

Fine-scale turbulent processes: mesoscale stirring and submesoscale instabilities

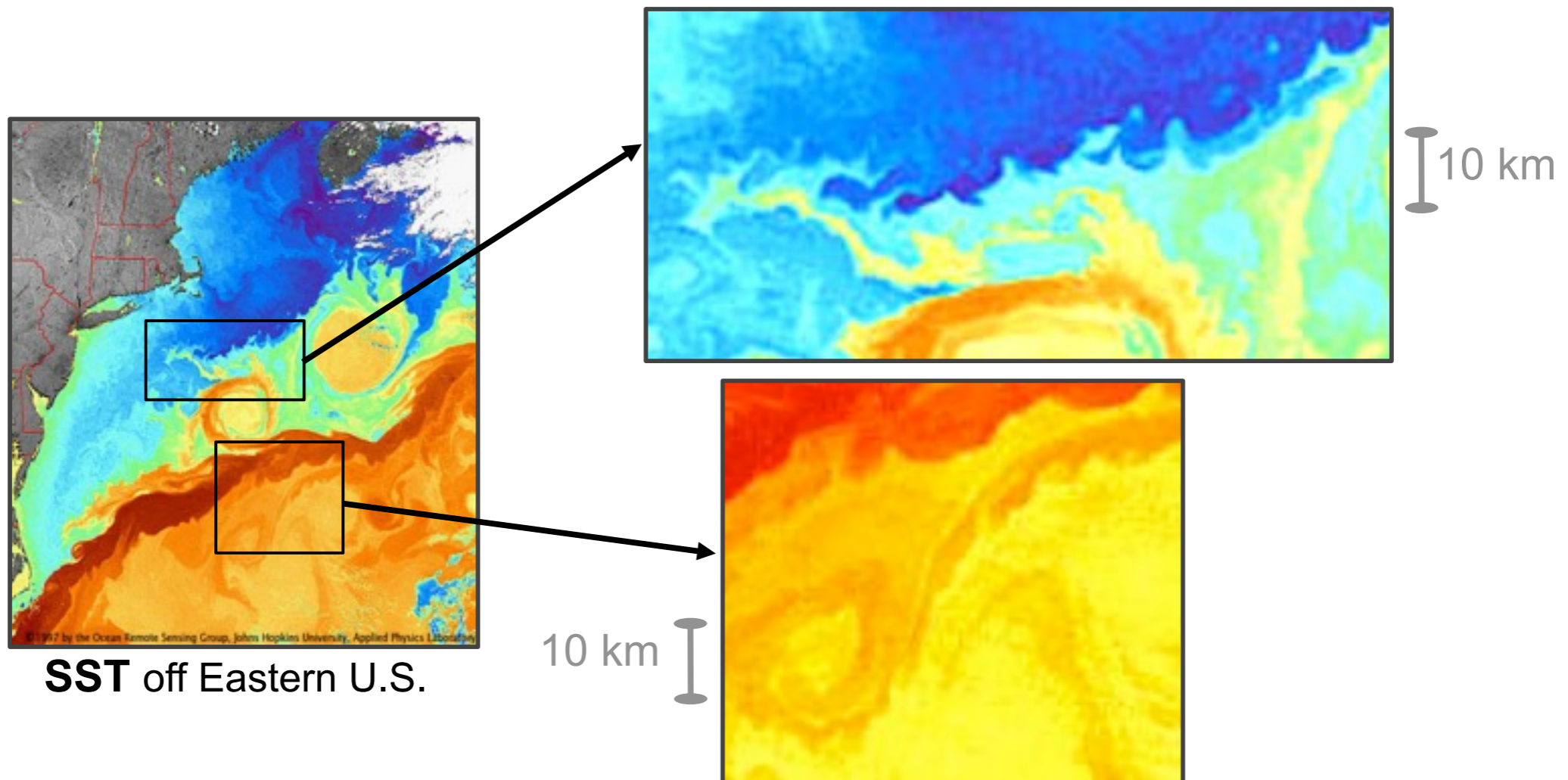
X. Capet (LOCEAN)

G. Roullet, P. Klein, A. Ponte, G. Maze (LPO)

ROMS Workshop, Oct. 2012, Rio de Janeiro

“**submesoscale turbulence**” shorthand for
“**upper-ocean frontal submesoscale turbulence**”

The upper ocean is filled with rapidly evolving structures that appear as tracer contrasts, *i.e.*, fronts. The cross-front length scale is of the order of 1-10 km.



Why is there an upper ocean submesoscale turbulence regime and how does it differ from region to region ?

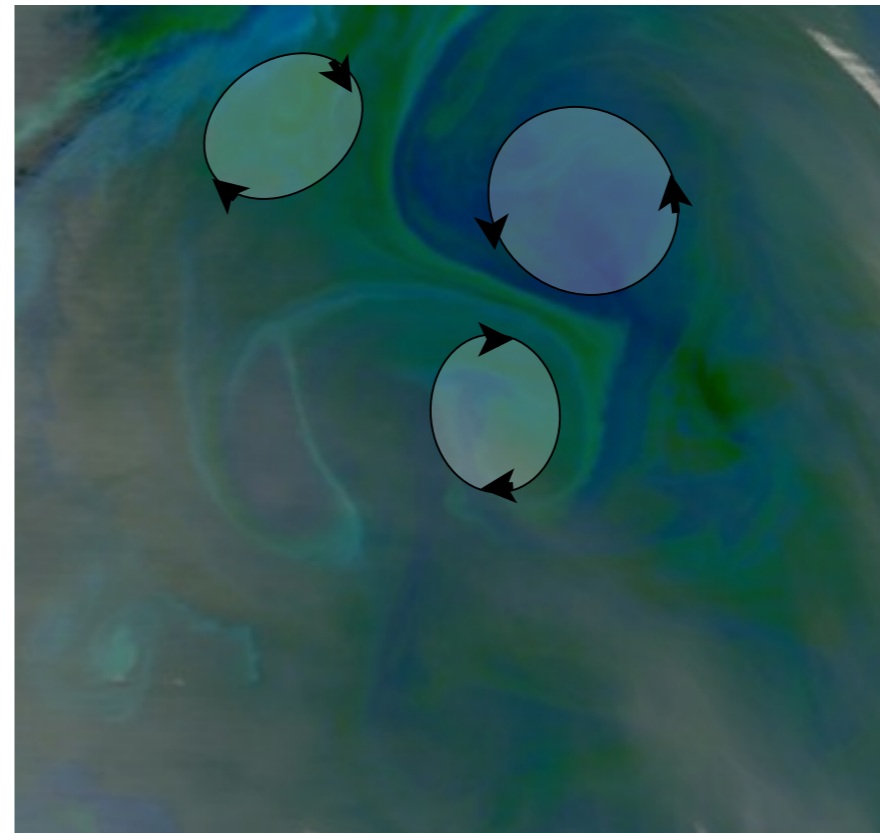
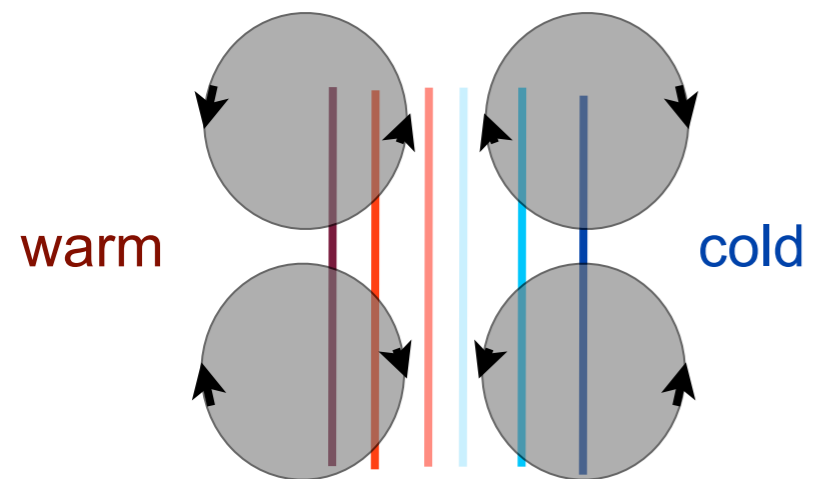
Most of our ROMS simulations resolve (at least partly) submesoscale turbulence nowadays. High order numerical schemes help in that regard (P. Marchesiello).

How do we interpret the type of turbulence that emerges at $dx \sim 1$ km or more ?

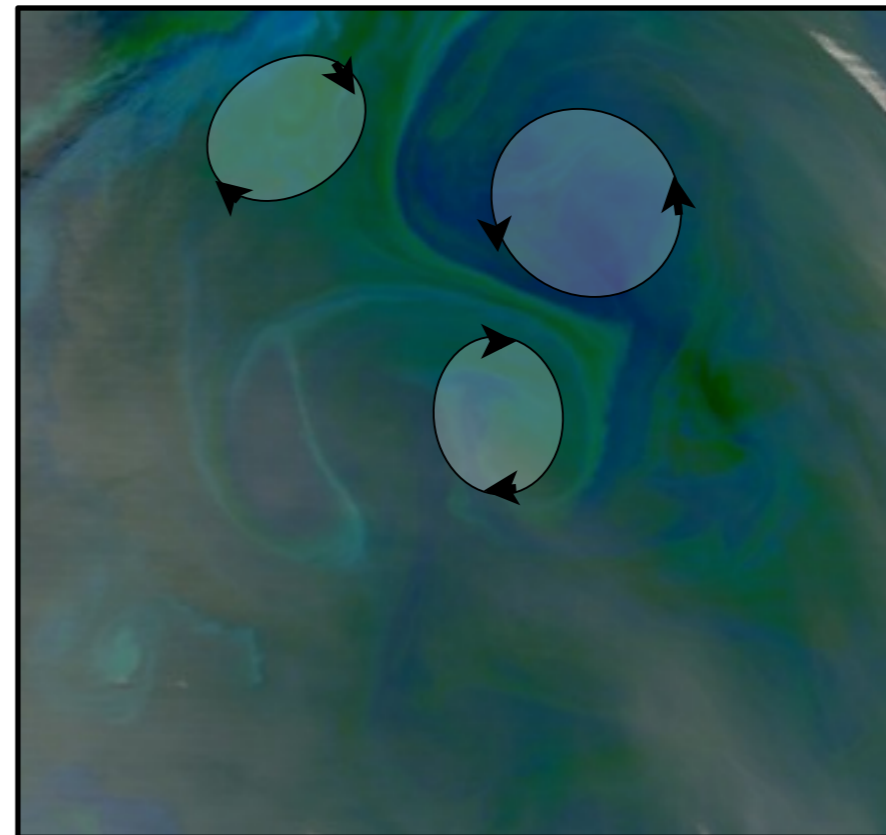
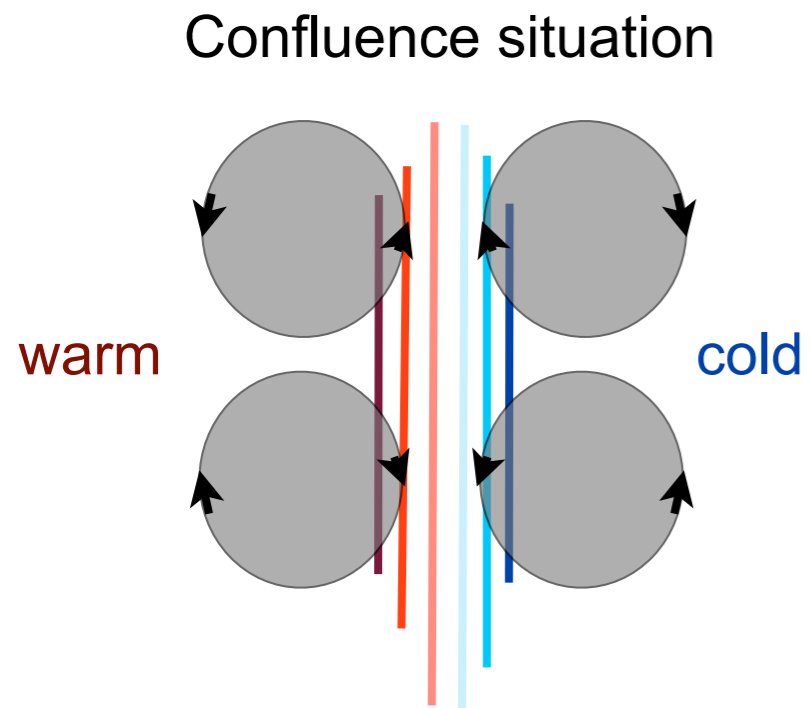
2 essential elements for submesoscale turbulence:

- frontogenesis
- flow instabilities (baroclinic)

