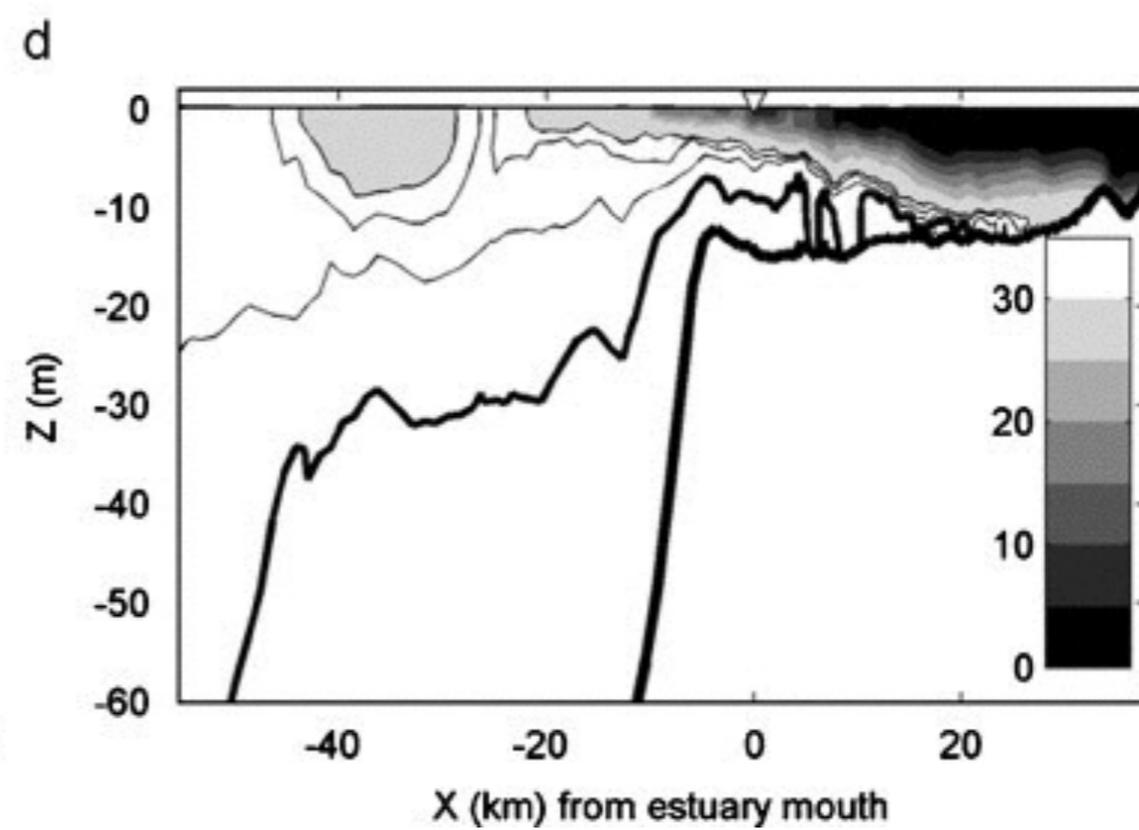
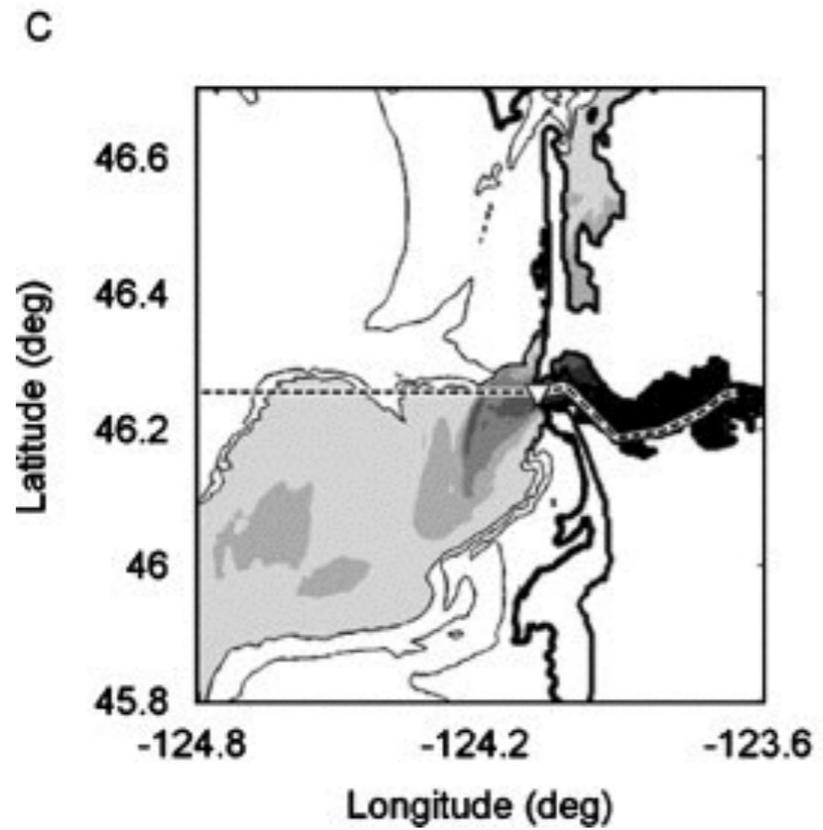
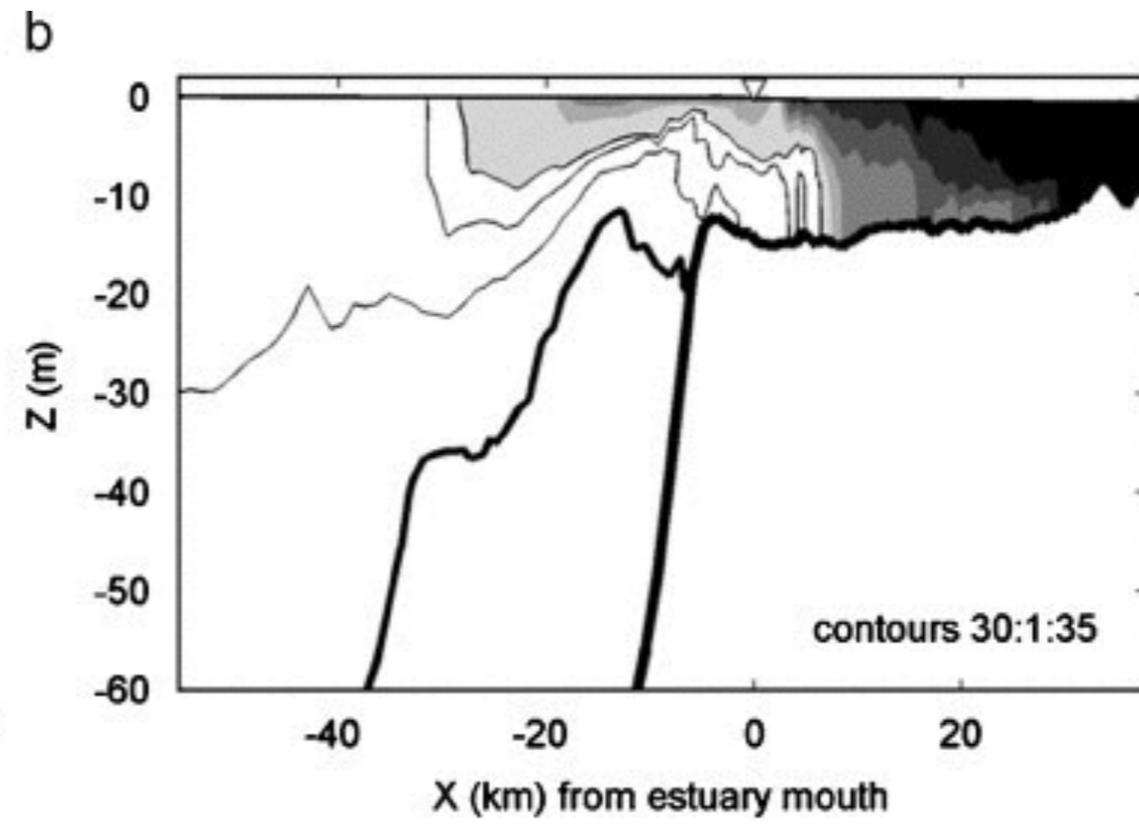
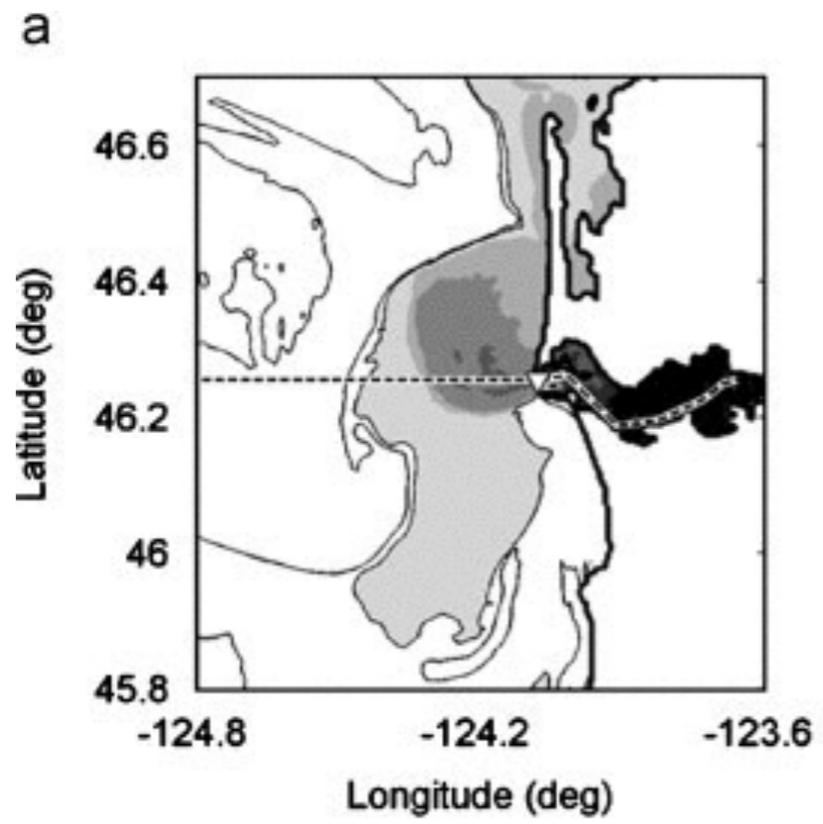
Submesoscale compared to 'mesoscale' Loop Current and Loop Current Eddies

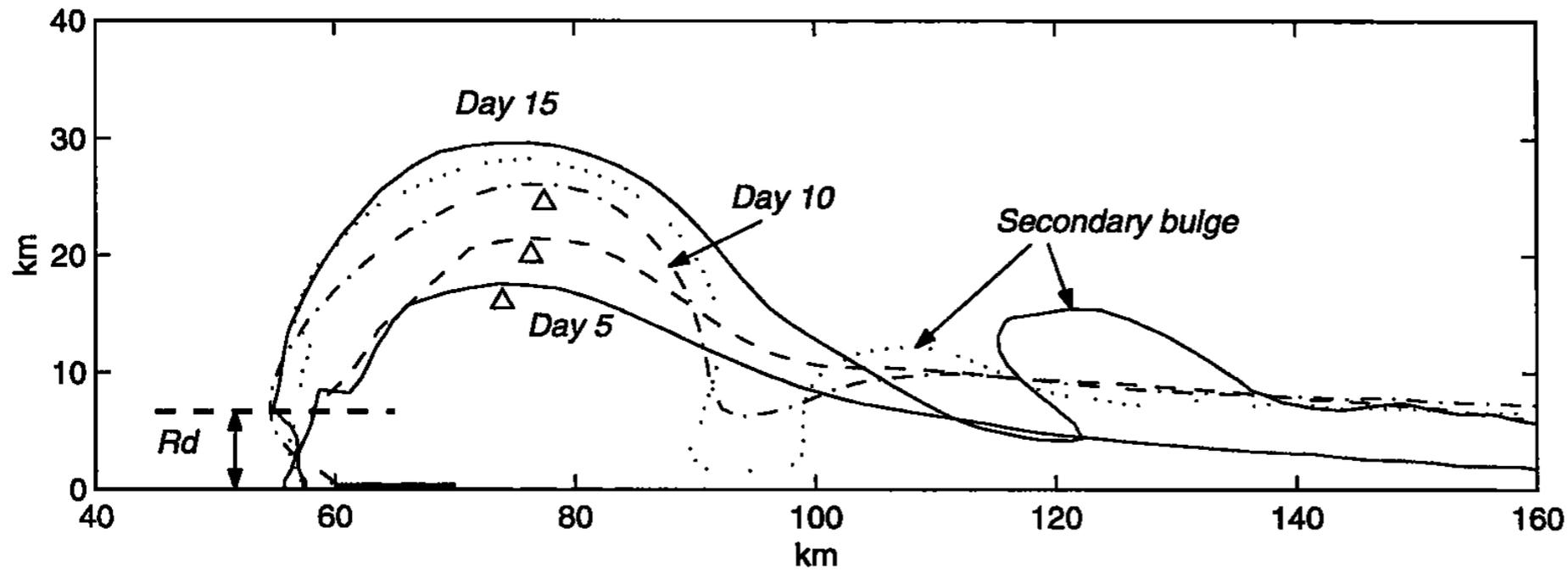
Similar to 'surface mixed layer instabilities':

- Similar relevant parameters ($Ri \sim 1-10$)
- Similar sub-observational scales, $O(10 \text{ km})$