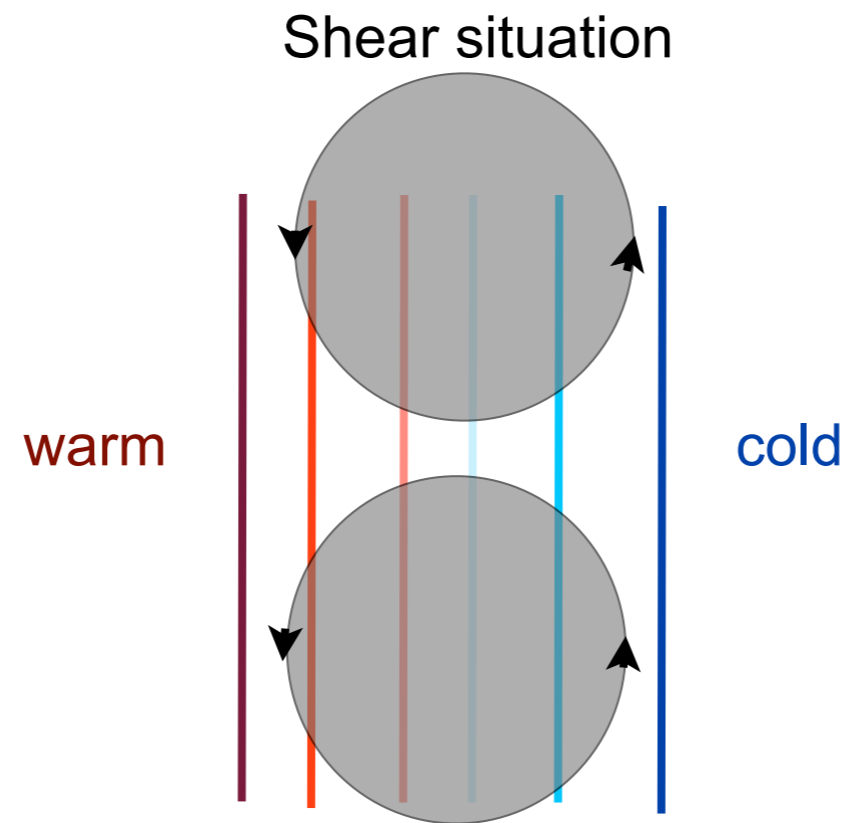
Confluence situation



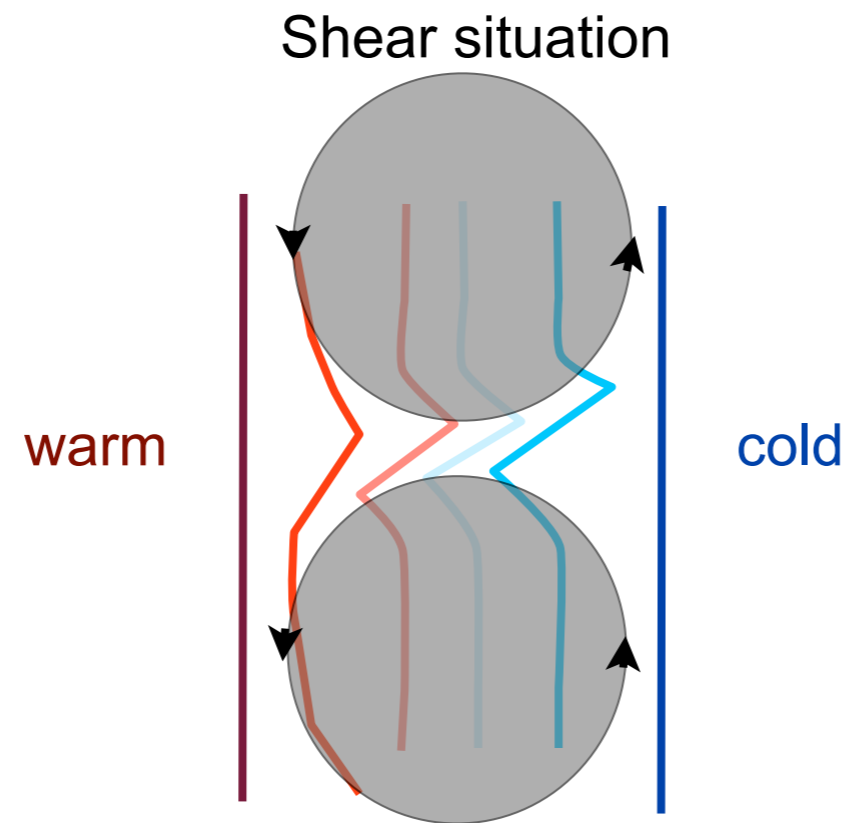
Kinematic argument: in a turbulent flow, **confluence** situations (due to mesoscale eddies) are common and will tend to increase preexisting tracer gradients.



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Kinematic argument: in a turbulent flow, **shear** situations (eg, due to mesoscale eddies) are common and will also tend to increase preexisting tracer gradients.

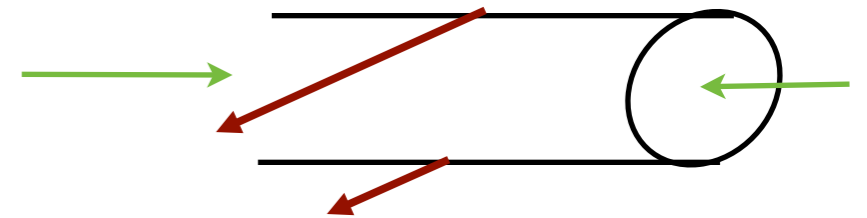
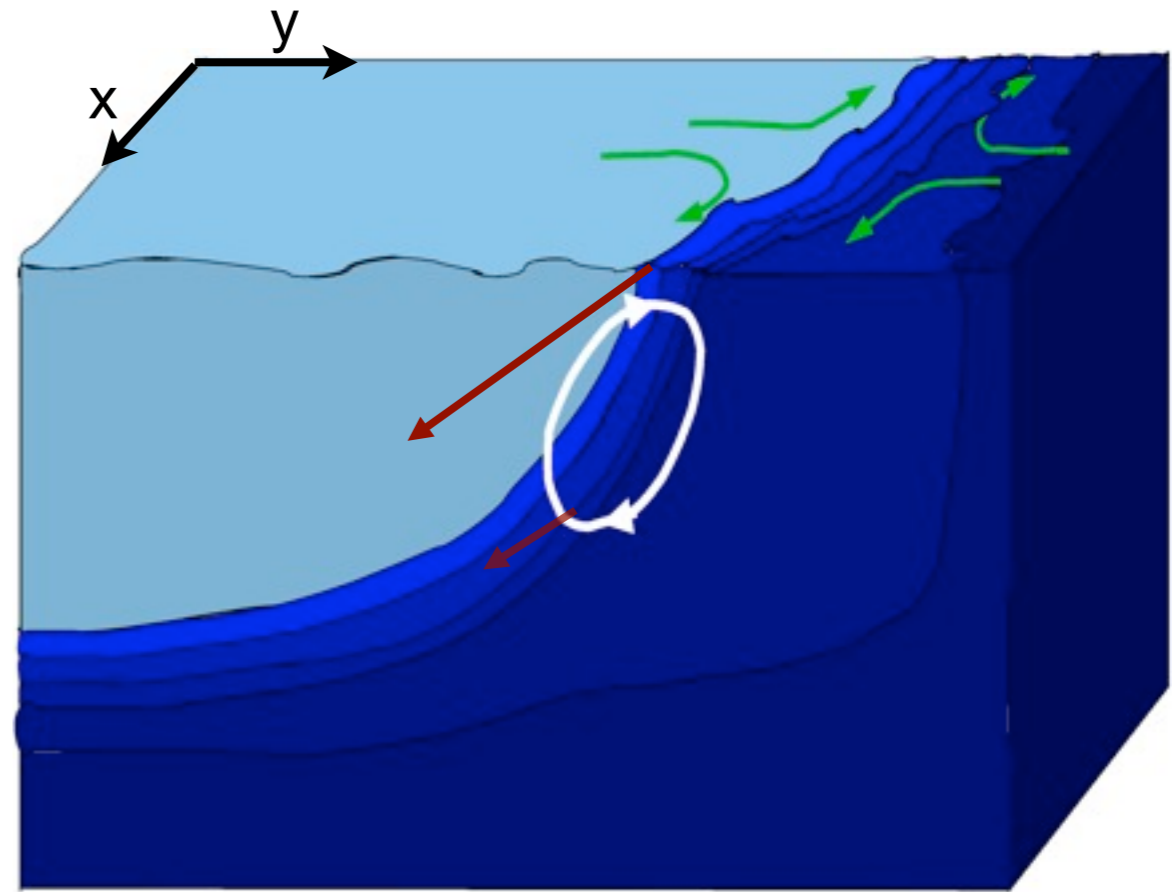


Kinematic argument: in a turbulent flow, **shear** situations (eg, due to mesoscale eddies) are common and will also tend to increase preexisting tracer gradients.

Dynamic argument for geophysical fluids in which Coriolis-pressure force balance is important:

$$f \frac{\partial u}{\partial z} = - \frac{\partial b}{\partial y}$$

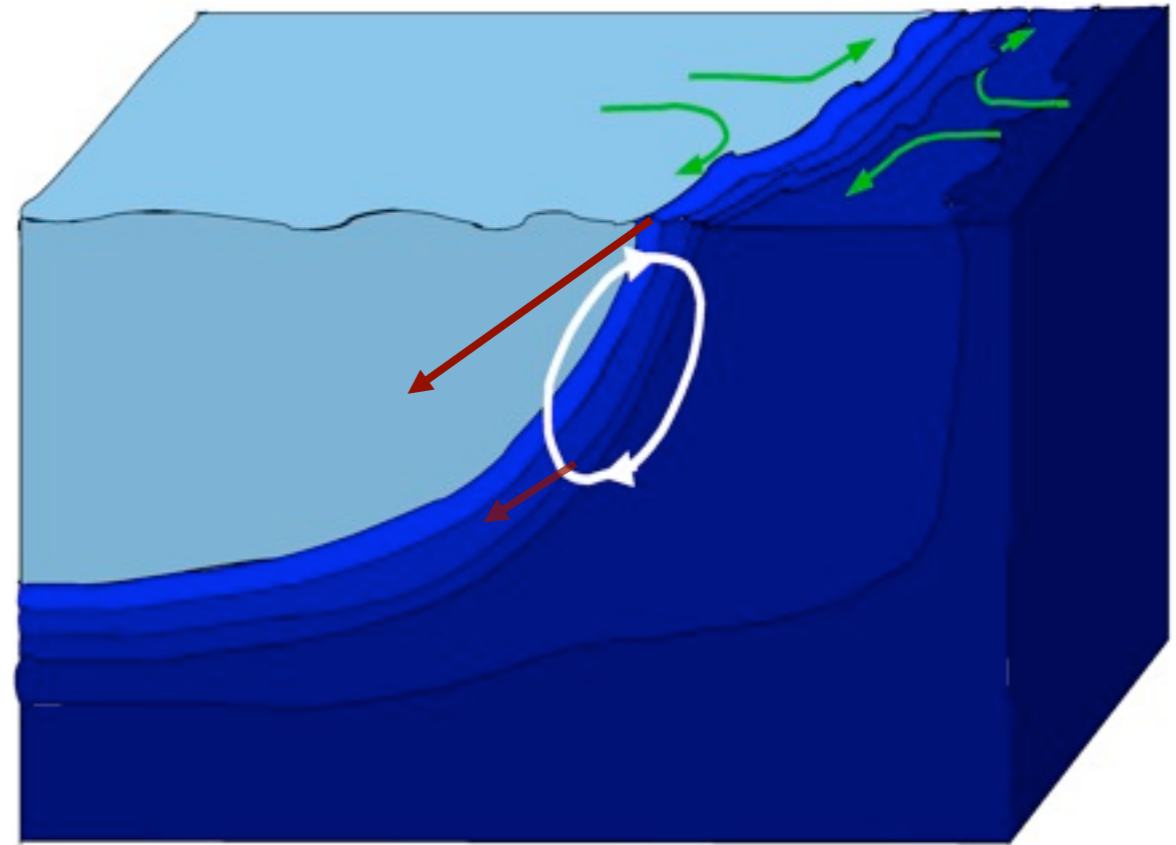
$$f \frac{\partial u}{\partial z} \searrow \quad - \frac{\partial b}{\partial y} \nearrow$$



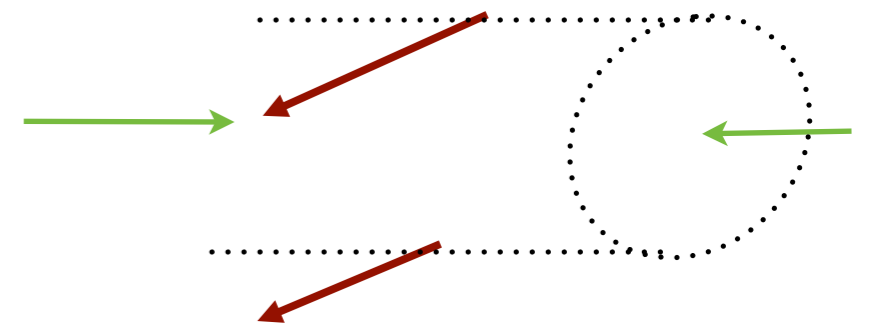
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Mesoscale strain tends to destroy thermal wind balance. Thermal wind balance disruption leads to the development of an **ageostrophic secondary circulation** (ASC) which limits imbalance by restoring shear and limiting frontal intensification.