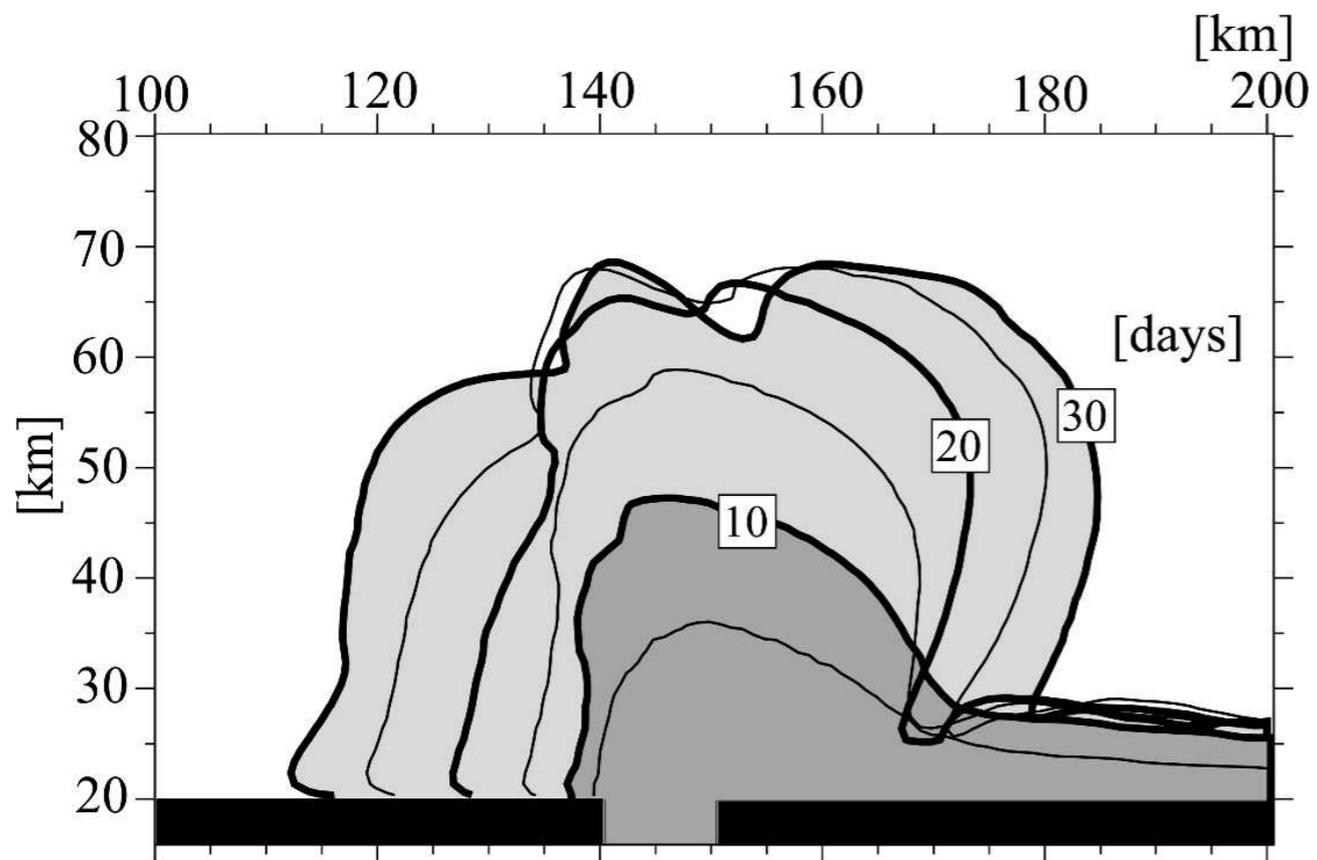
Horner-Devine *et al.* (JFM, 2006)





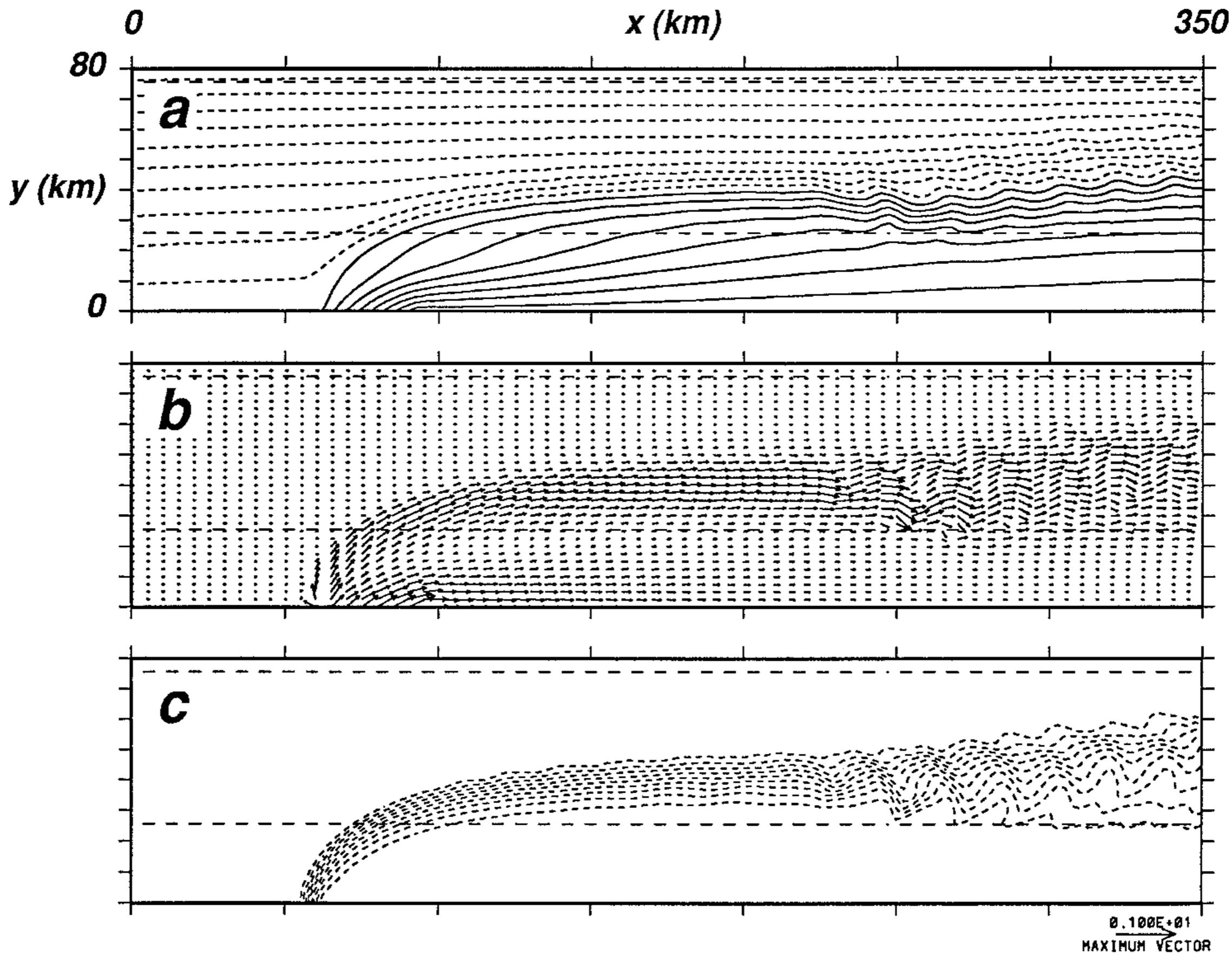
Yankovsky *et al.*
(*JGR*, 2001)

Unsteady discharge

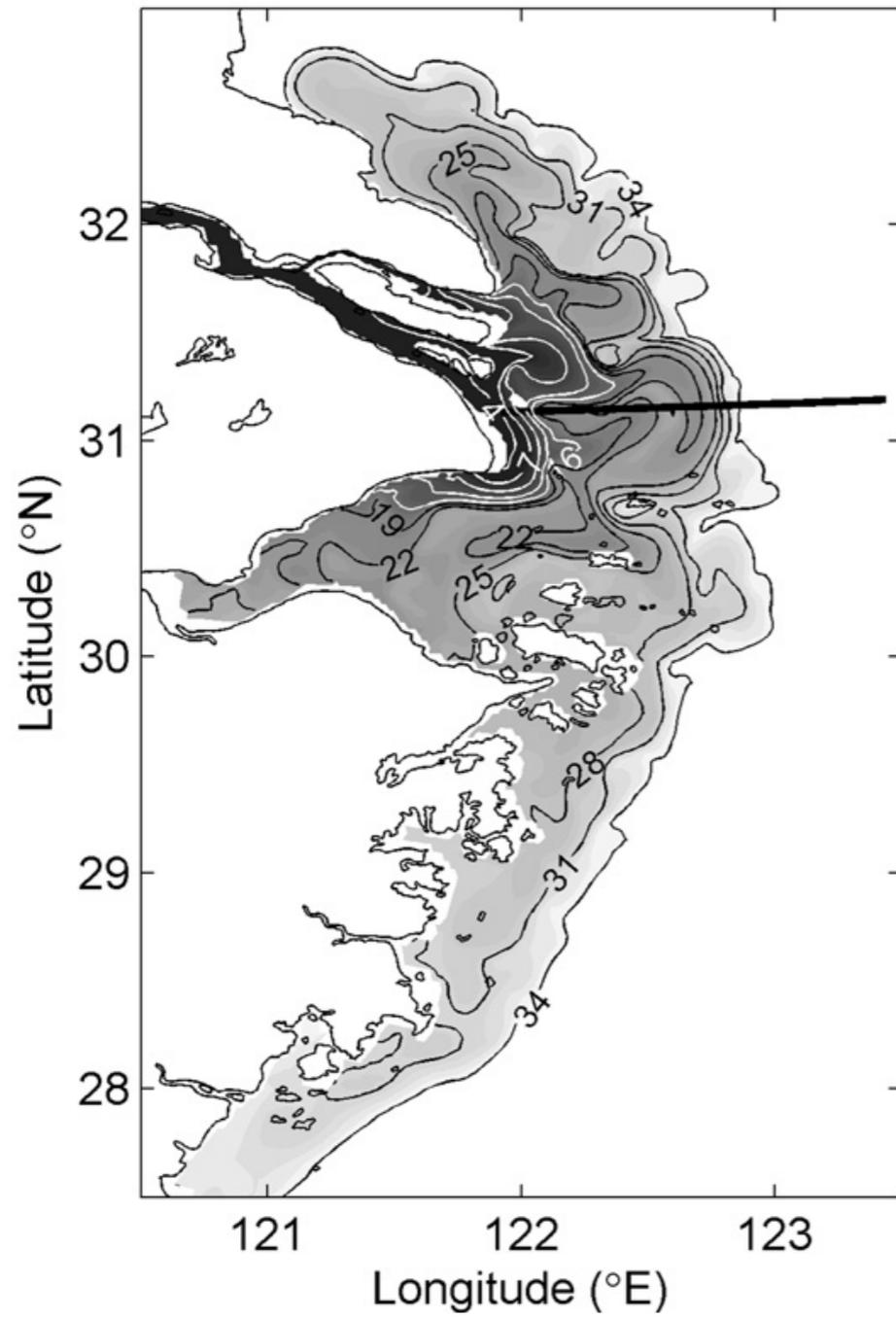


Isobe (*JPO*, 2005)