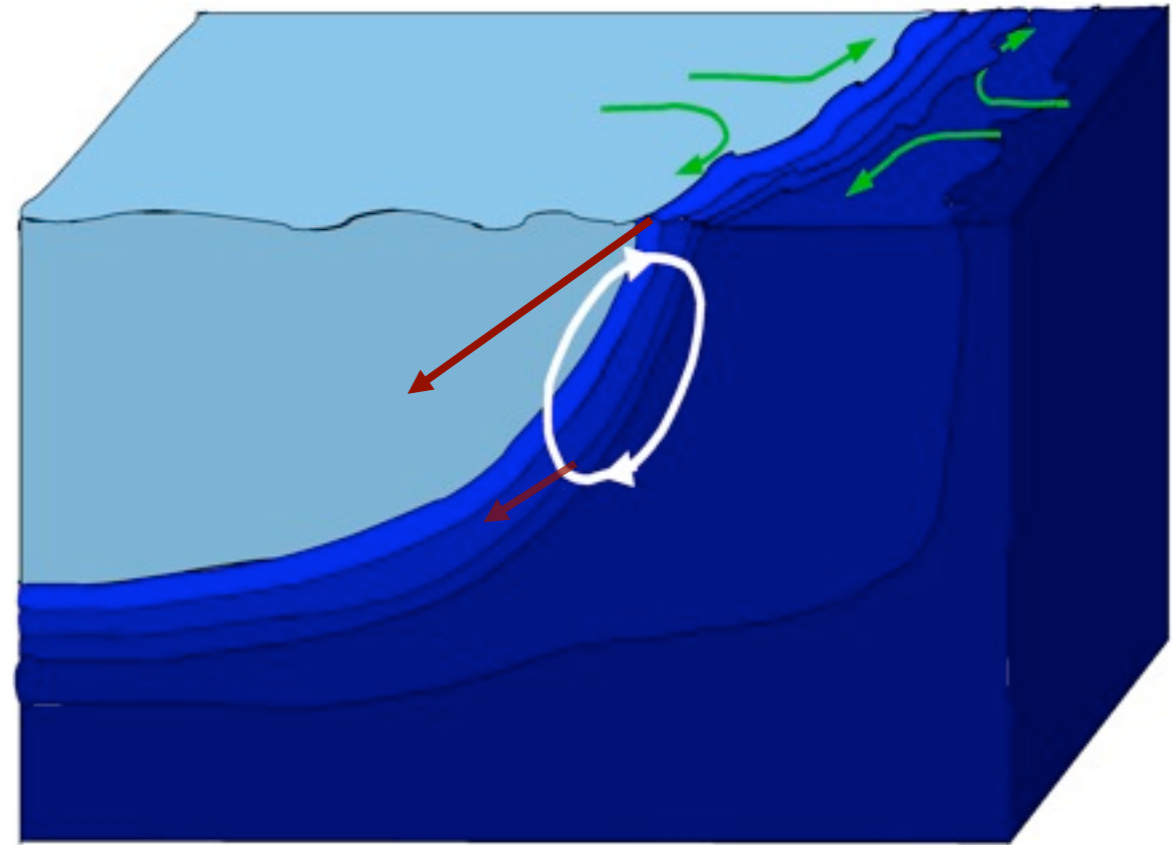


At the surface $w \sim 0$ limits the efficiency of ASCs \rightarrow frontal intensification is difficult to halt \rightarrow intensification of ASCs with potentially large w .

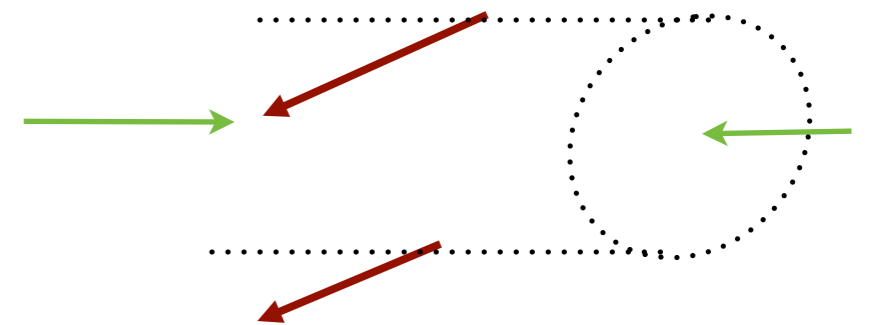
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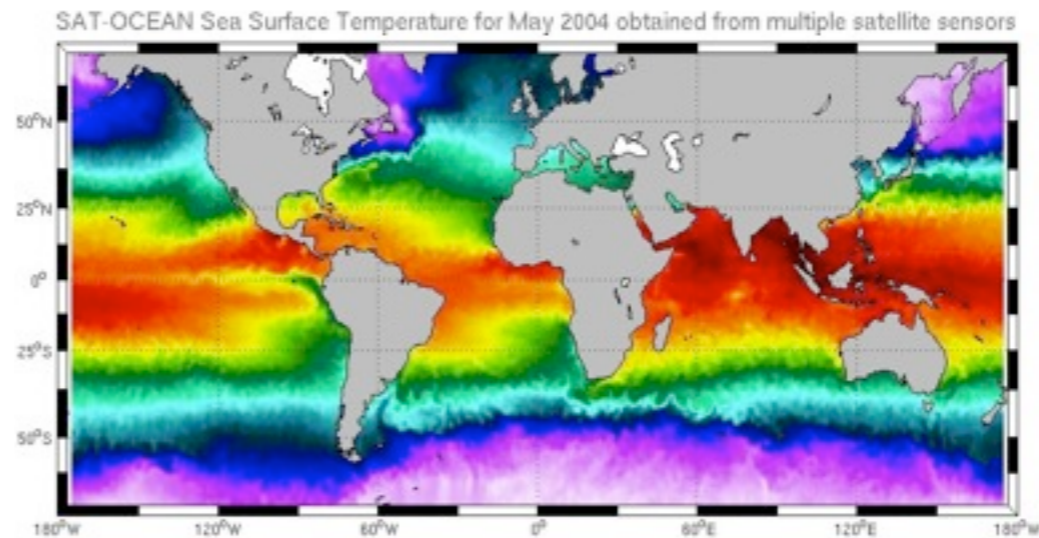
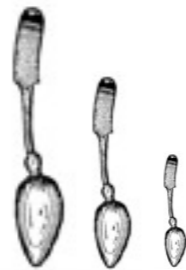


Vertical velocities develop that are concentrated within narrow frontal regions. Nonlinearities tend to enhance/concentrate the downward branch and weaken/broaden the upward branch. **The processes that limit the magnitude and the shrinking of cross-front length scale are still being investigated.**

Ingredients for frontogenesis:

Surface density gradients

Stirring spoons

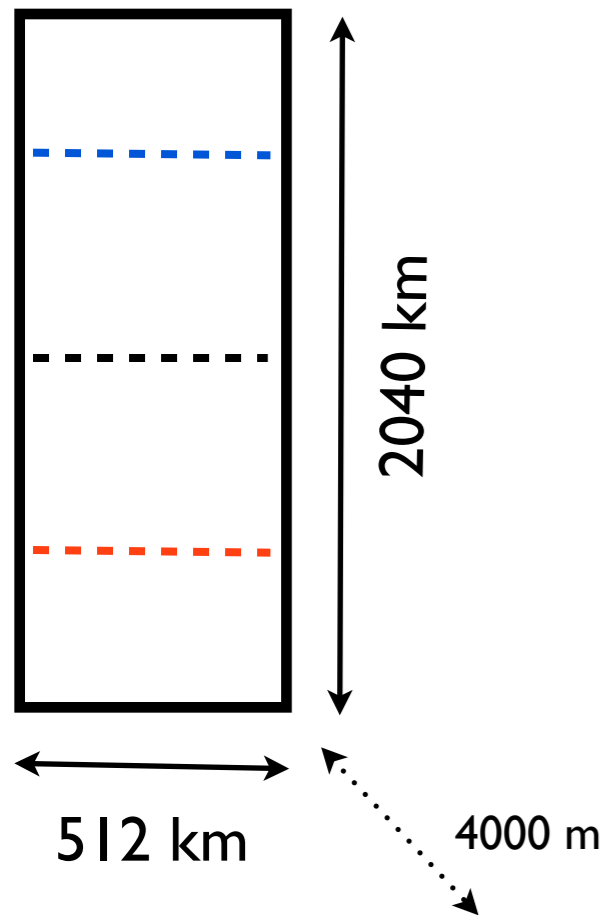


Purpose of this study: present and compare two situations with different distributions of stirring spoons associated with different flow instabilities.

2 essential elements concepts to understand submesoscale turbulence:

- frontogenesis
- flow instabilities (baroclinic)

Reentrant channel with a baroclinic jet (3D restoring of zonally averaged u and ρ)

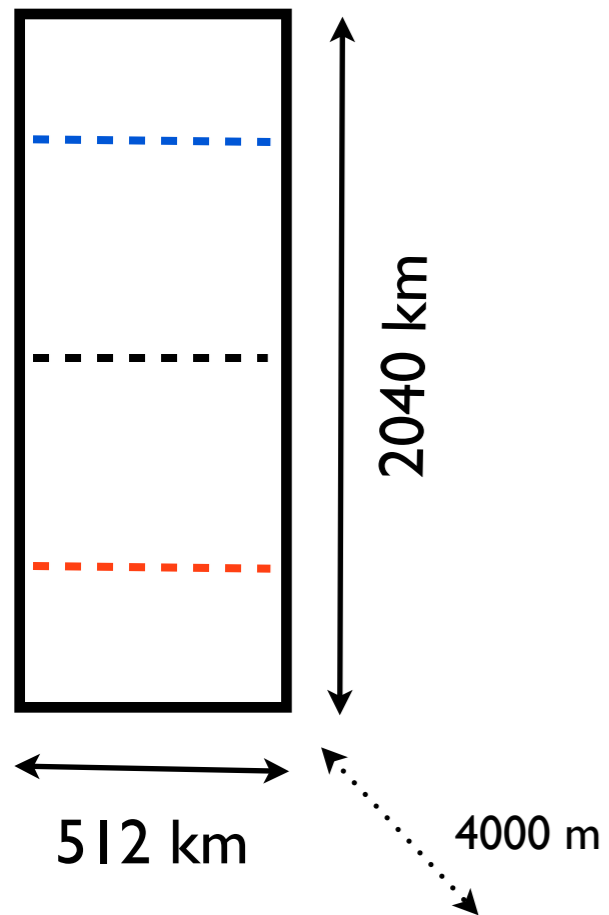


512 x 2040 x 200 grid points.

1 month: 9h wall clock time of 64 Sandy Bridge nodes (2 x 8 cores/node, 2.3 GHz)

ROMS UCLA, Ri based vertical diffusion.

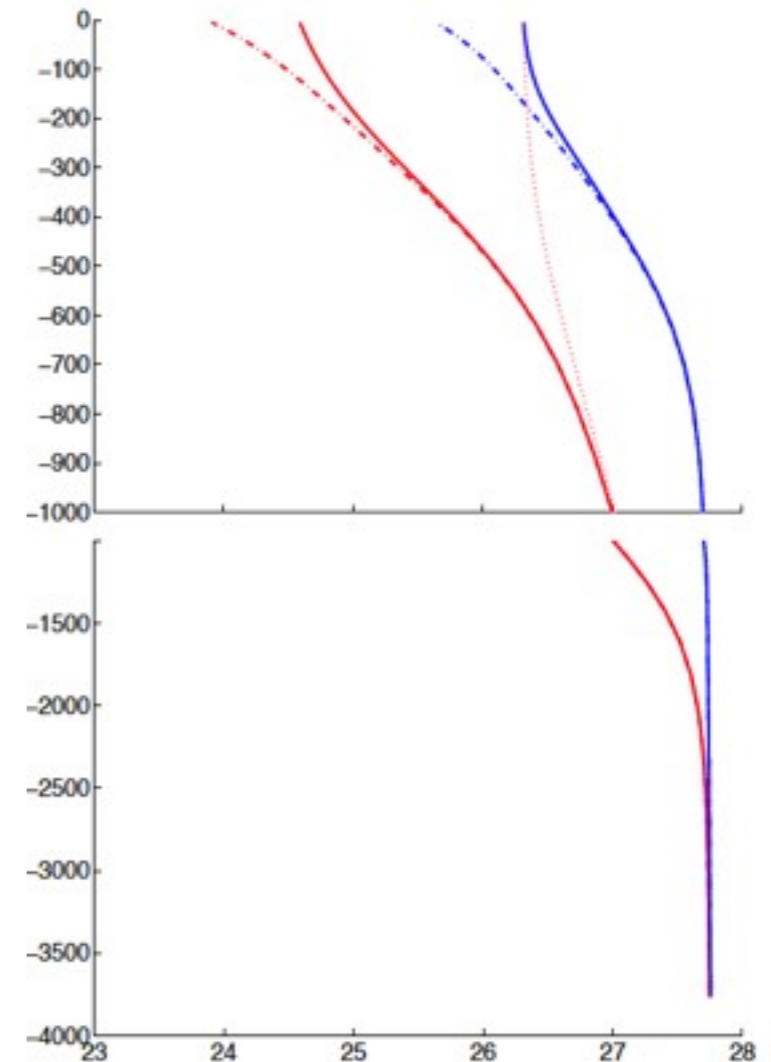
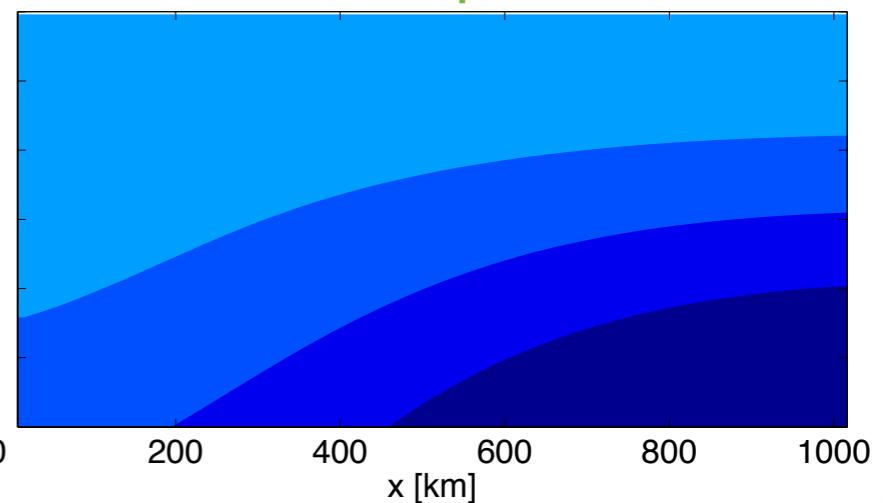
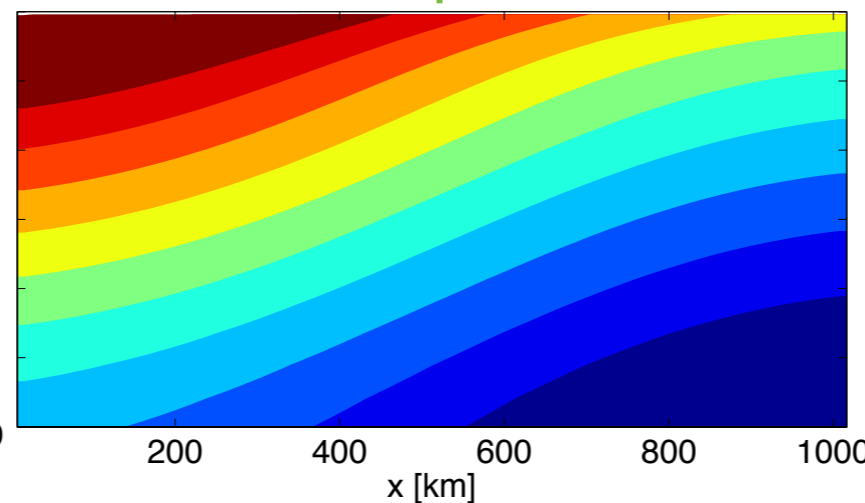
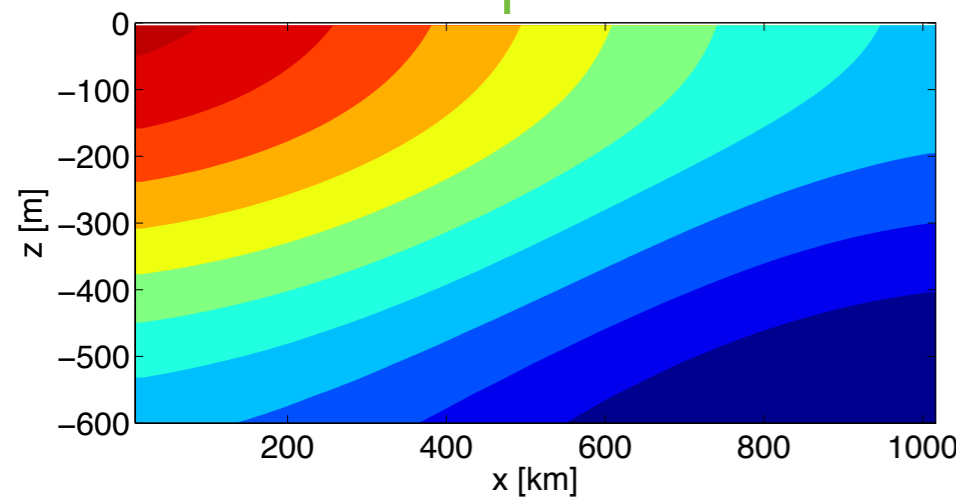
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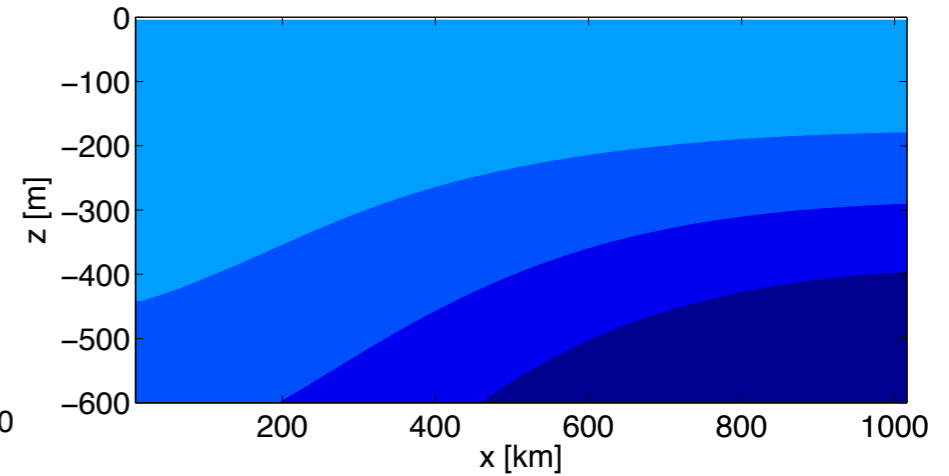
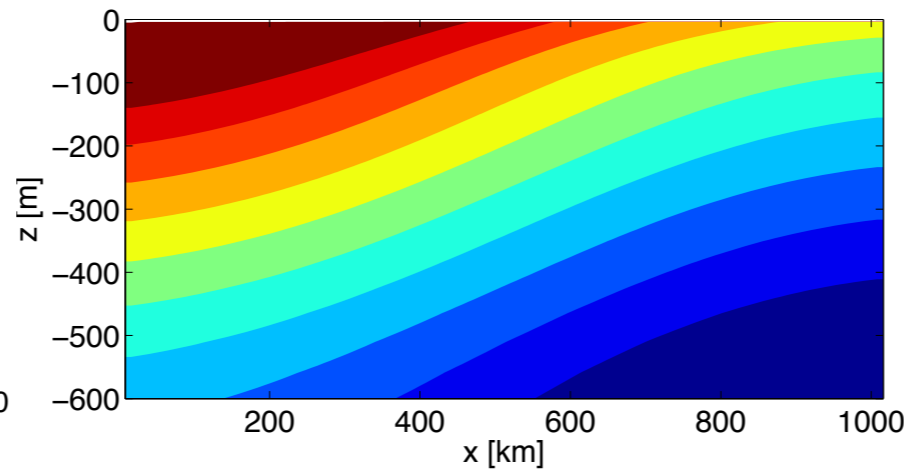
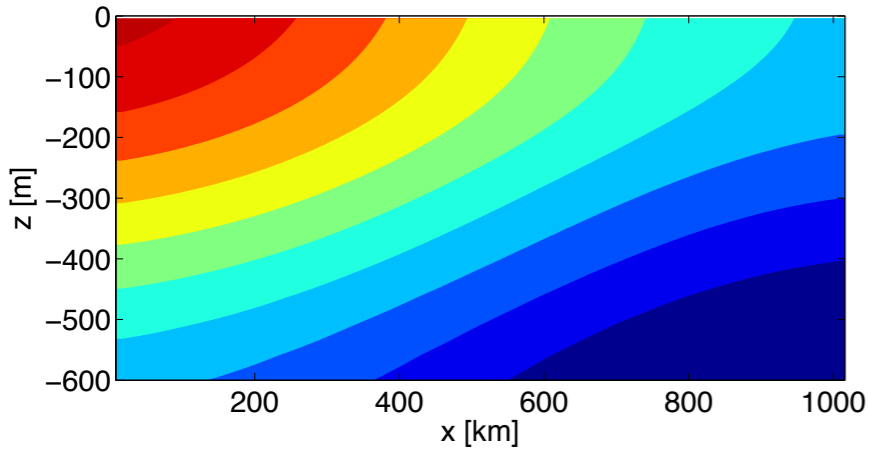


setup 1

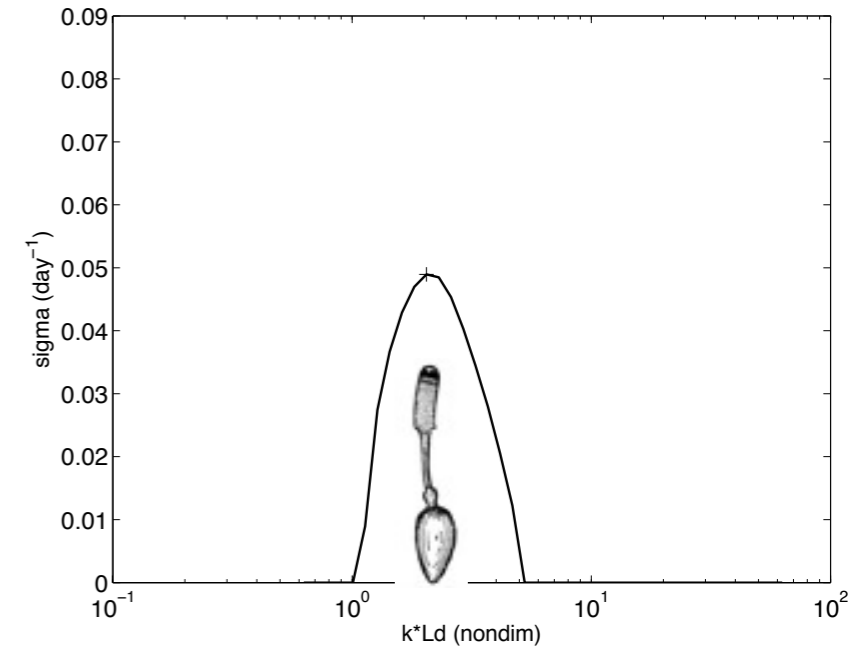
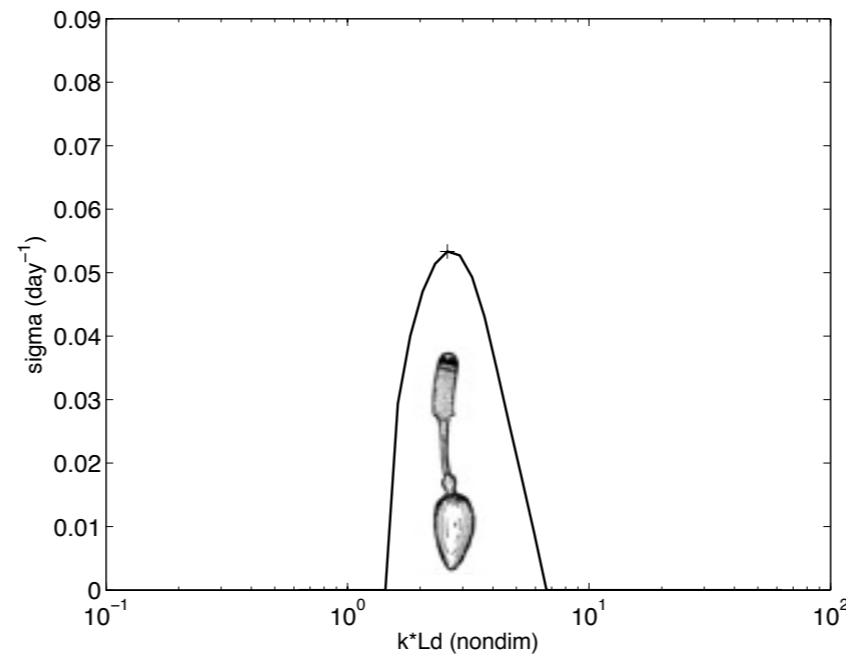
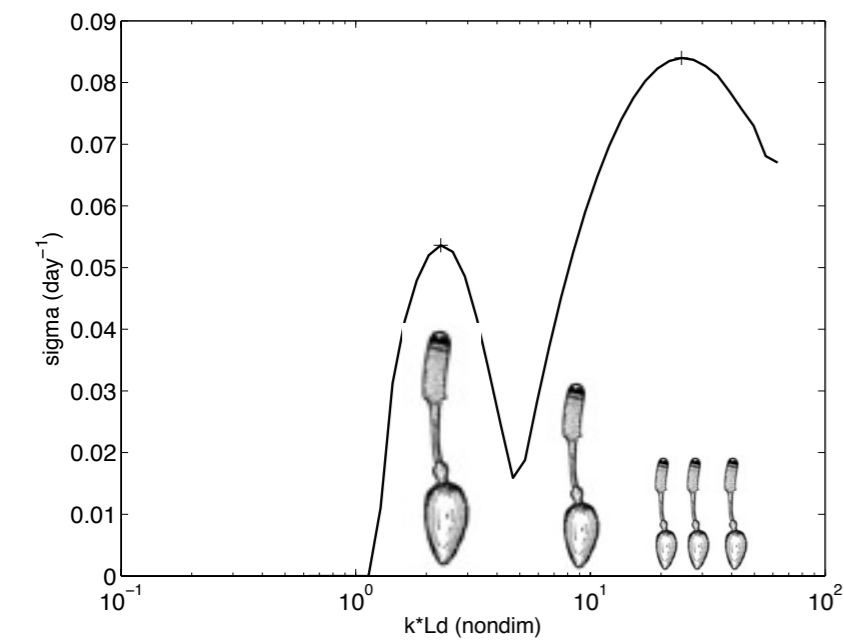
setup 2