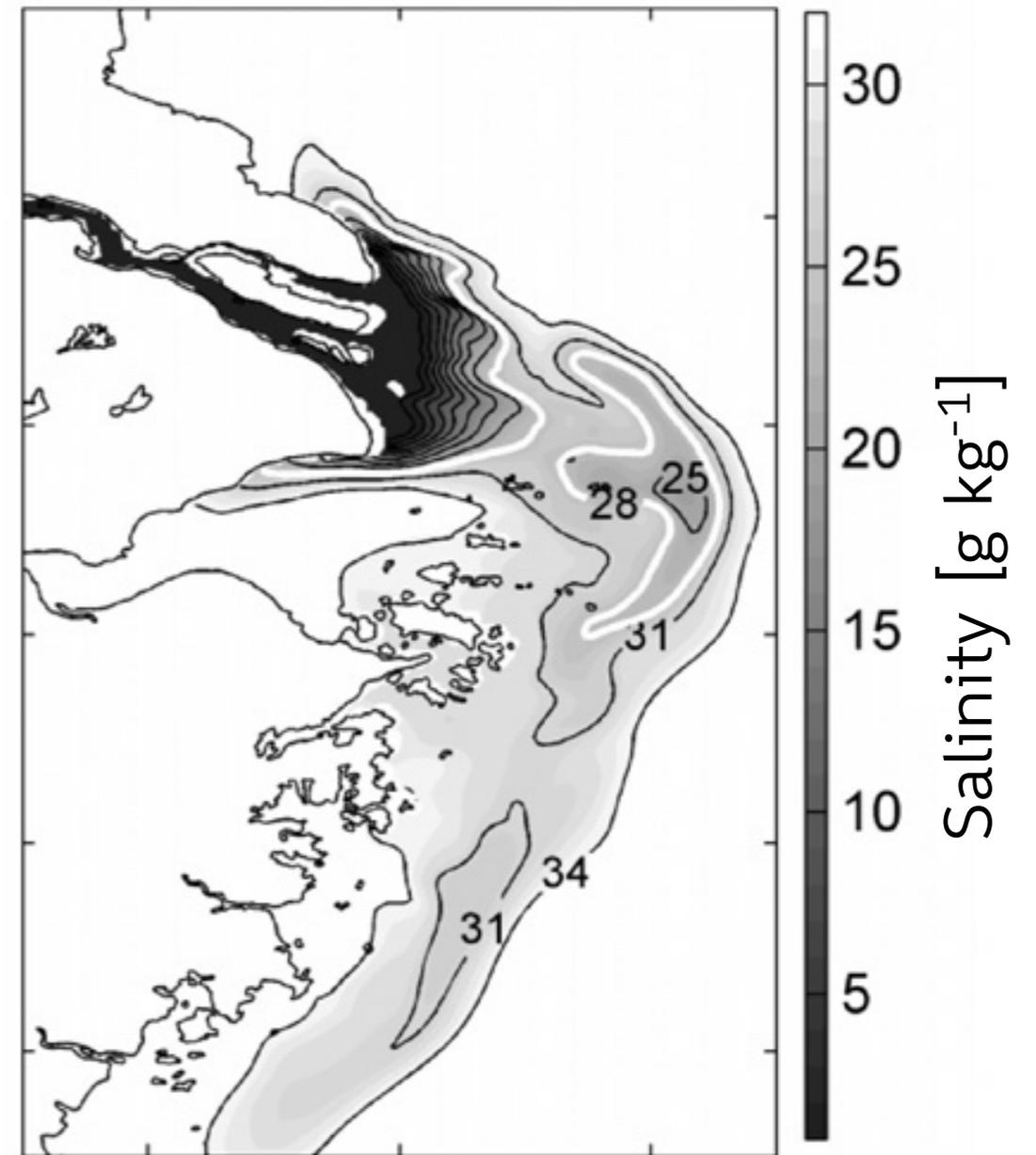
Inertial instability

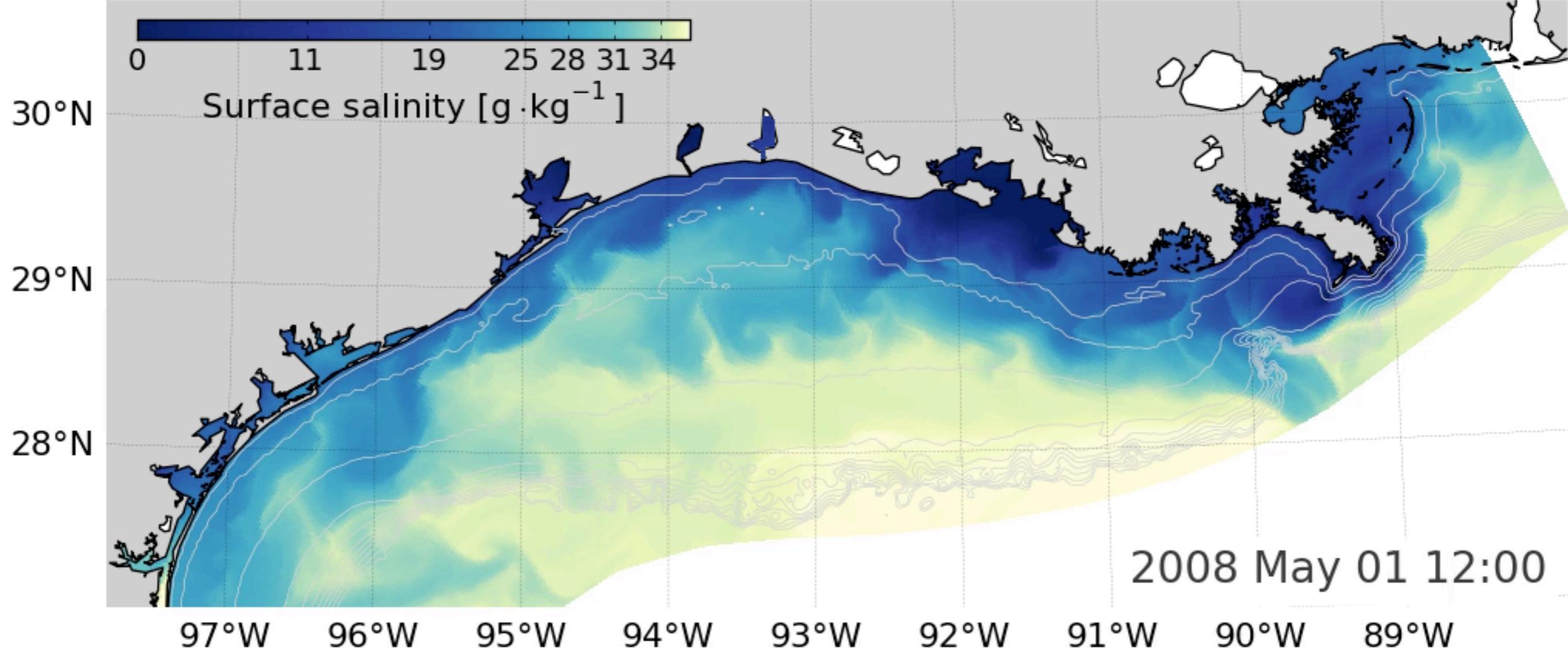


NO TIDE



WITH TIDE





ROMS - TXLA model

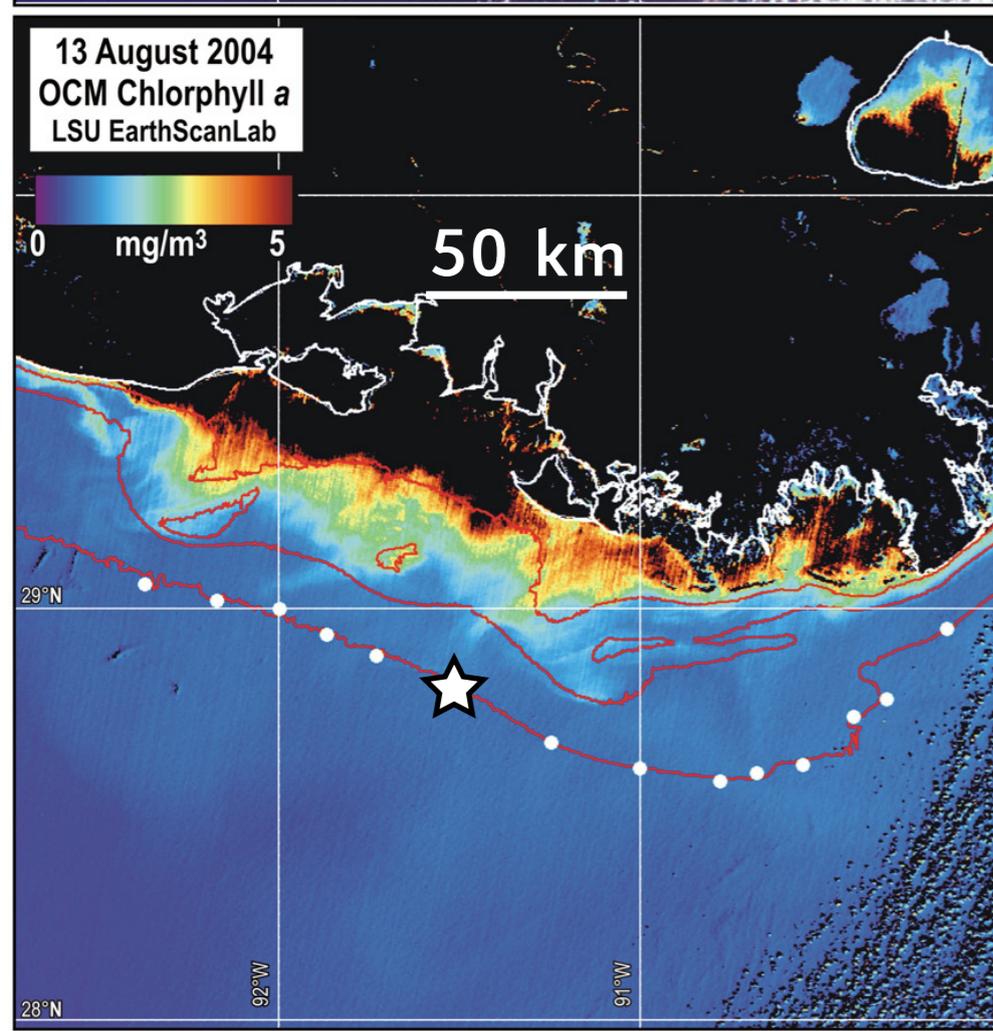
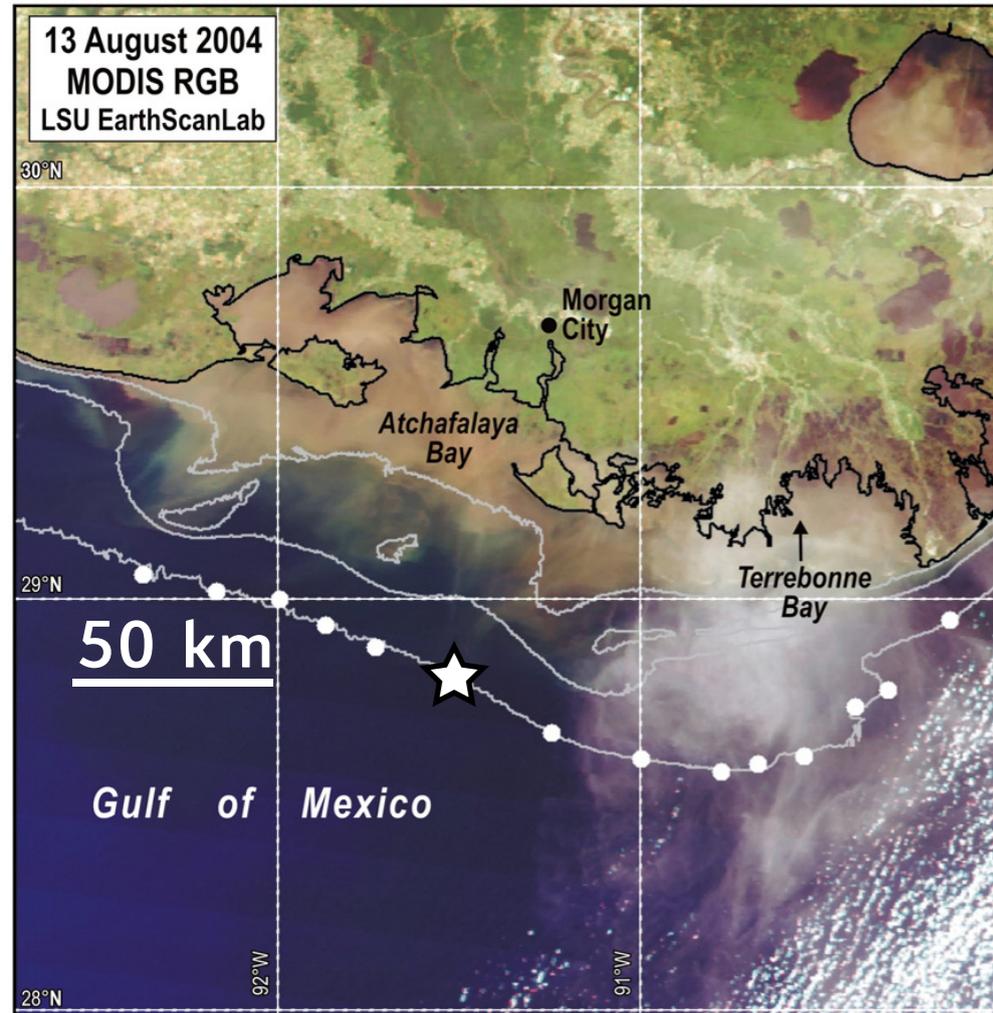
~1 km resolution

Forced with NARR surface fluxes

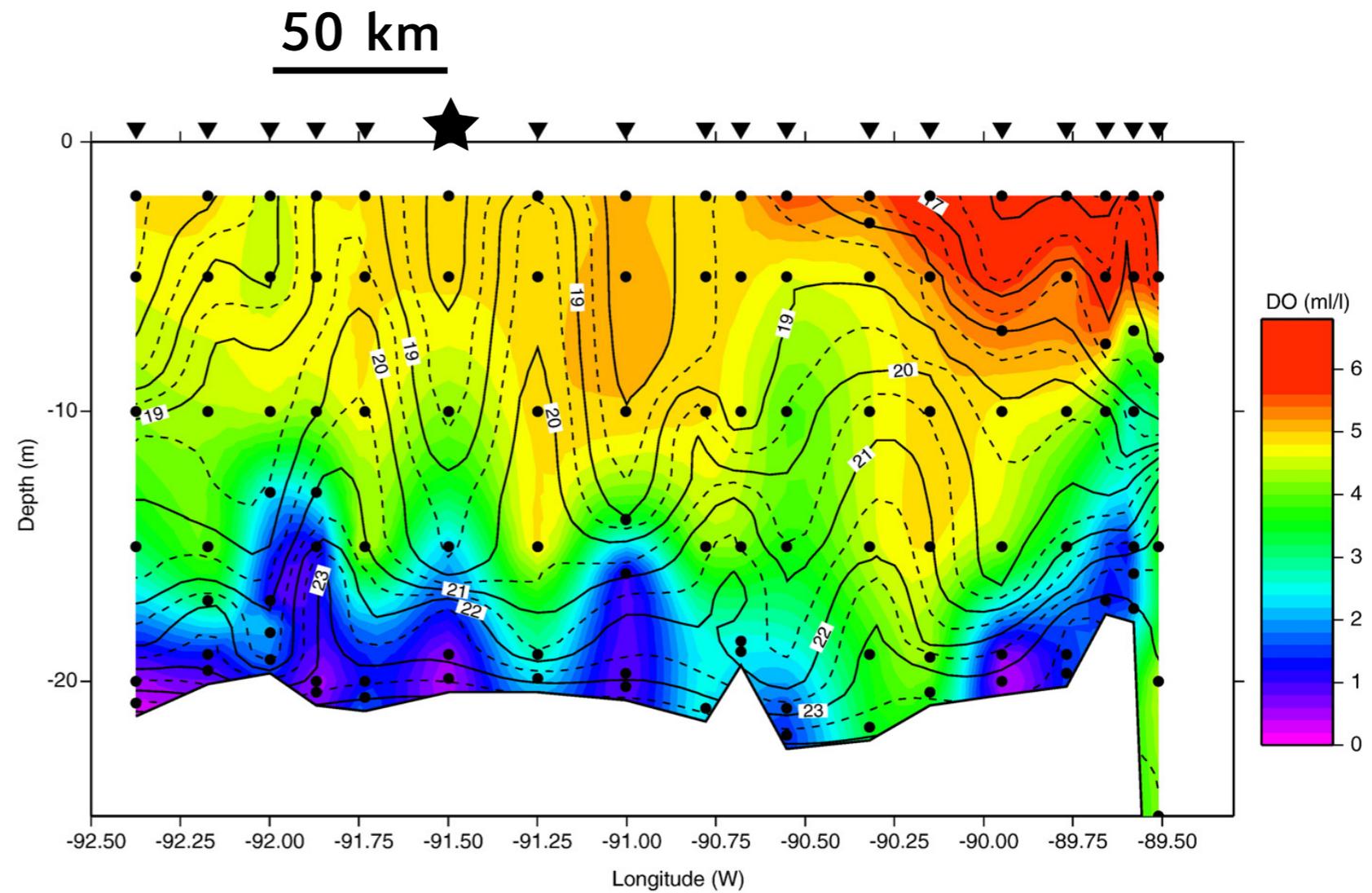
All major regional river included

Nested in GMX HYCOM

X. Zhang *et al.* (*JGR*, 2012)



Topography or instability?



Dimarco et al (2010 JMS)

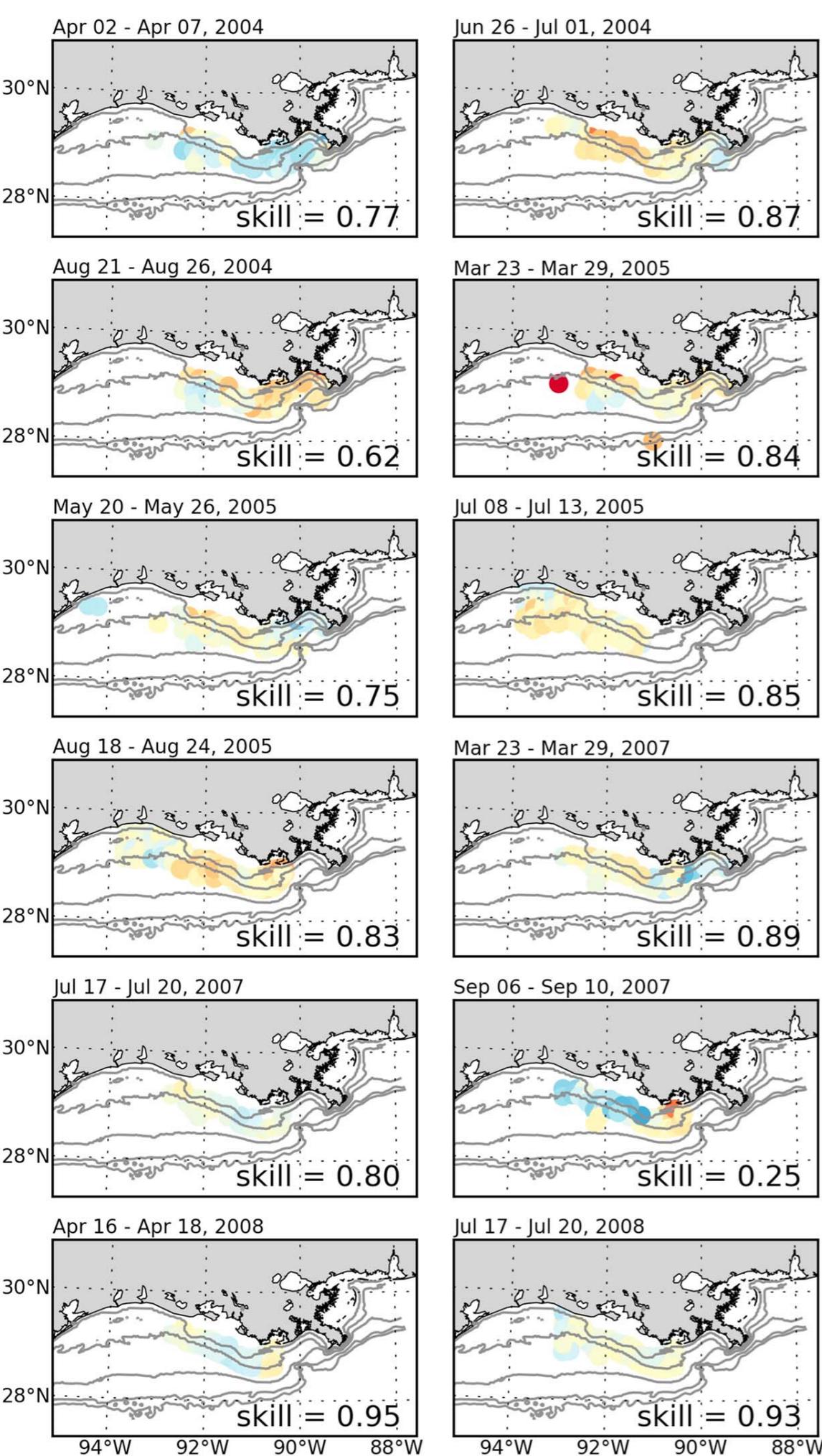


Table 2. Root-Mean-Square (RMS) of the Observed Salinity Relative to Modeled Salinity and Climatological Salinity, Respectively (Equations (2) and (3)) for Each Hypoxia Cruise During Years 2004 to 2008

Hypoxia Cruise	RMS (Obs – Mod)	RMS (Obs – Clim)
2004 MCH01	1.61	3.33
2004 MCH02	1.24	3.38
2004 MCH03	1.52	2.47
2005 MCH04	1.52	3.79
2005 MCH05	1.69	3.42
2005 MCH06	1.29	3.35
2005 MCH07	1.40	3.42
2007 MCH08	1.44	4.38
2007 MCH09	1.56	3.51
2007 MCH10	1.65	1.90
2008 MCH11	1.39	6.13
2008 MCH12	1.36	5.00

Climatology WOA09
1°x1° resolution

$$\text{skill} = \frac{\sum(\text{obs}_i - \text{mod}_i)^2}{\sum(\text{obs}_i - \text{clm}_i)^2}$$

X. Zhang *et al.* (JGR, 2012)

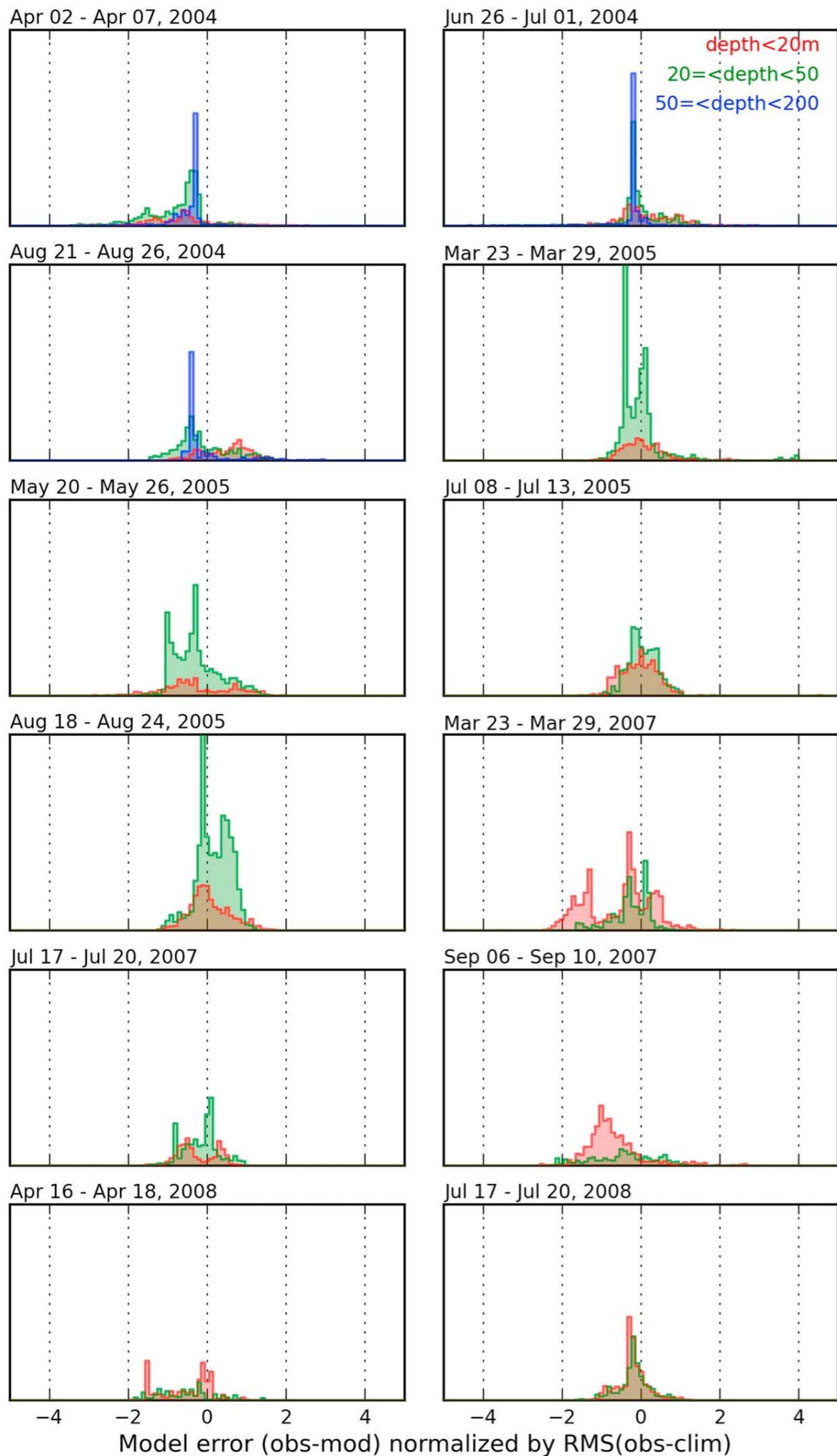
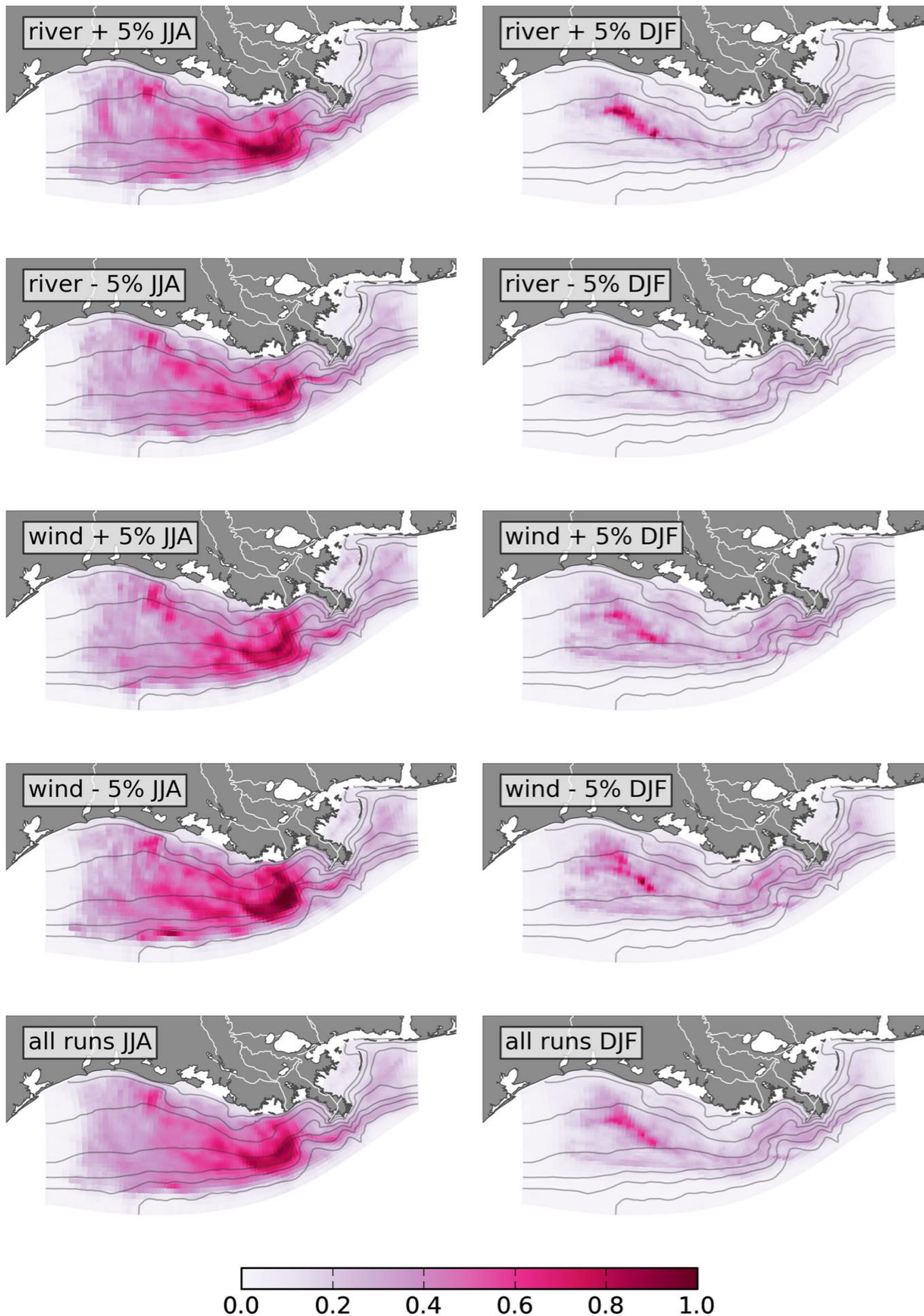


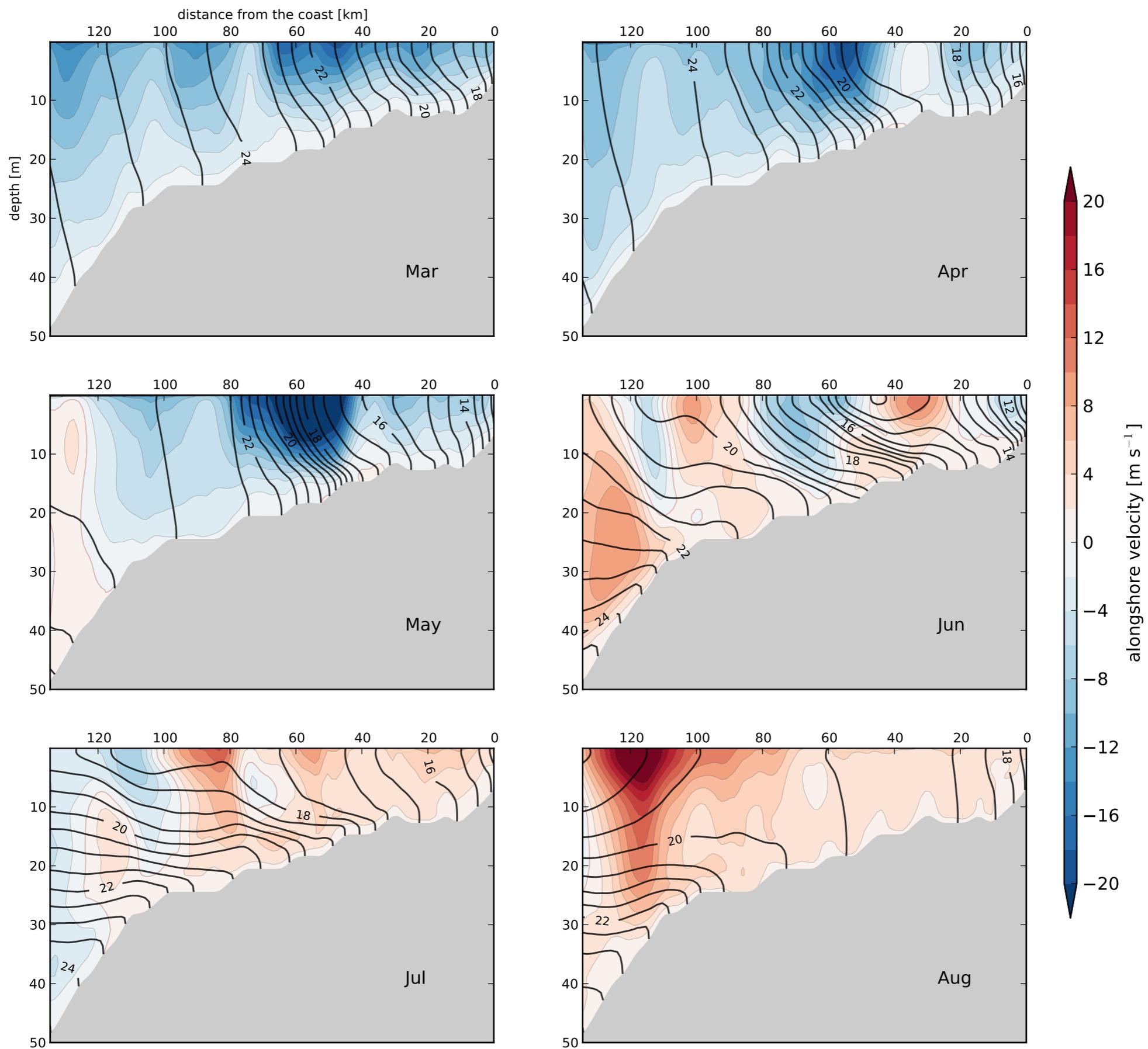
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2008 MCH12	1.36	5.00



Simulations suggest:

- eddy field is chaotic,
- strongest in summer, and
- may be triggered by a variety of processes.