setup 3





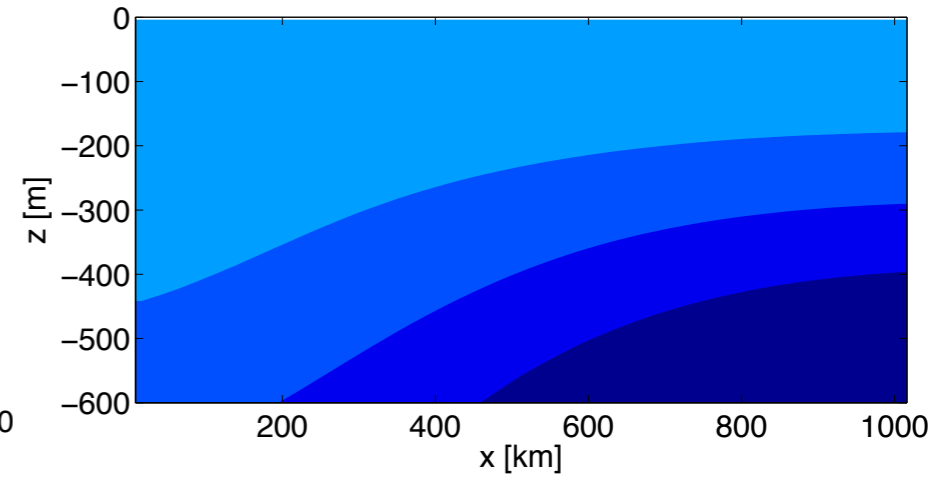
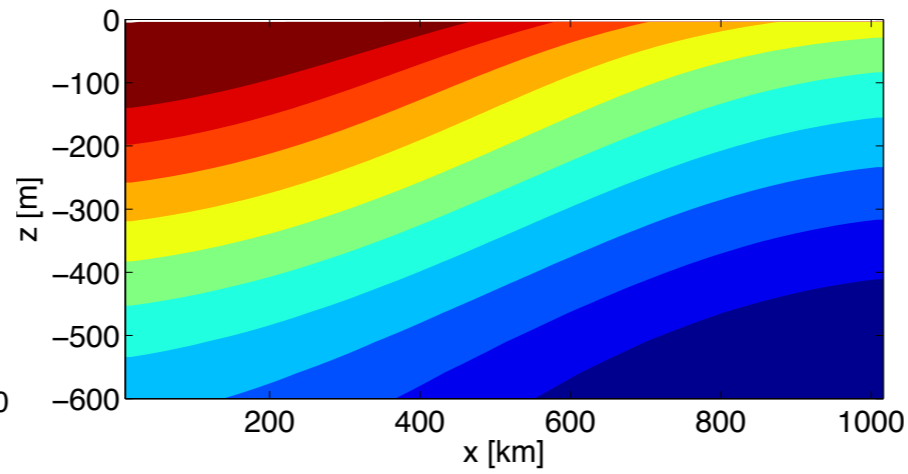
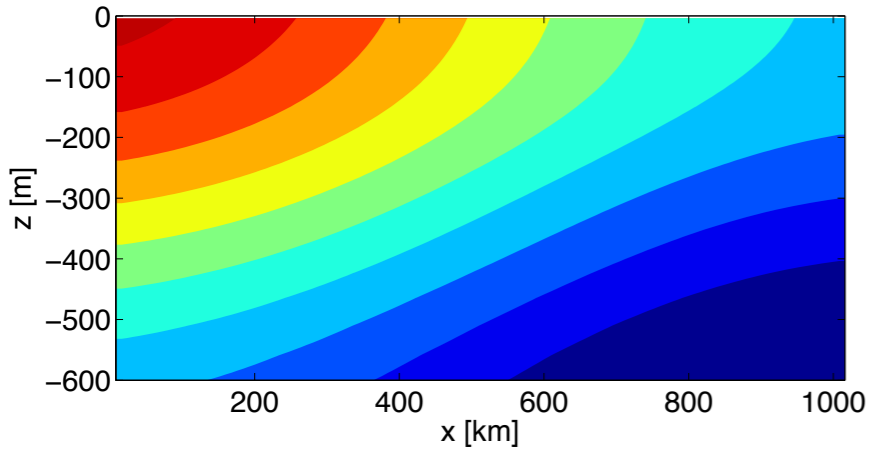
Linear instability growth rates



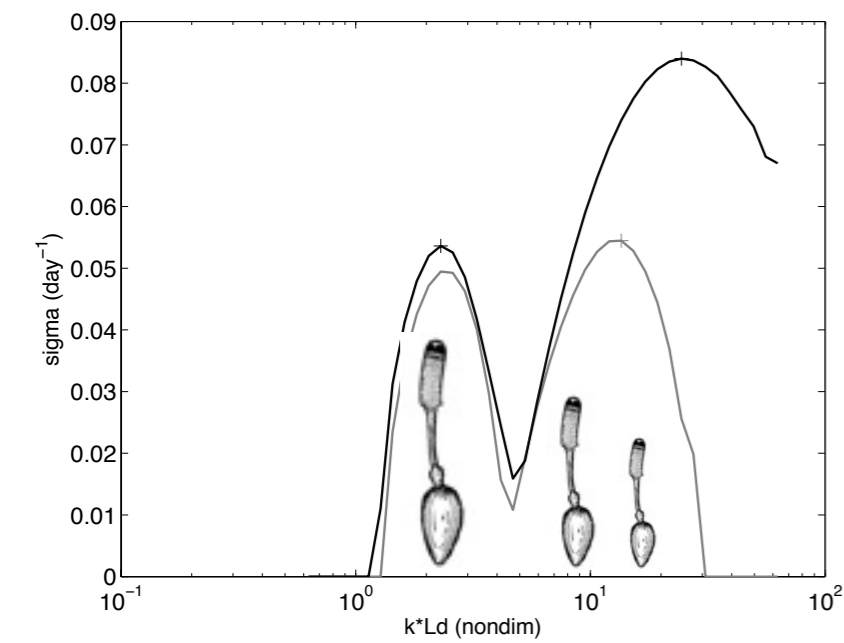
*classical interior BCI +
surface density gradient +
Charney-like instability mode*