Base state

Mean state:

$$fU = -\frac{\rho_y}{\rho_0}$$

$$-b = -\frac{\rho_z}{\rho_0}$$

so that,

$$U_z = -\frac{b_y}{f} = \frac{M^2}{f}$$

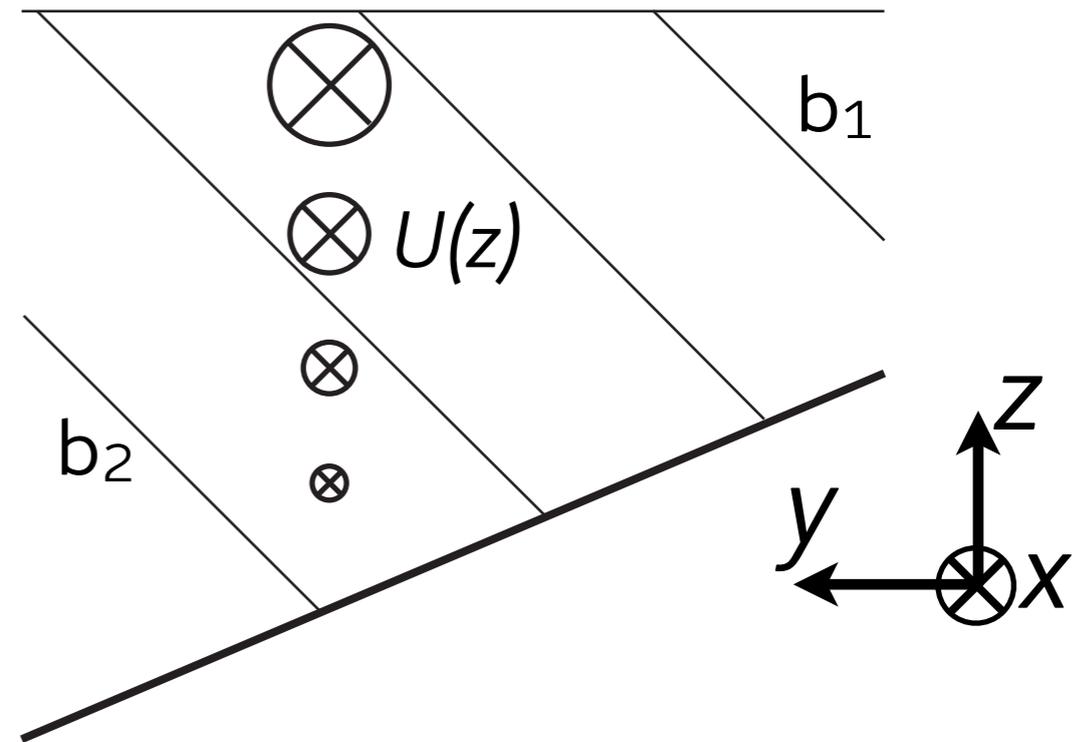
where:

$$N^2 = -b_z$$

$$M^2 = -b_y$$

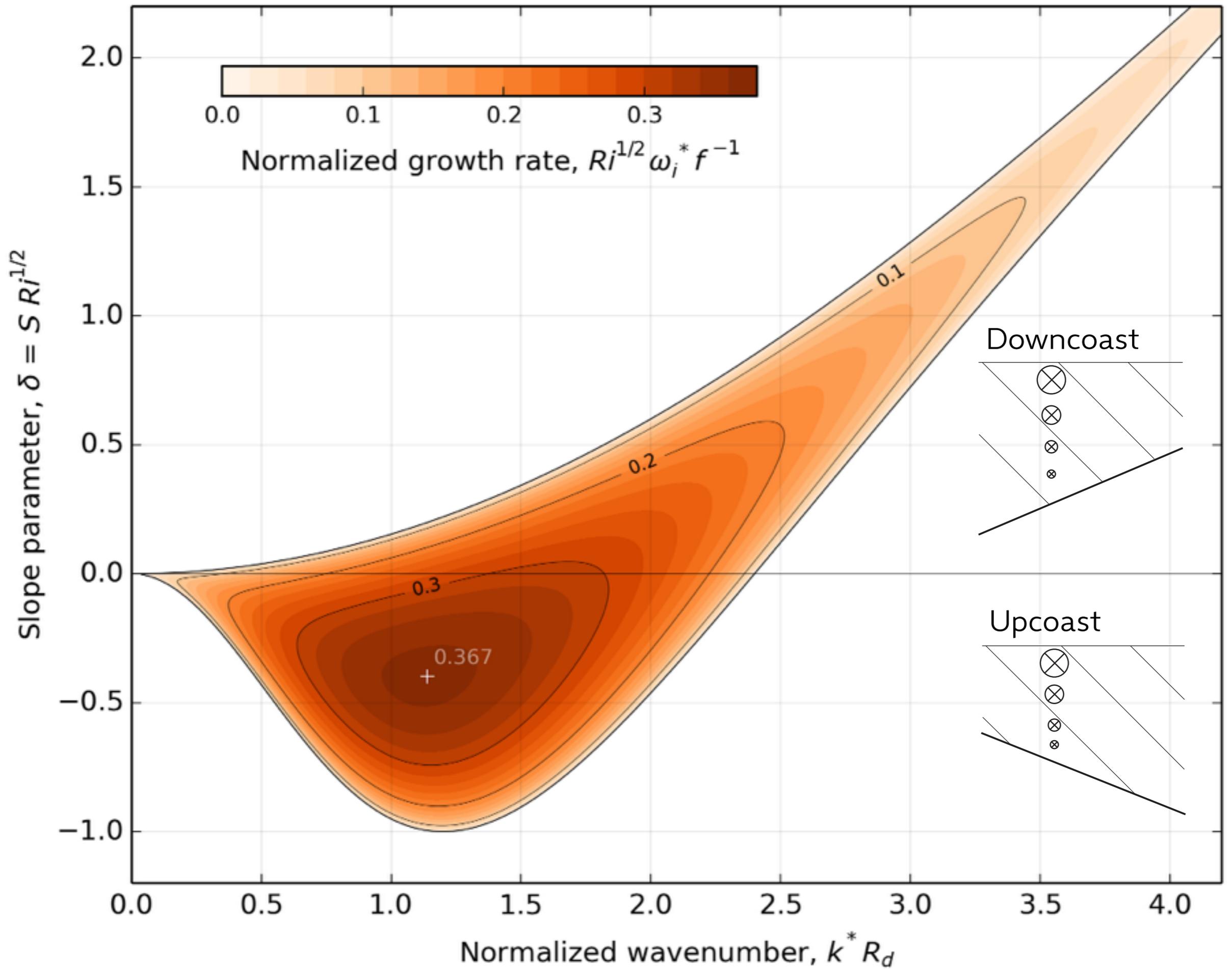
Bouyancy:

$$b = g \left(\frac{\rho_0 - \rho}{\rho_0} \right)$$

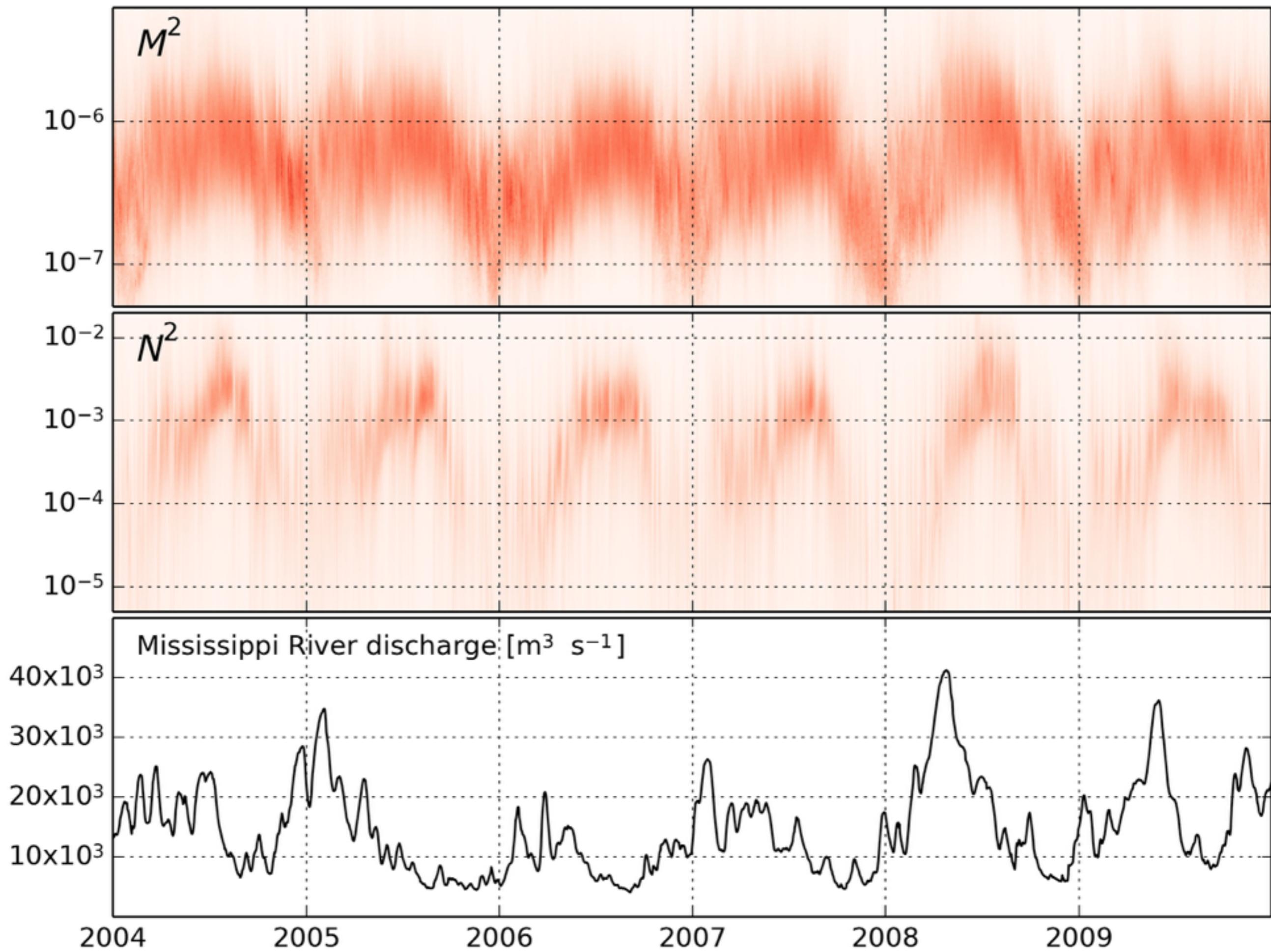


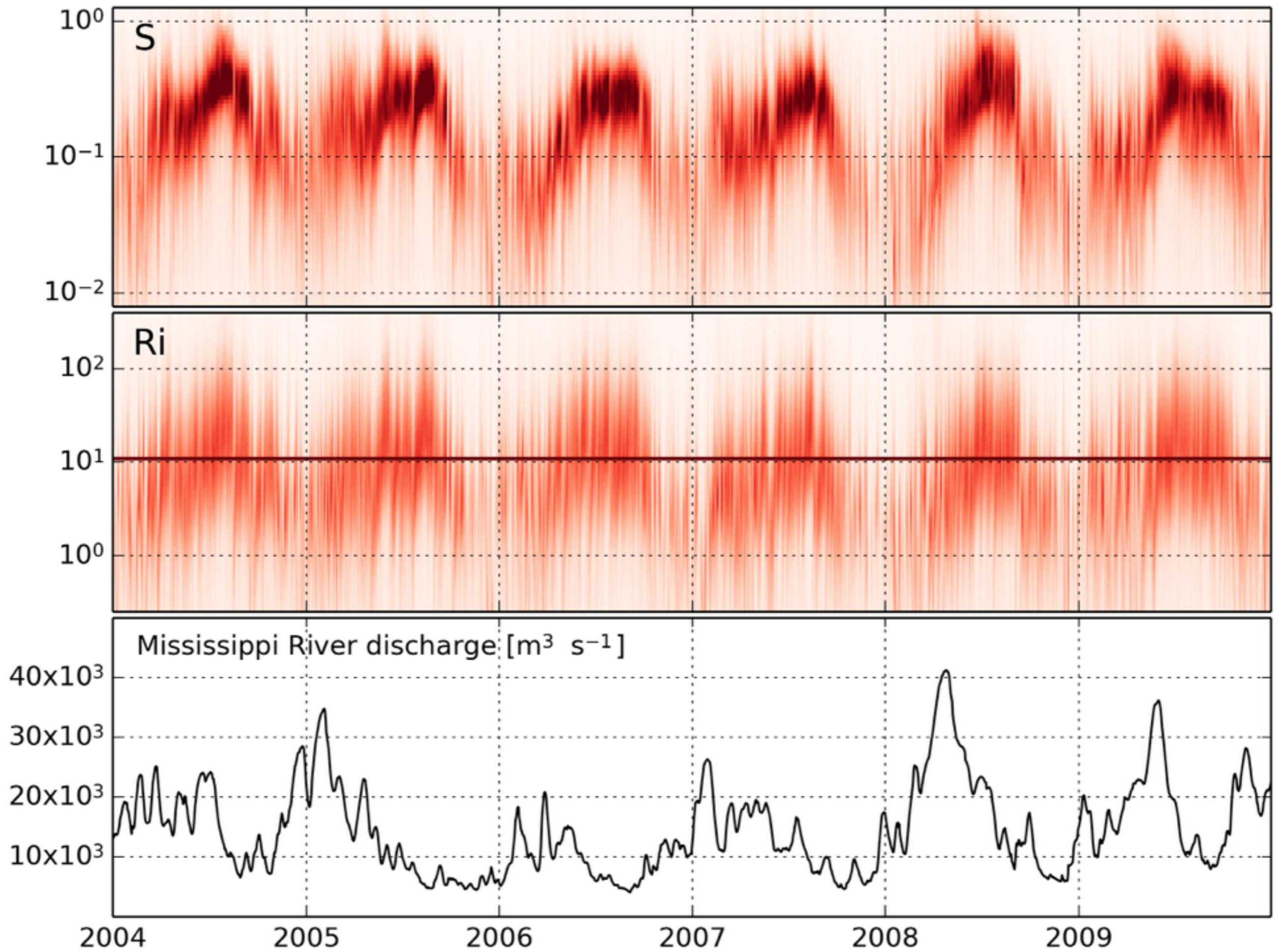
Relevant parameters:

$$Ri = \frac{N^2 f^2}{M^4} \quad S = \frac{N}{f} \alpha$$



Adapted from Blumsack and Geirasch (*JAS*, 1972)



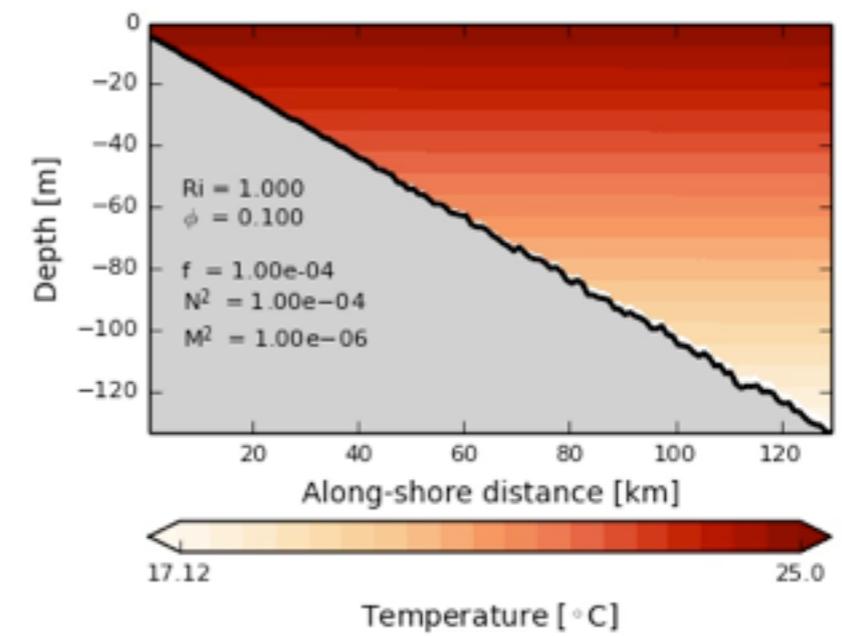
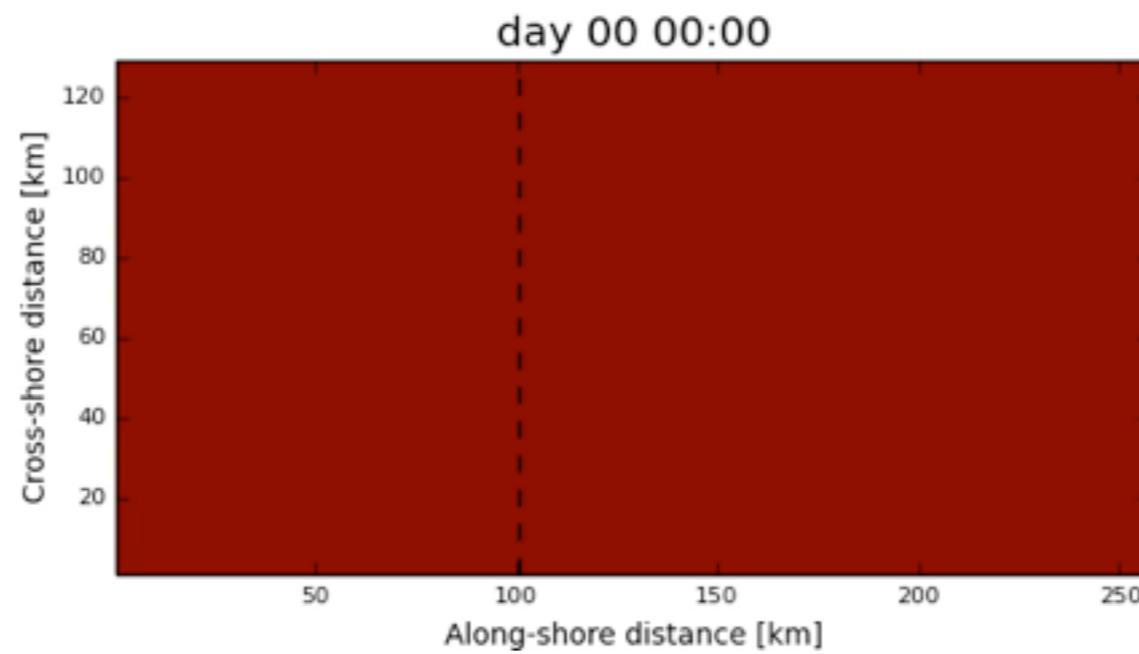


Base case:

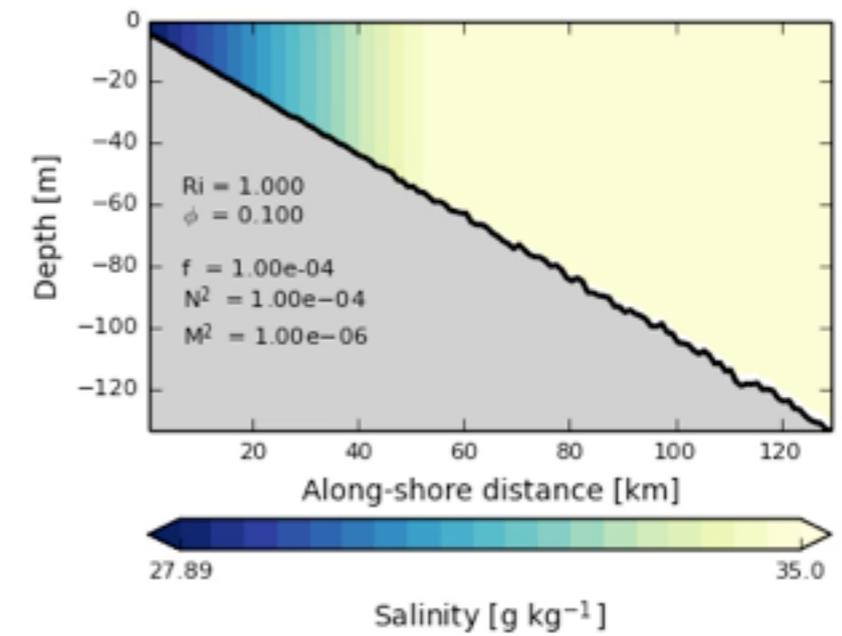
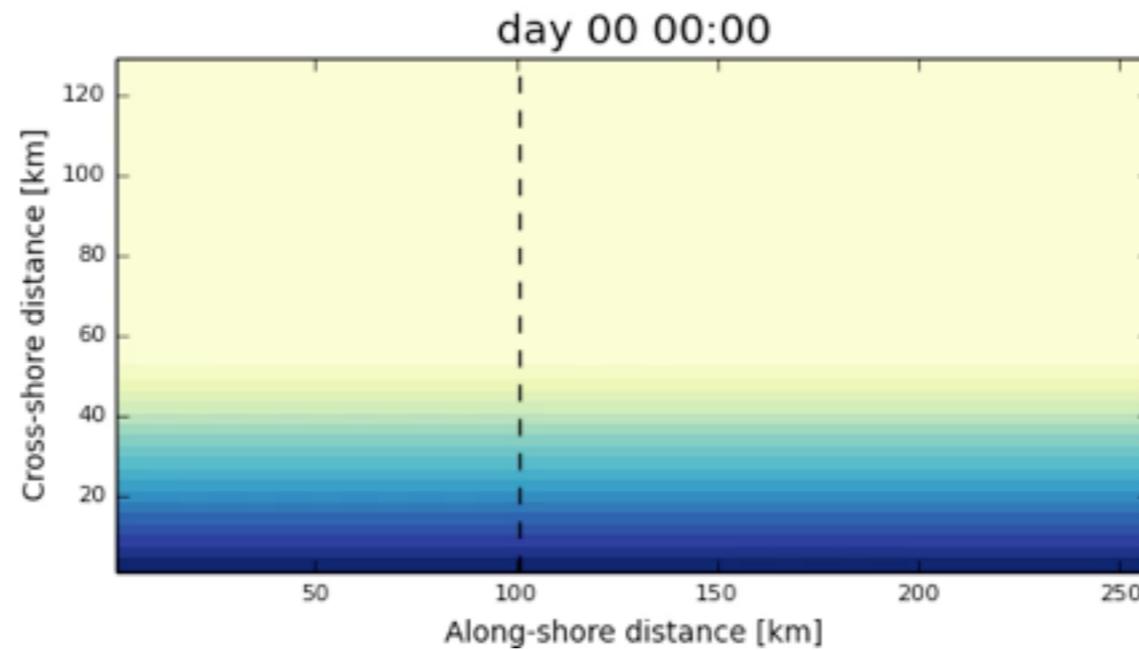
$$Ri = 1.0$$

$$S = 0.1$$

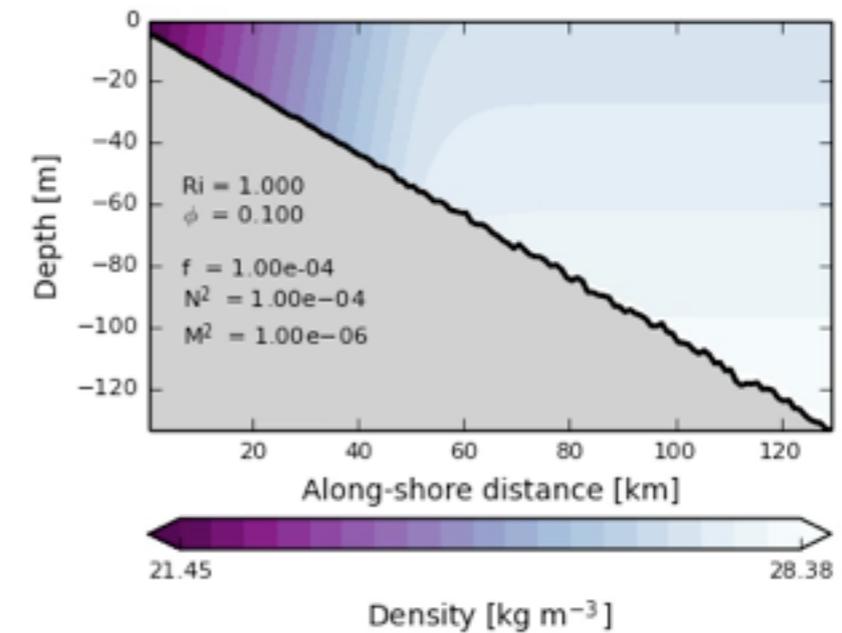
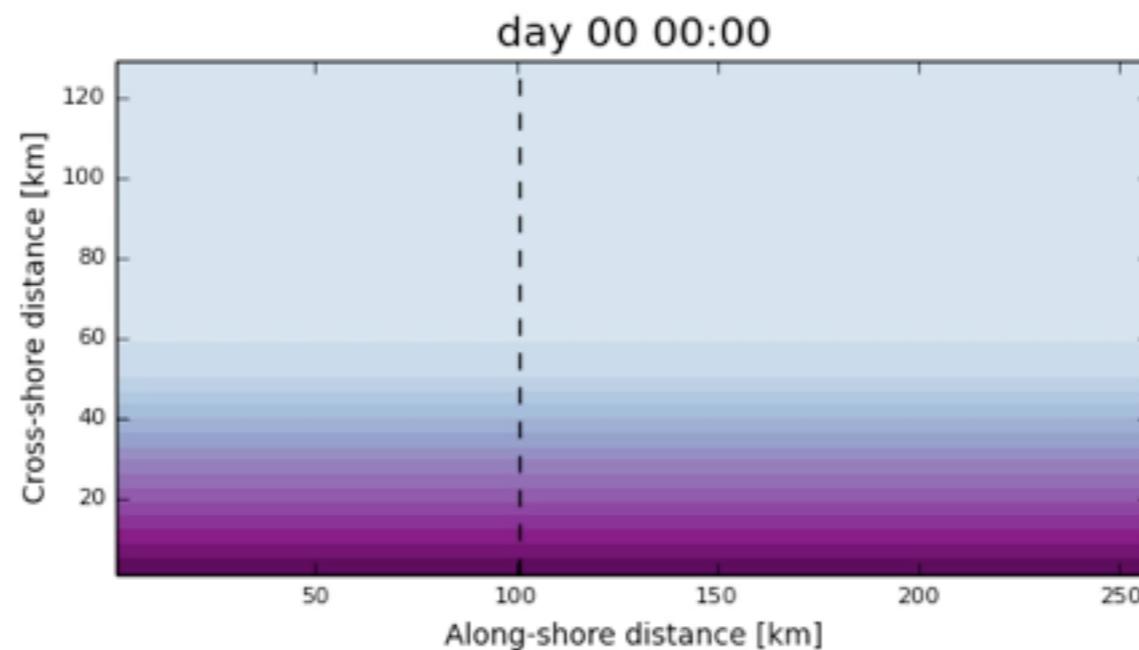
Temperature