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surface density gradient*

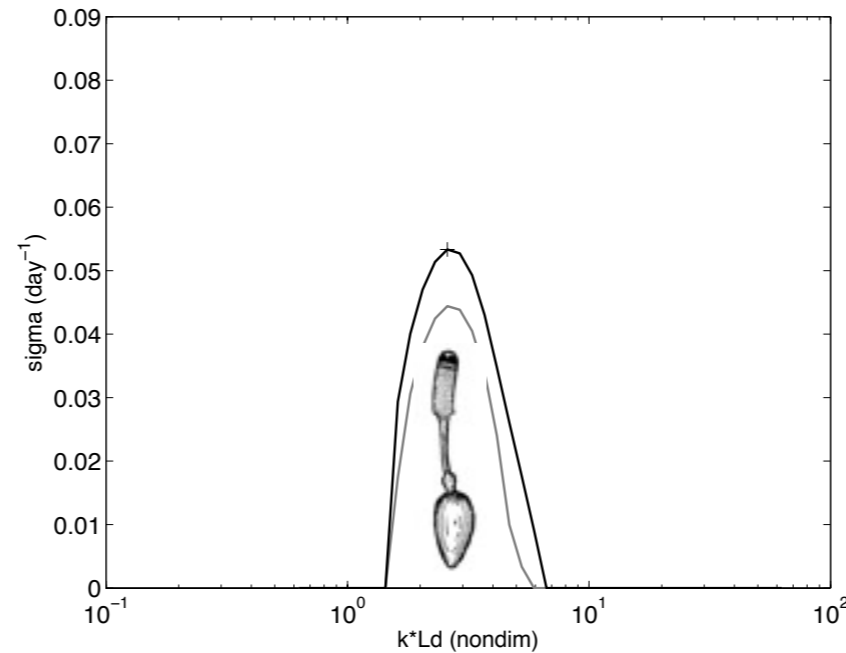
classical interior BCI



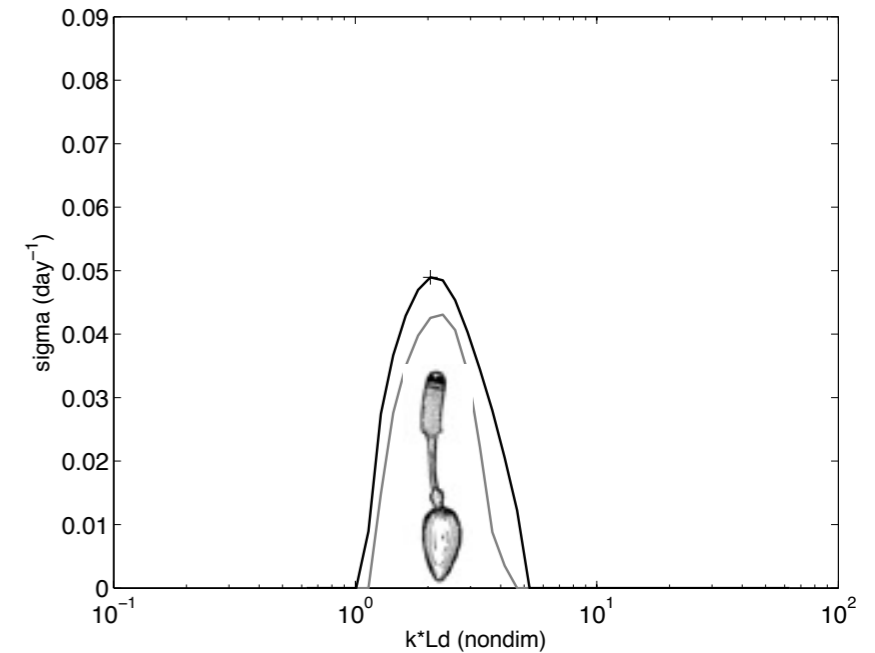
Linear instability growth rates



*classical interior BCI +
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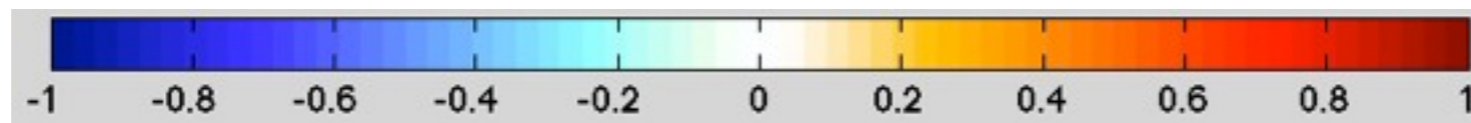
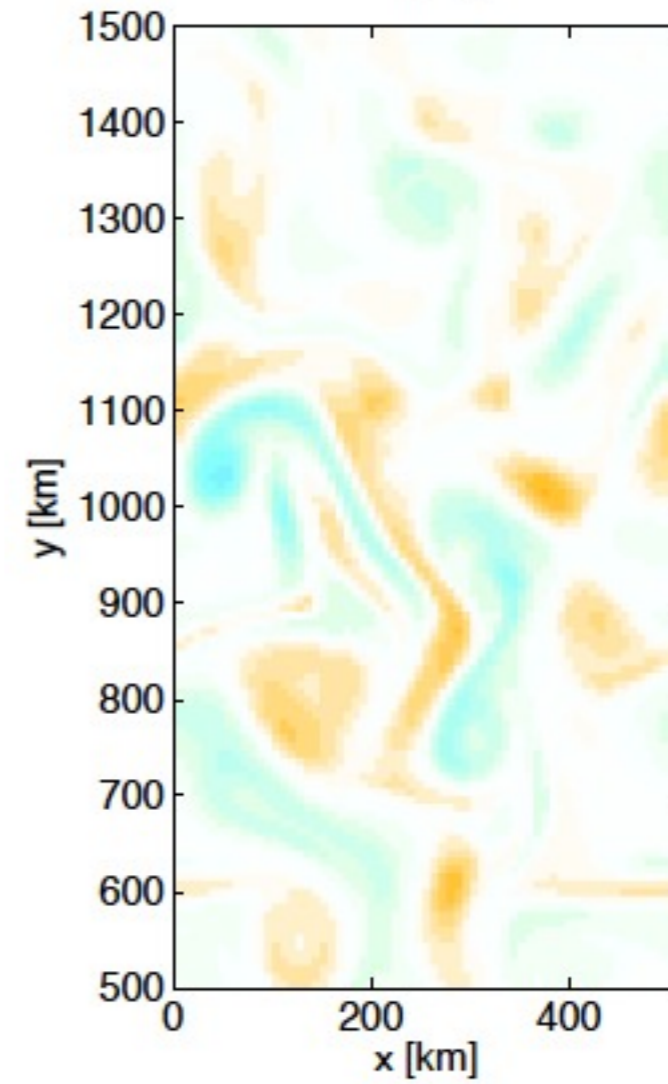
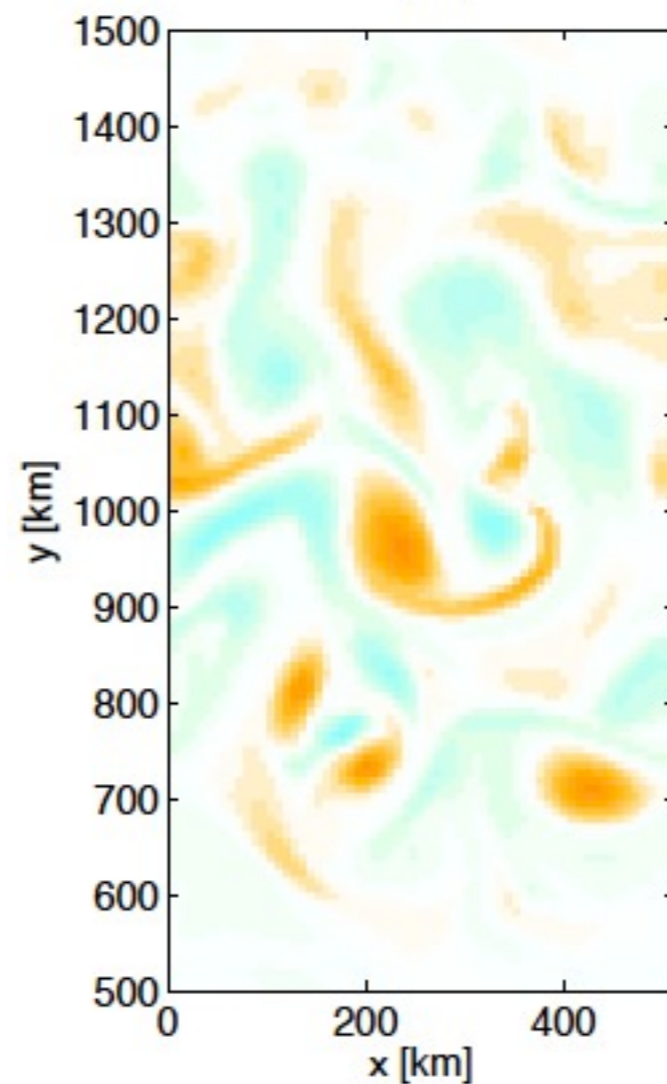
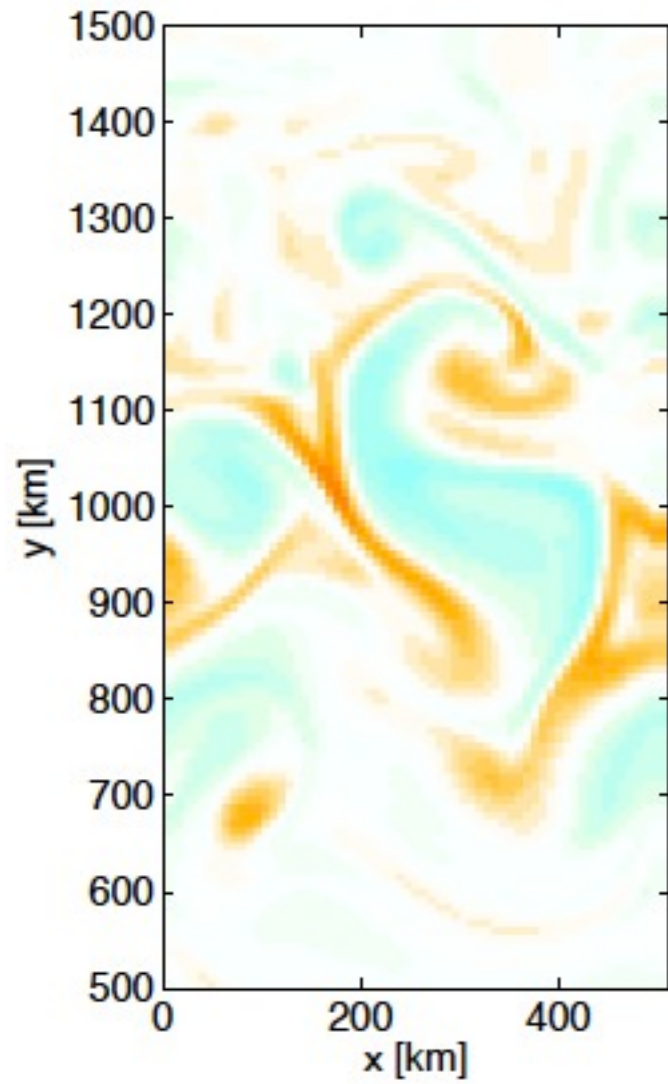
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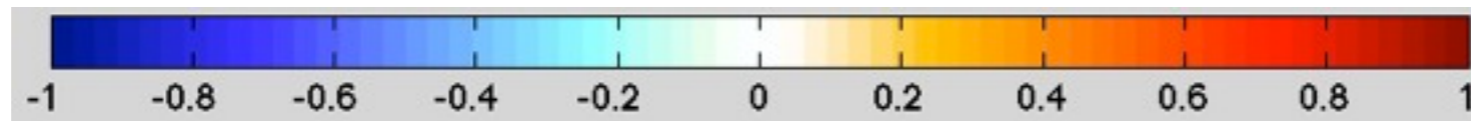
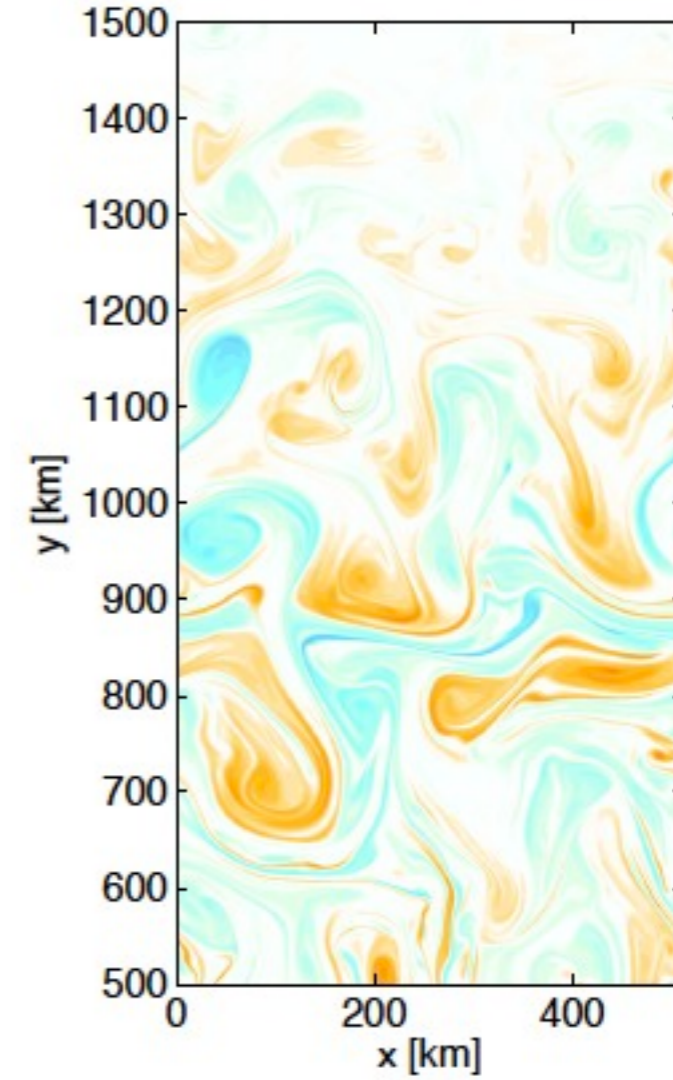
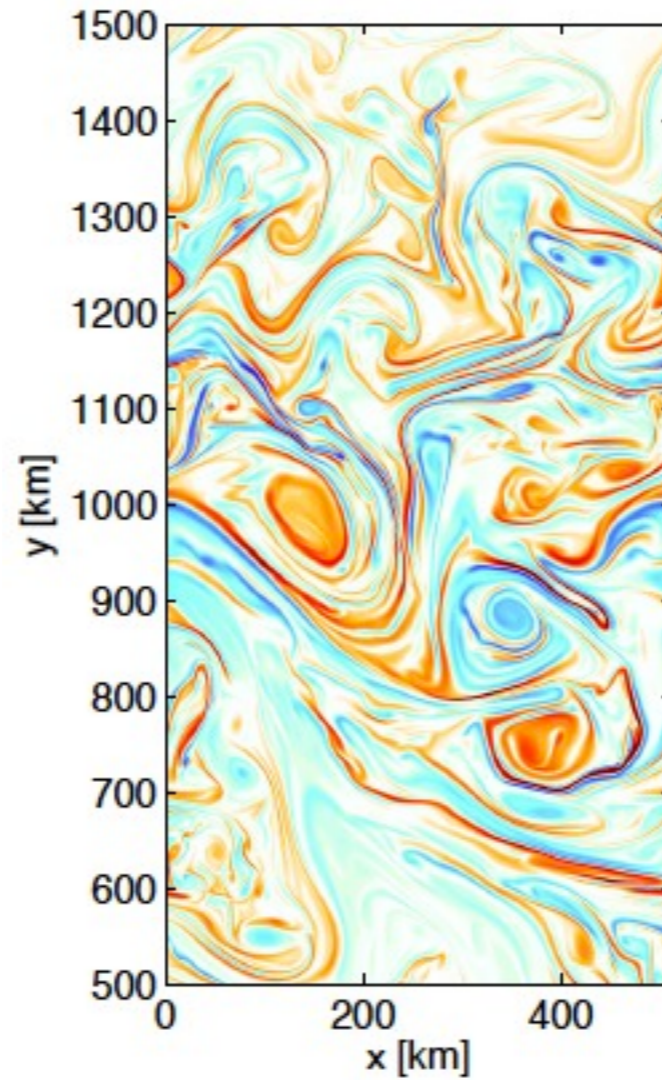
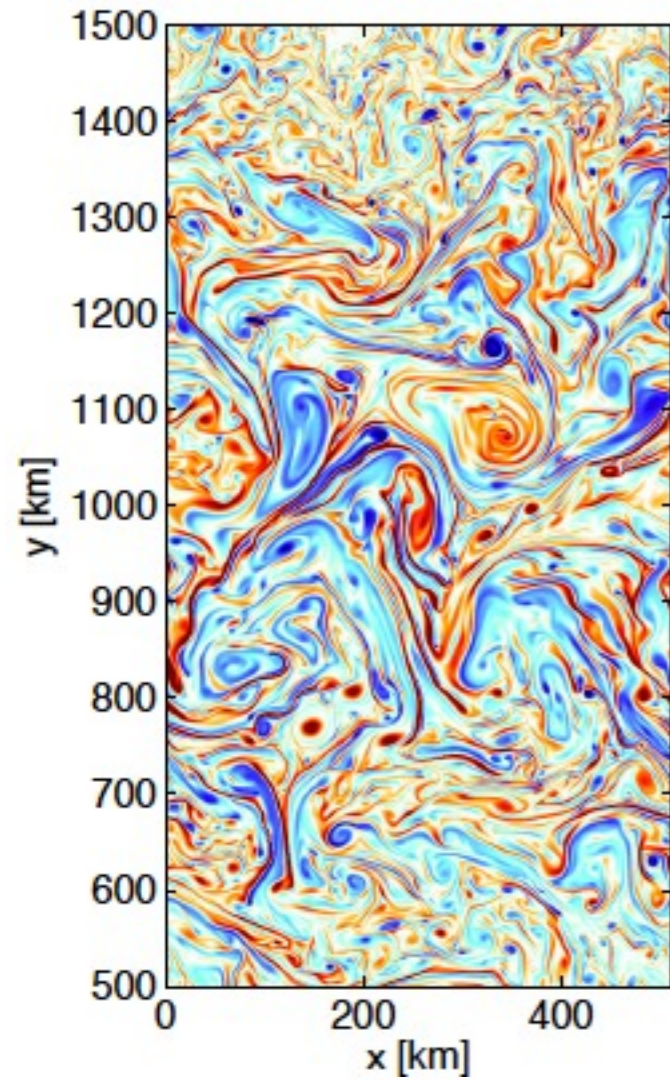
classical interior BCI

The Charney mode arises from the simultaneous presence of a i) poleward surface density gradient ii) a weakly stratified upper ocean

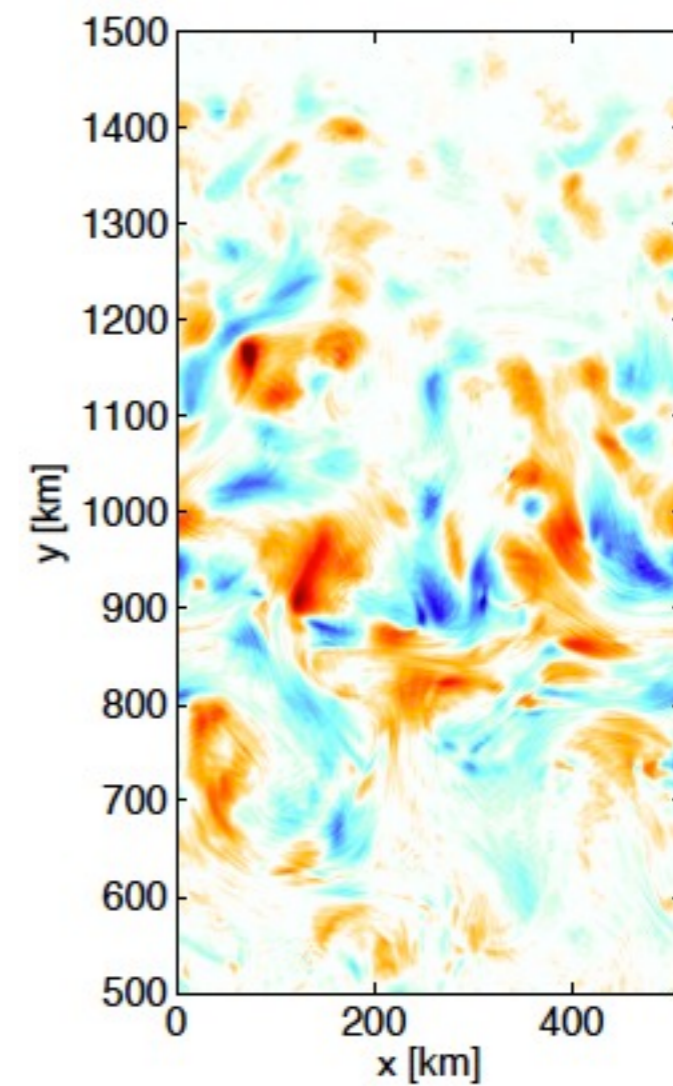
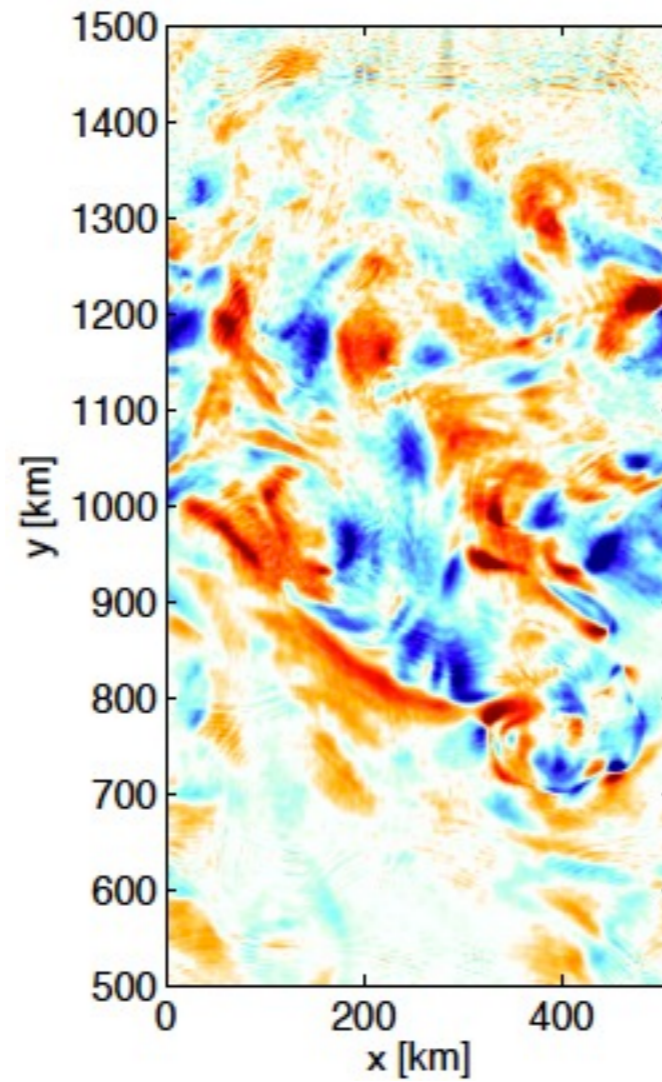
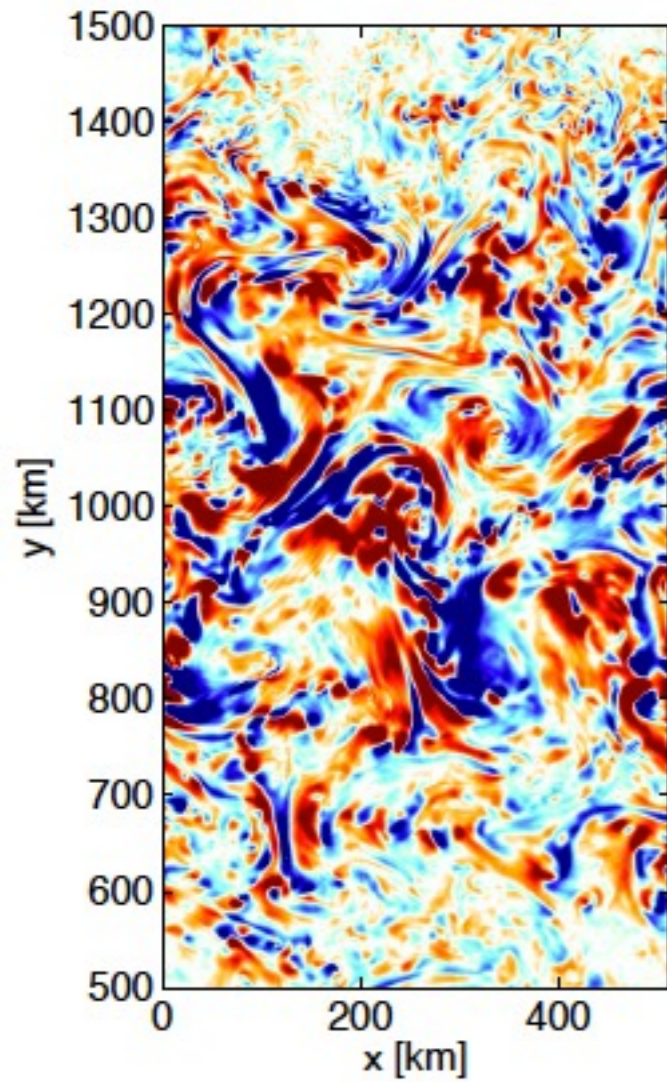
Surface ζ - 8 km



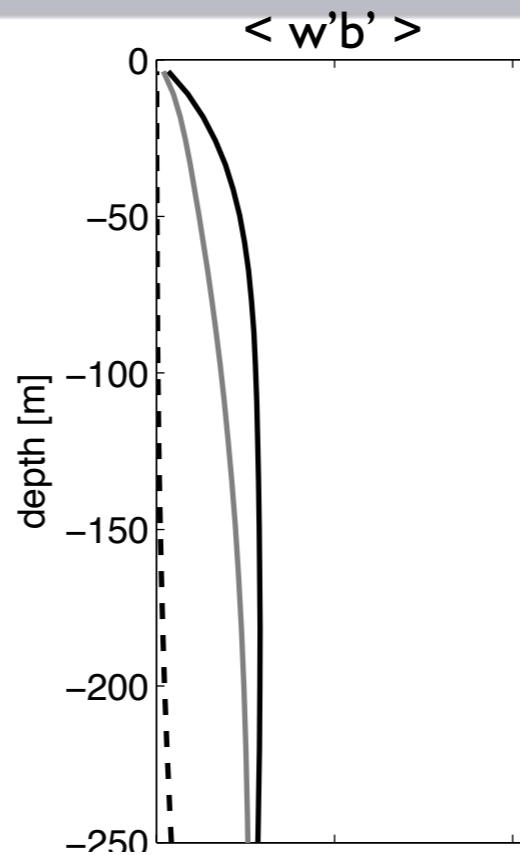
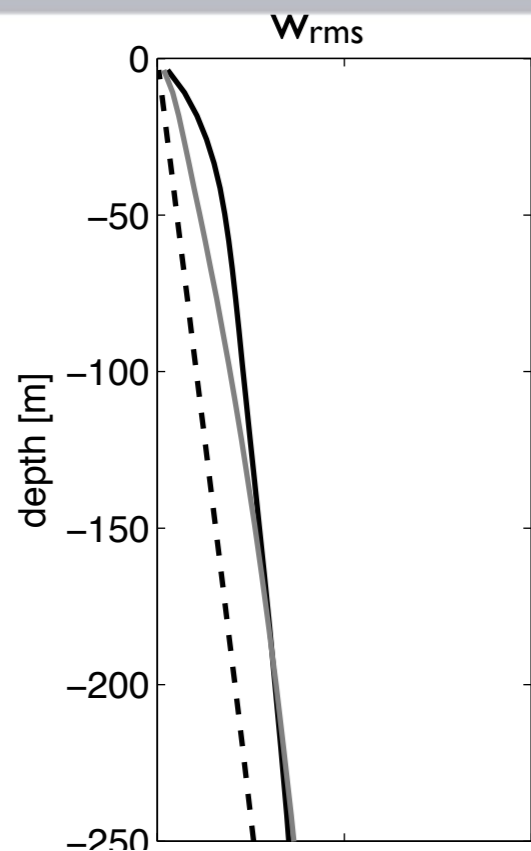
Surface ζ - 1 km



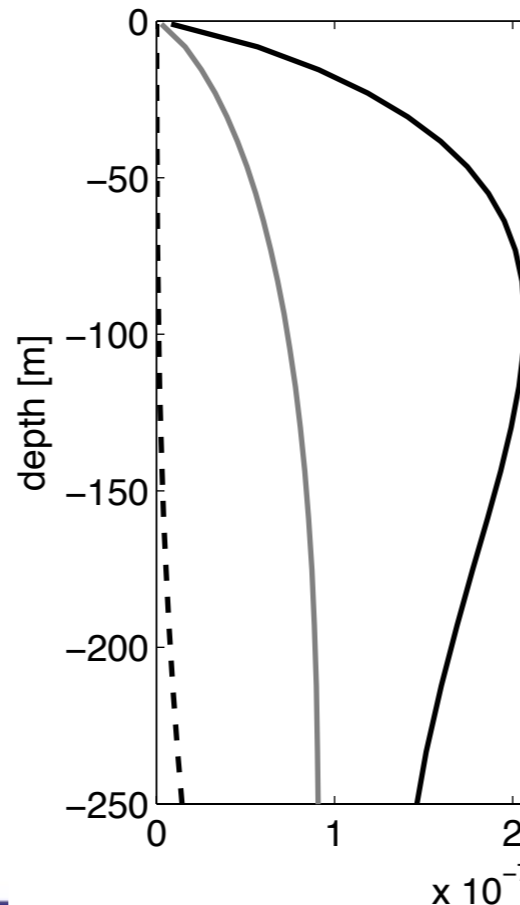
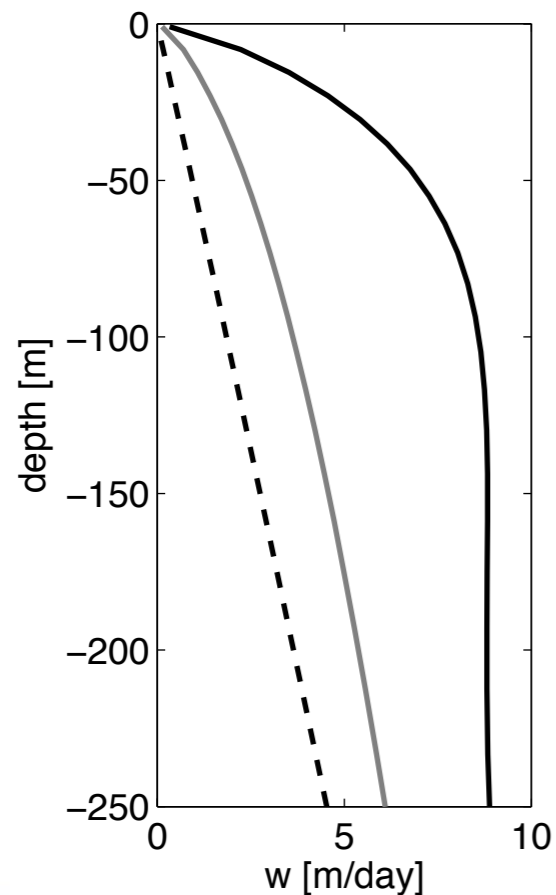
120 m w - 1 km



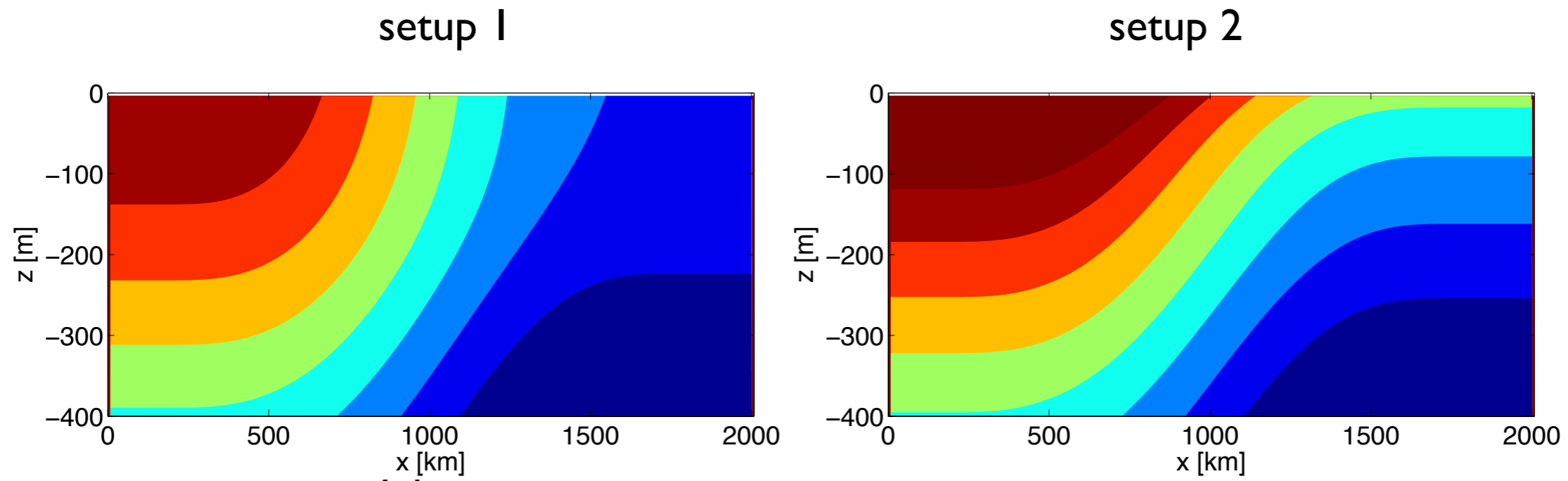
$dx = 8 \text{ km}$



$dx = 1 \text{ km}$



Setup 3 is not sensitive to resolution between beyond mesoscale resolving. Very large increase in w_{rms} and $\langle w'b' \rangle$ for setup 1 with increasing resolution; much less dramatic for setup 2.