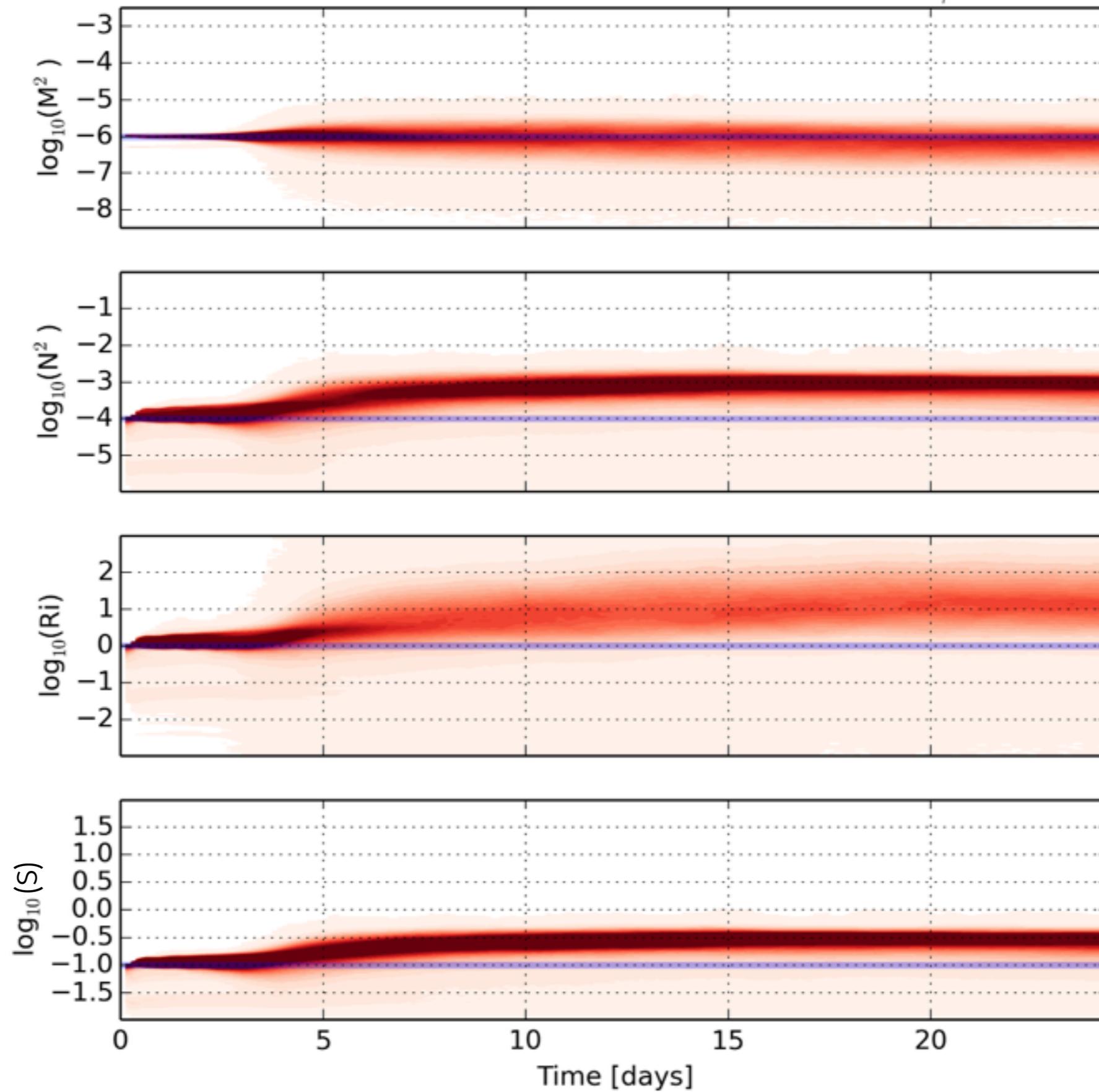
Salinity



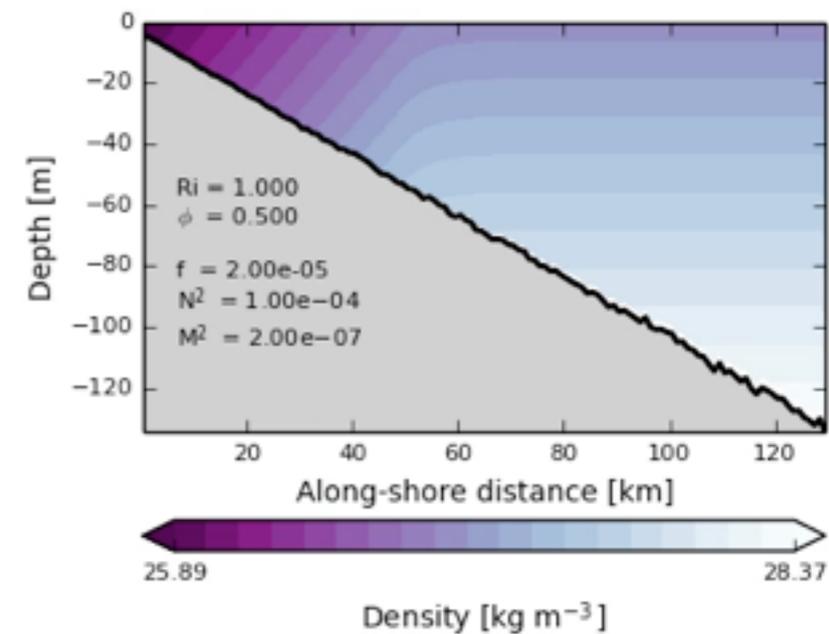
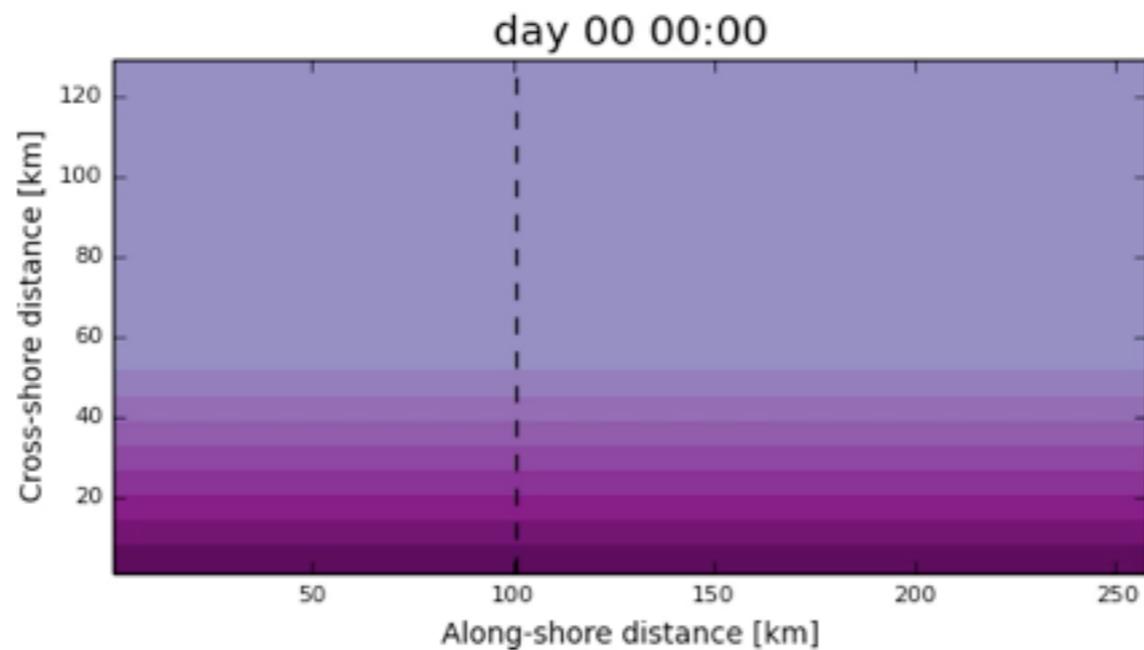
Density



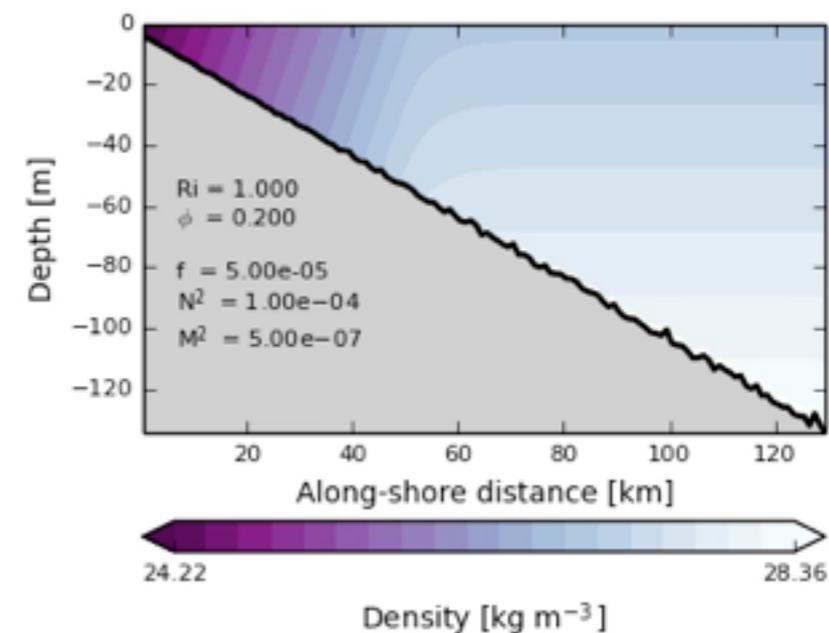
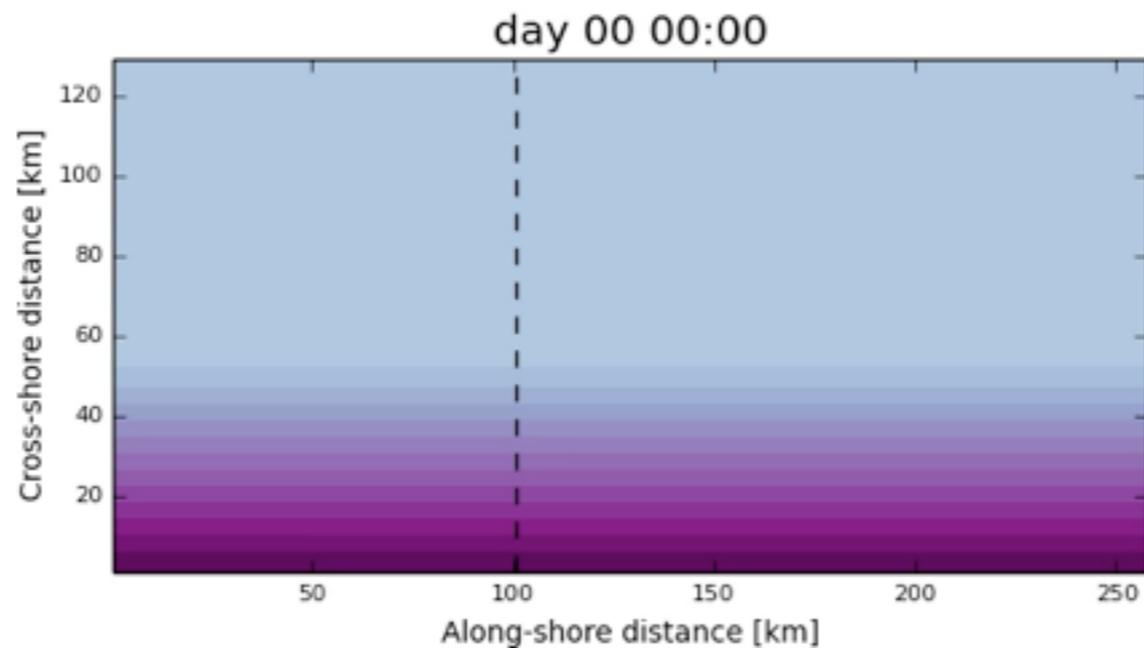
$M^2 = 1.00e-06$ $N^2 = 1.00e-04$ $Ri = 1.00e+00$ $\phi = 1.00e-01$