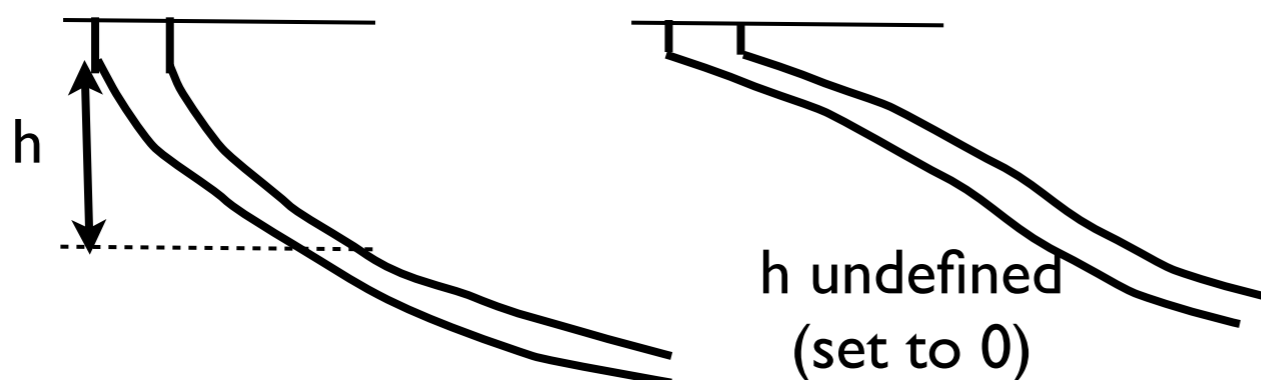
$$\partial_y q = \beta - f \partial_z s + \left[\frac{gf}{\rho_0 N(0)^2} \partial_y \rho \delta(0) \right]$$

Charney-Stern criterion: $\partial_y q$ must change sign (in the vertical for baroclinic instability)

- The Charney mode arises from the simultaneous presence of
- a poleward surface density gradient
 - a weakly stratified upper ocean (precisely an increasing density slope)

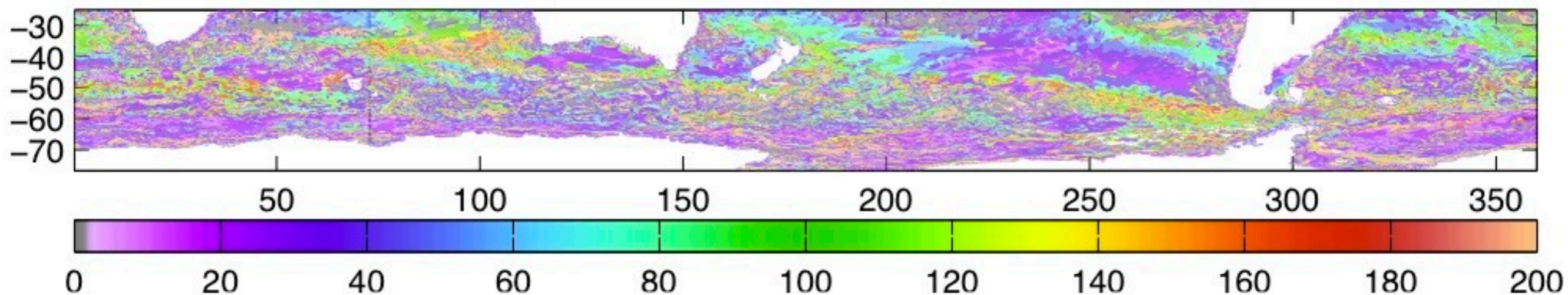
- The Charney mode arises from the simultaneous presence of
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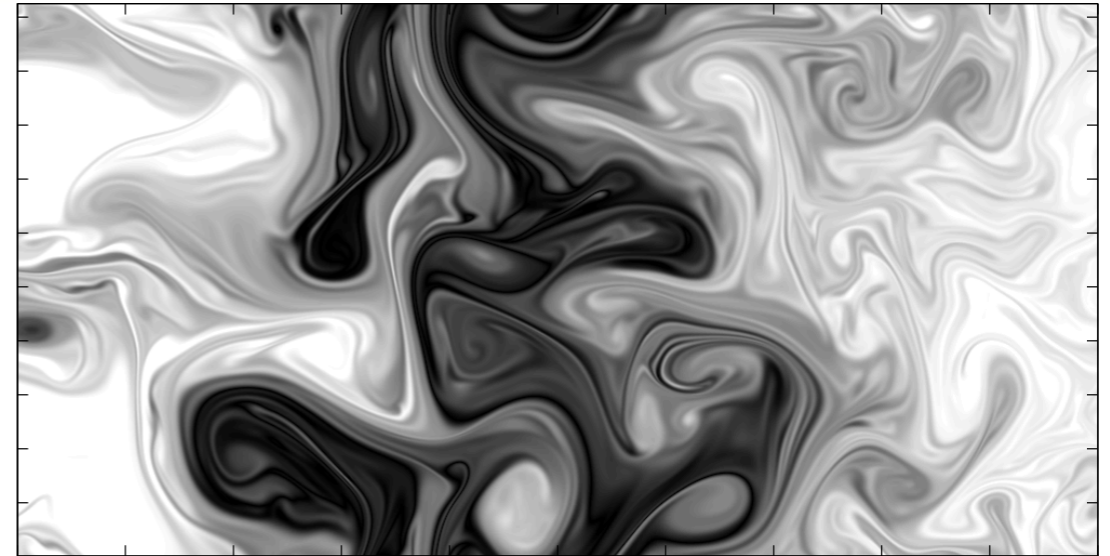
We search for thick subsurface layers close to the surface and in which $ds/dz > 0$ in the southern hemisphere



We find them in austral spring in places of mode water formation.

y1999m11d21 - 65





View of ocean turbulence until ~ 2000: there are mainly big spoons in the ocean. SST/density are passive tracer mainly stirred by these big spoons → tracer filaments wrapped around mesoscale eddies with no dynamical importance.

Present view: SST/density are active tracers that help energize the flow at scale smaller than the mesoscale through several processes:

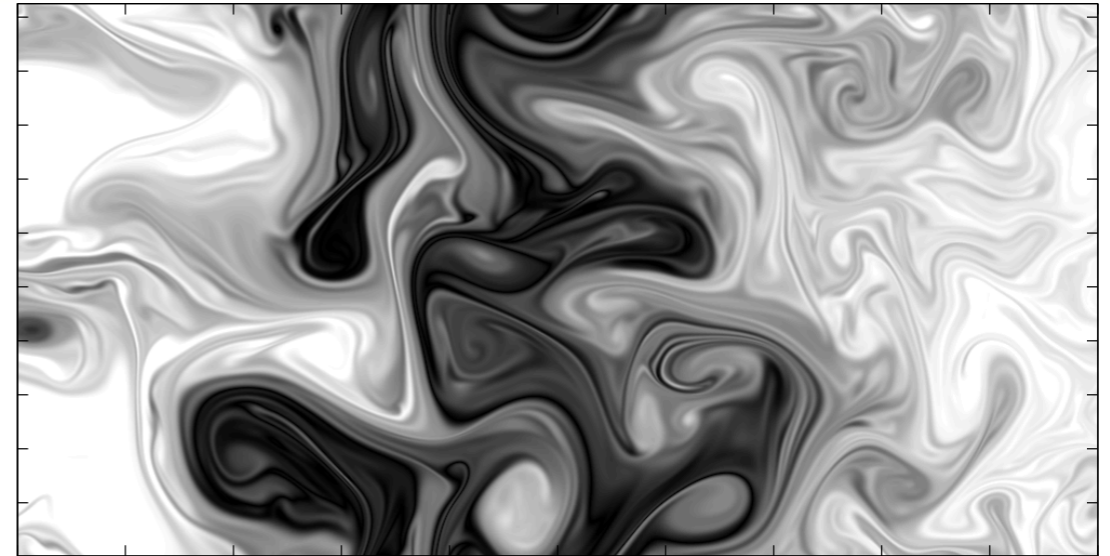
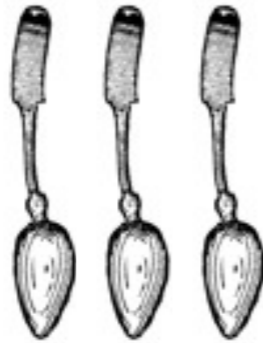


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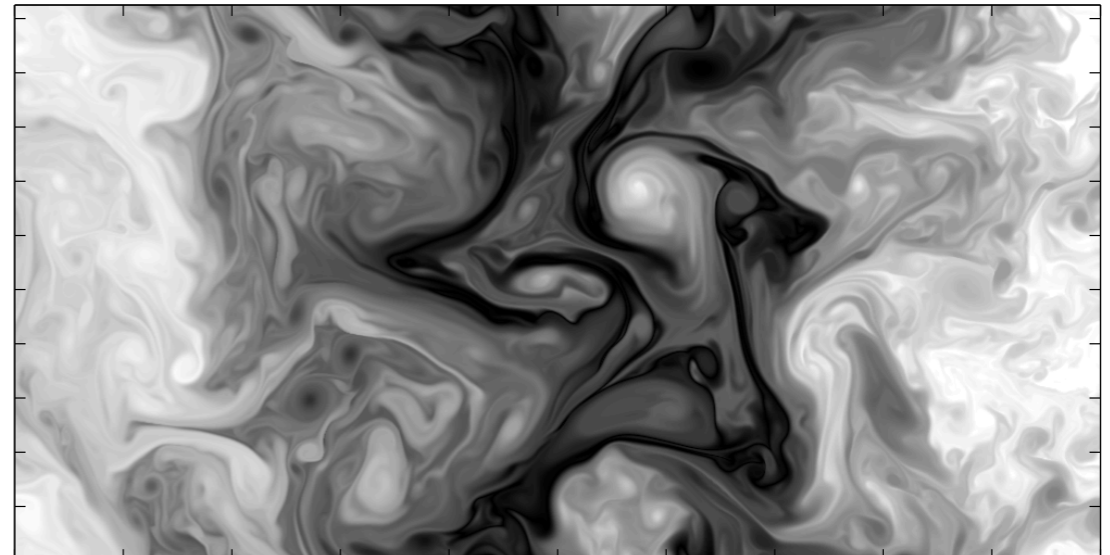
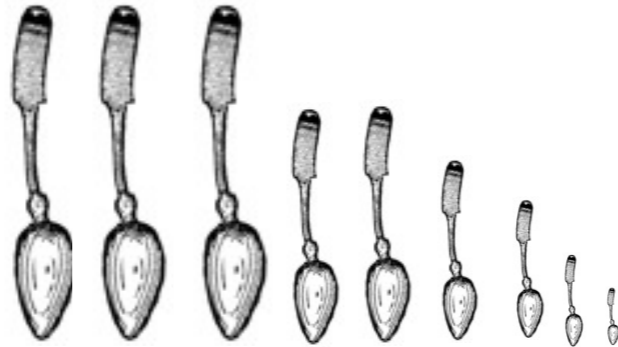
Present view: SST/density are active tracers that help energize the flow at scale smaller than the mesoscale through several processes:

- mesoscale driven frontogenesis
- $O(1)$ km scale linear instabilities (MLI, **Charney type BCI** \rightarrow symmetric, inertial, shear ...)
- wind can contribute to submesoscale energization in subtle ways (see Thomas, 2008, 2010 ...)

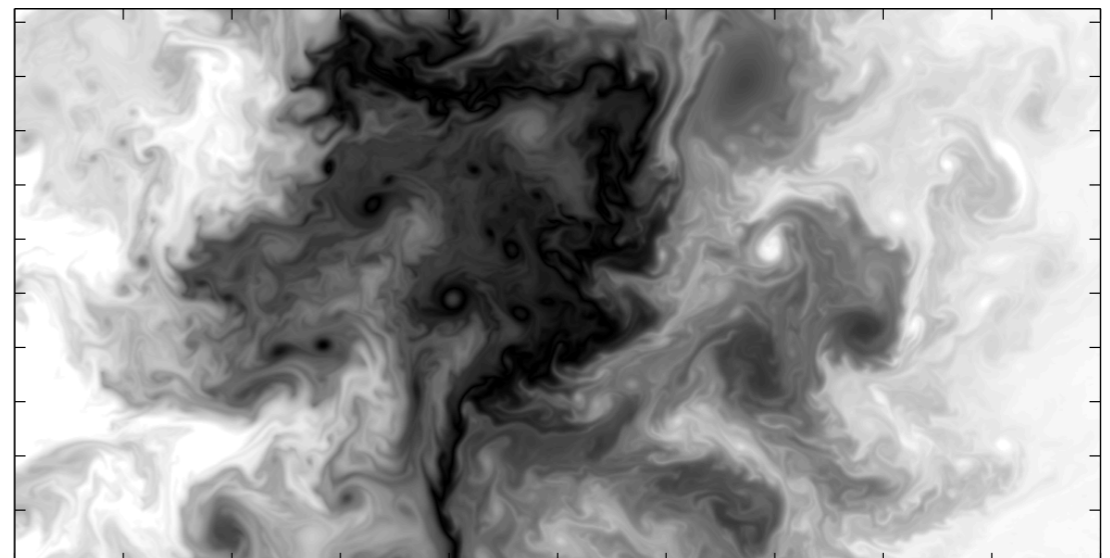
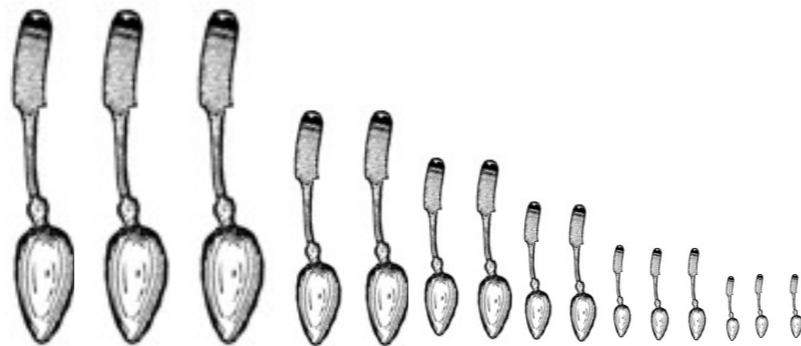
Interior BCI



Interior BCI + Charney



Interior BCI + MLI



- Intense lateral density gradients are no guarantee for large vertical fluxes of tracers (and neither is resolution increase in numerical models)
- Subsurface stratification is a key parameter controlling these fluxes. One important way it does so is by controlling the existence and intensity of small scale instability modes (eg, Charney).
- Mode water regions are good candidates for Charney type instability and thus for effective connections between the subsurface and near surface
- MLE (FK08) are only one aspect of the dynamics not resolved in mesoscale-resolving numerical models. → what do we do with other aspects ?
- To be resolved: interactions between submesoscale turbulence and more complex/finer scale environmental processes (Langmuir cells, wind forcing ...)

Difficulty in estimating the role played by submesoscale turbulence in fluxing tracers between the surface layer (typically the mixed layer) and the interior (typically below the nutricline):

- Fluxes are strongly sensitive to the tracer sink and source functions
- Eulerian w are not very informative and often misleading (Lagrangian quantity)
- Fluxes can be very large within the mixed layer but they generally decrease rapidly below its base (\rightarrow ML instability)

We have become accustomed to the fact that submesoscale turbulence and its accompanying near-surface vertical tracer fluxes are highly sensitive to near-surface stratification (i.e. ML depth): MLI intensity scales as h_{bl}^2

(Fox-Kemper et al, 2008)