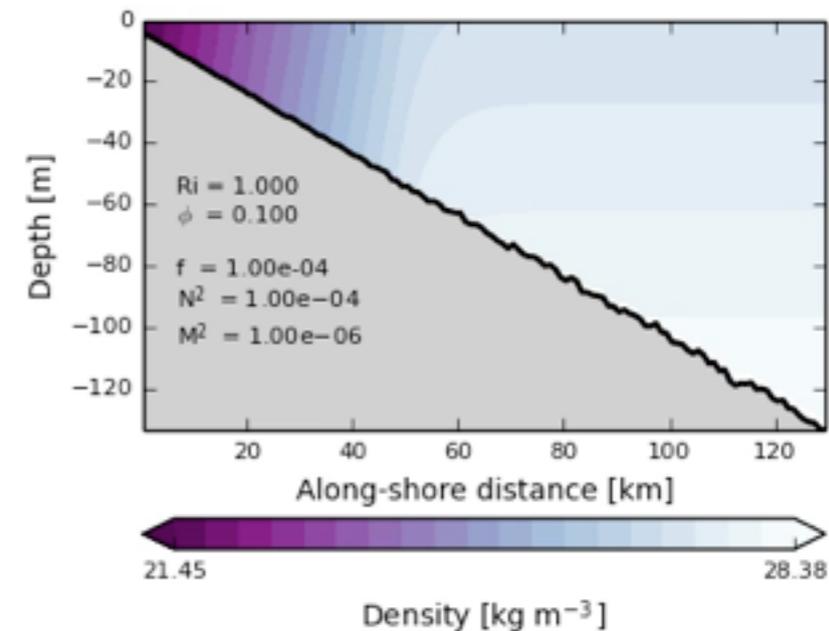
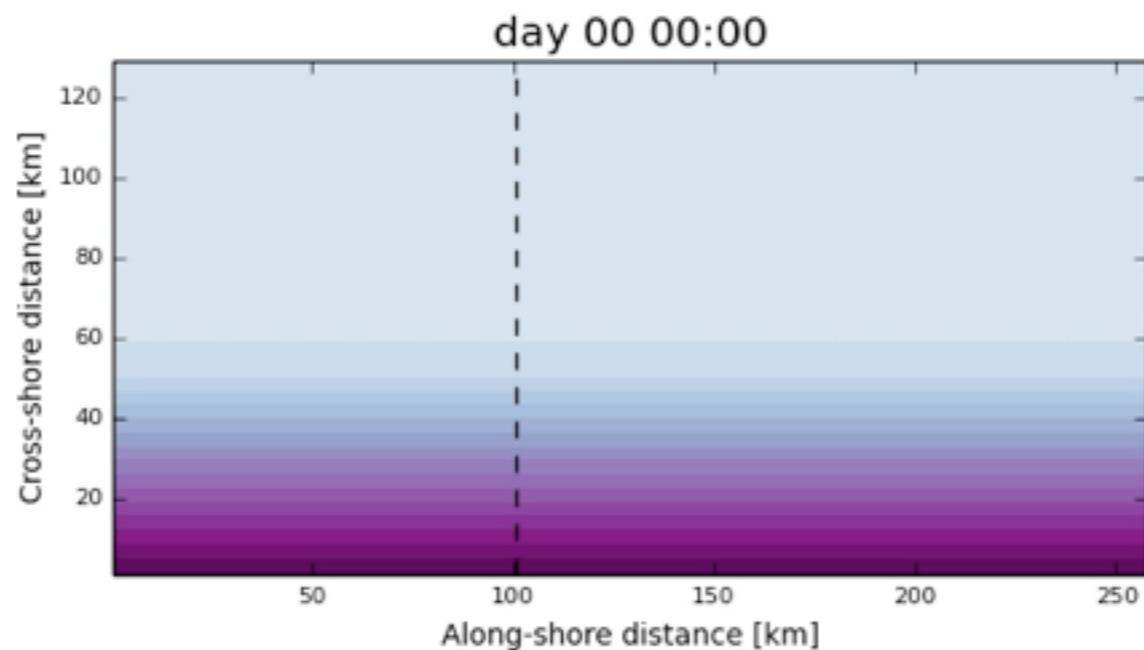
$S = 0.5$

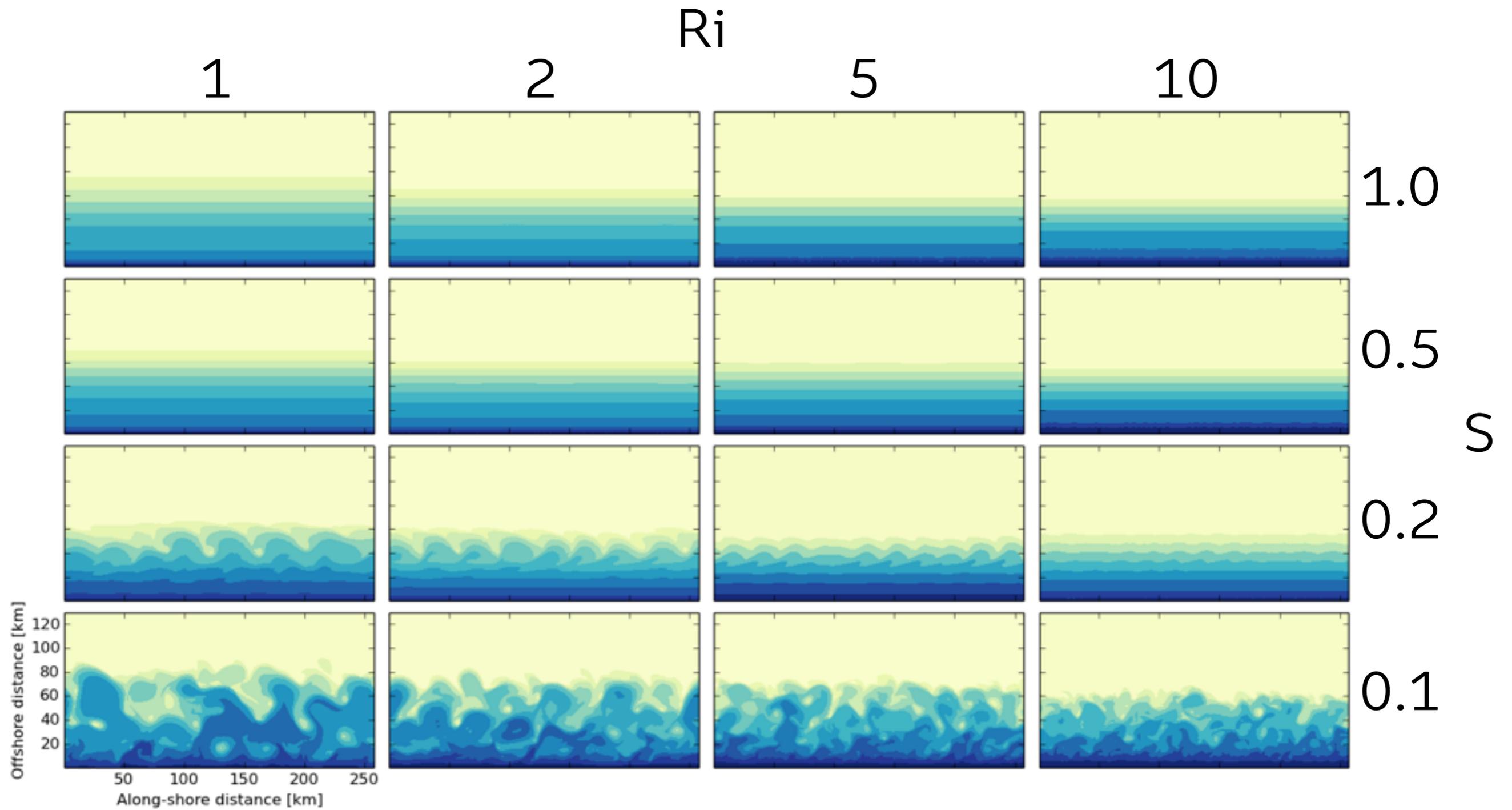


$S = 0.2$



$S = 0.1$

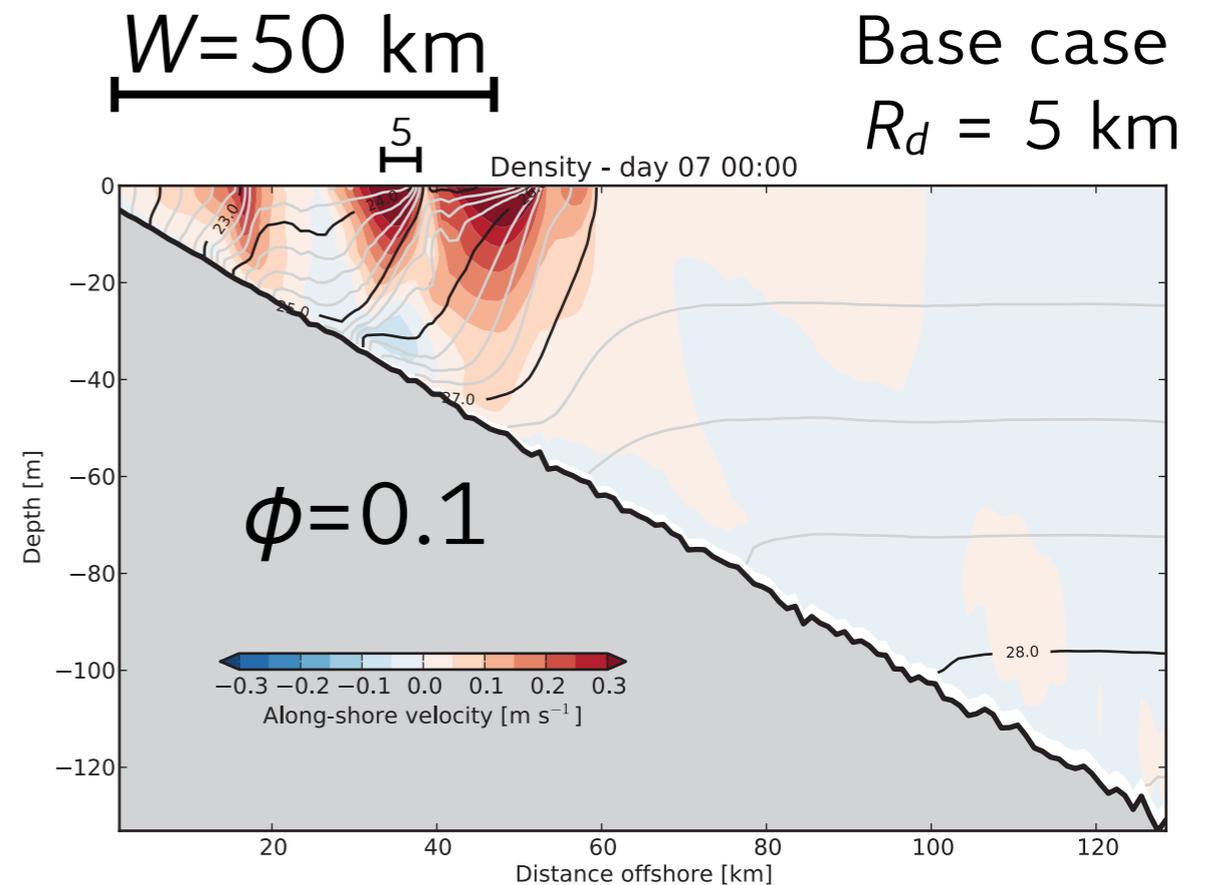




Eddies only form for $S \leq 0.3$, or $3Rd < W$.

In order to form,
eddies need to fit.

$$W > R_d = \frac{NH}{f}$$

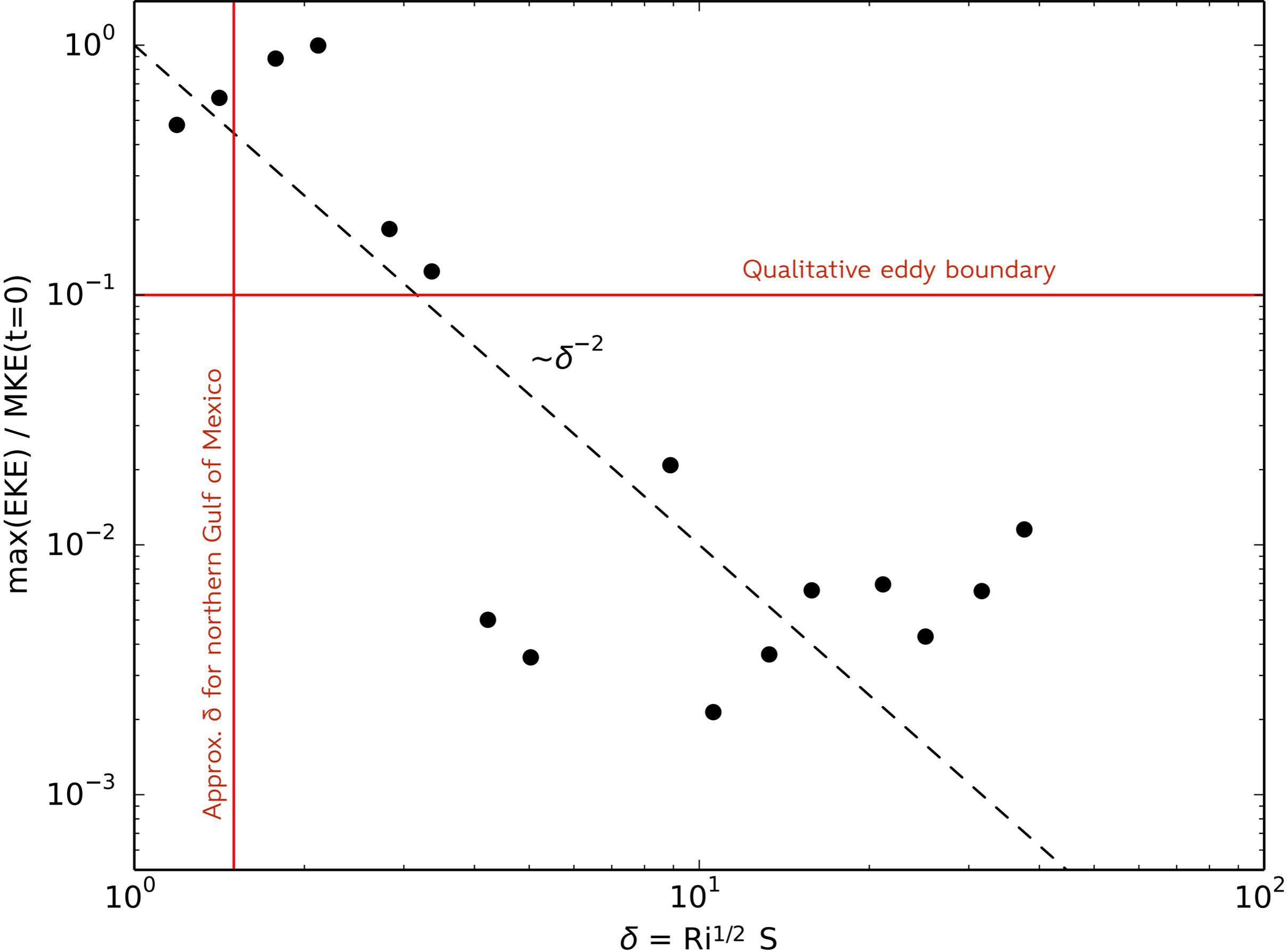


Where the aspect ratio, α , is used to define

$$1 > \frac{NH}{fW} = \frac{N}{f}\alpha \quad \text{and using} \quad Ri = \frac{N^2 f^2}{M^4}$$

$$1 > S \equiv \frac{R_d}{W} = \frac{N}{f}\alpha = \frac{M^2}{f^2} \alpha \sqrt{Ri}$$

Maximum EKE normalized by initial mean KE



Conclusions:

Eddies need to fit – $S < 0.3$.

This may be why baroclinic instabilities are rare in river plumes.

When eddies form, the field evolves to $S \sim 0.3$.

This suggests that eddies cannot export fresh water from the shelf alone.